

Hidden interactions of quarks

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$$\begin{pmatrix} u_L \\ d_L \\ u_R \\ d_R \end{pmatrix} \quad \begin{pmatrix} c_L \\ s_L \\ c_R \\ s_R \end{pmatrix} \quad \begin{pmatrix} t_L \\ b_L \\ t_R \\ b_R \end{pmatrix}$$



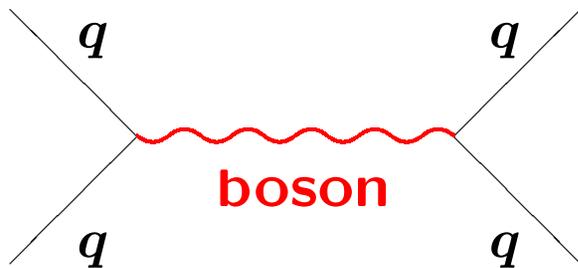
“Everything” around—and in—us is built of *up* and *down* quarks (plus gluons and electrons).

- What are the known interactions of the quarks?
- What additional interactions could the quarks feel?

Particles are manifestations of quantum fields.

Local quantum field theory:
interactions at a distance mediated by particle exchange.

One-boson exchange:



Quarks are fermions (spin 1/2).

1. Electromagnetic interaction

Coulomb force: $\sim \frac{1}{r^2}$ (*massless spin-1 particle exchange: photon*)

Laws of physics have a $U(1)$ gauge symmetry if the Lagrangian is invariant under a gauge transformation:

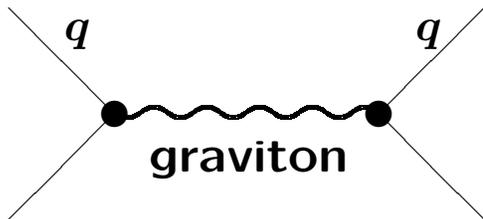
spin 1/2 (quark) field: $u(x^\mu) \rightarrow e^{2/3 i \alpha(x^\mu)} u(x^\mu)$

spin-1 (photon) field: $A_\mu(x) \rightarrow A_\mu(x) - \partial_\mu \alpha(x^\mu)$

Photon mass term, $A_\mu A^\mu$, is not gauge invariant
 $\Rightarrow U(1)_{\text{e.m.}}$ symmetry keeps the photon massless!

2. Gravitational interaction

Newton's law: $\sim \frac{1}{r^2}$ (massless spin-2 exchange: "graviton")



Classical limit is identical to General Relativity!

Graviton-quark interaction is a dimension-5 operator suppressed by the Planck scale

\Rightarrow theory is valid at energies $\ll 10^{18}$ GeV

3. Strong interaction

One-gluon-exchange force: $\frac{1}{r^2} \left[1 - b \ln \left(\frac{r}{10^{-14} \text{ cm}} \right) \right]^{-1}$

Approximation breaks down at distances $\gtrsim 10^{-14} \text{ cm}$,

bound states form: proton, neutron, pions, ...

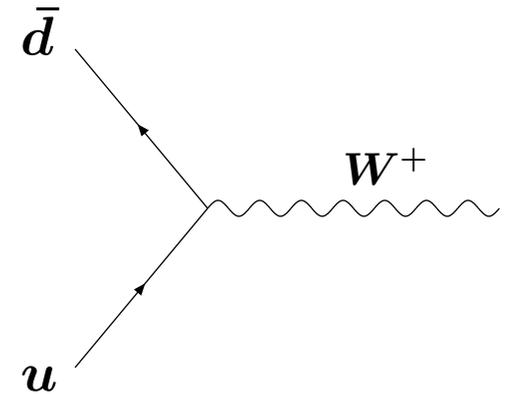
Short-range interaction due to confinement!

4. Weak interaction

“ Z^0 -exchange force”: $\sim \frac{1}{r^2} \exp \left[-2r / (10^{-16} \text{cm}) \right]$

Short-range interaction due to the Z^0 mass.

W^\pm interaction changes the quark flavor:



Weak & electromagnetic interactions require $SU(2)_W \times U(1)_Y$
gauge symmetry: 4 gauge bosons $\rightarrow W^\pm, Z^0, \gamma$

If the laws of physics are gauge invariant,

where are the W^\pm and Z^0 masses coming from?

Vacuum = the ground state of a quantum field theory.

**Lagrangian has an $SU(2)_W \times U(1)_Y$ gauge symmetry,
vacuum has only a $U(1)_{\text{e.m.}}$ gauge symmetry.**

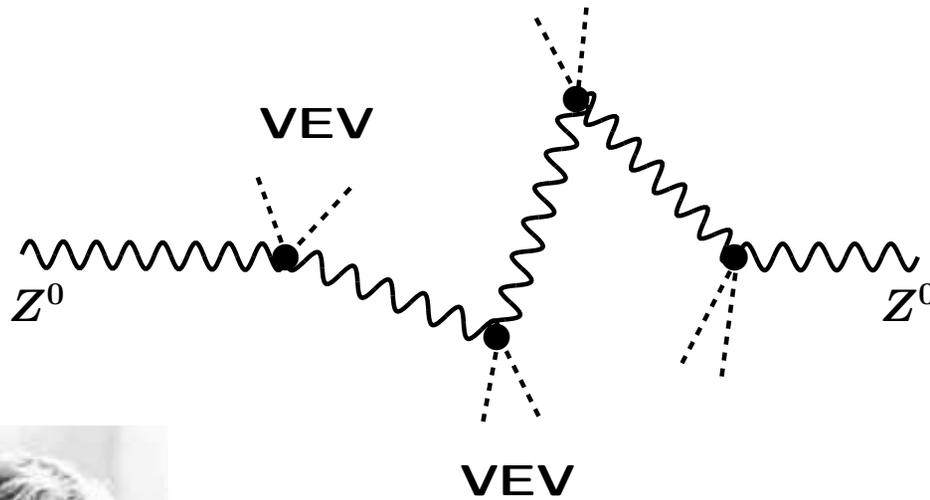
Similar to superconductivity: photon is massive inside superconductors.

**Need something to
populate the vacuum ...**



Higgs field 'condenses', has a Vacuum Expectation Value:

$$v_H = (\sqrt{2}/g)M_W \approx 174 \text{ GeV}$$



Z⁰ acquires a mass!

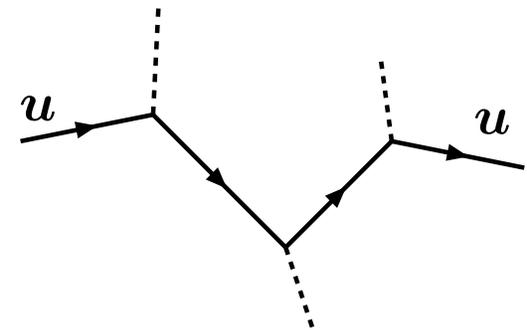


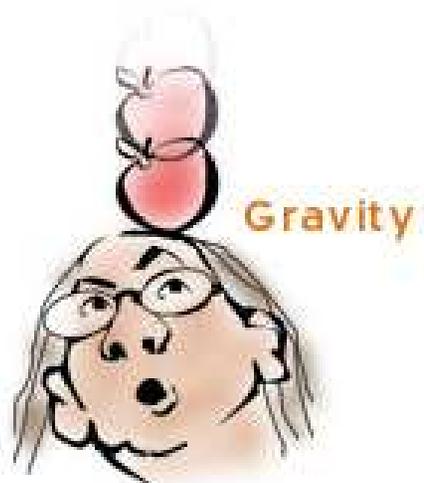
Steven Weinberg

"A Theory of Leptons"

1967

Quark masses also arise due to the Higgs field:





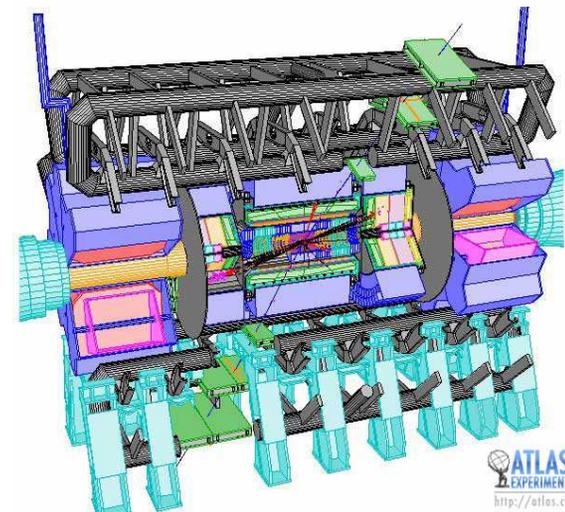
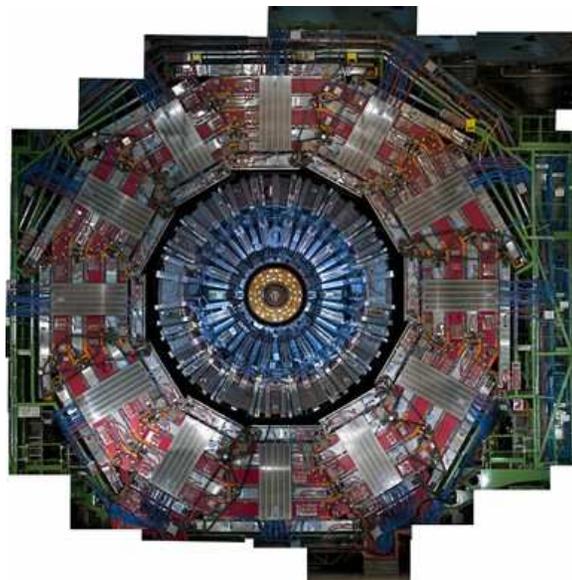
+ ?

Higgs field implies the existence of a particle carrying the same quantum numbers as the vacuum:

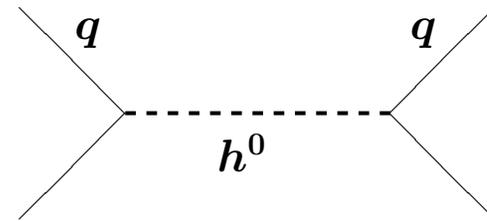
the Higgs boson (h^0) is a particle of spin 0 that couples to W, Z , quarks and leptons, proportional to their mass:

$$\frac{\sqrt{2}}{v_H} h^0 \left(M_W^2 W^\mu W_\mu + \frac{M_Z^2}{2} Z^\mu Z_\mu + \sum_q m_q \bar{q}_L q_R + \dots \right)$$

Discovered in 2012 with the CMS and ATLAS detectors:



Fifth “force”: Higgs exchange!



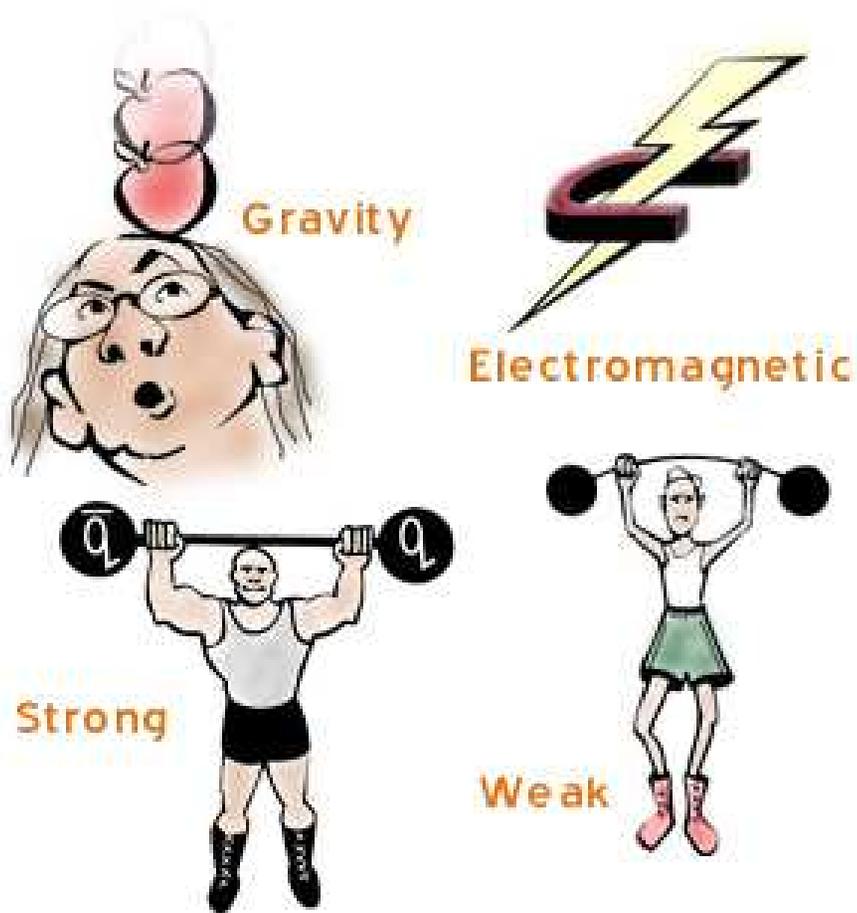
$$F \approx \left(\frac{m_q}{174 \text{ GeV}} \right)^2 \frac{1}{r^2} \exp \left[-2.5r / (10^{-16} \text{ cm}) \right]$$



tiny for ordinary stable matter... ($m_d \approx 5 \text{ MeV}$, $m_u \approx 3 \text{ MeV}$)

There is also a momentum-dependent interaction induced at two loops (involving a top quark and a gluon).

Scalar exchange \rightarrow spin-independent force.



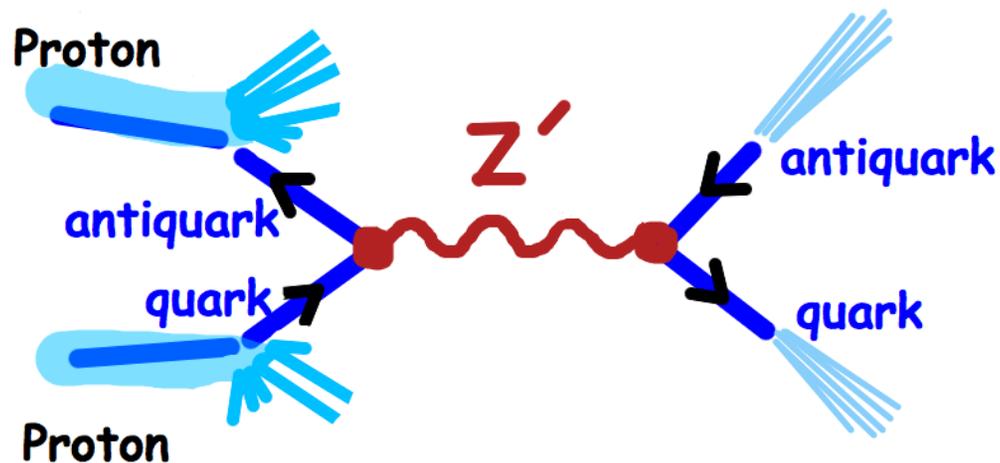
+ h^0 + ?

Could there exist additional quark interactions?

Yes, if the new bosons are sufficiently heavy, or sufficiently weakly coupled.

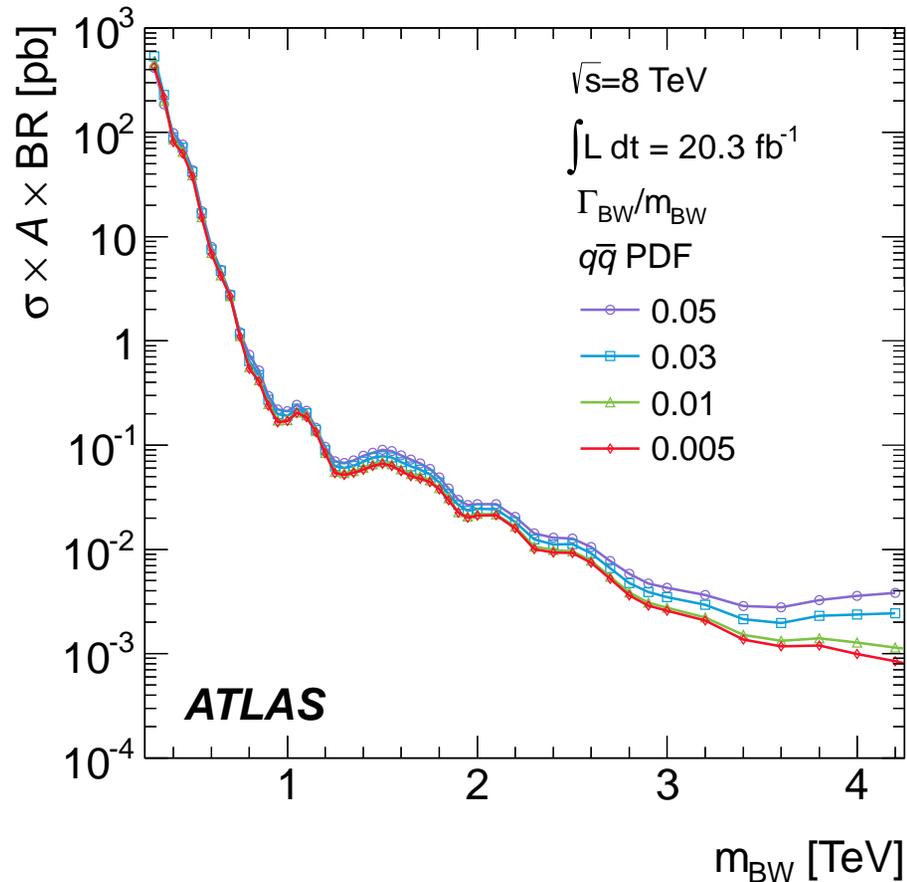
Hypothetical heavy particle of spin 1 and charge 0: Z' boson.

If Z' couples only to quarks (“leptophobic”), then it can be produced at hadron colliders (SpS, Tevatron, LHC) and decays back to quark-antiquark pairs, which give rise to jets of hadrons:



Two final-state jets, hard to separate from the background ...

However, the two hadronic jets form a resonance that can show up above the background if $M_{Z'}$ is large enough.

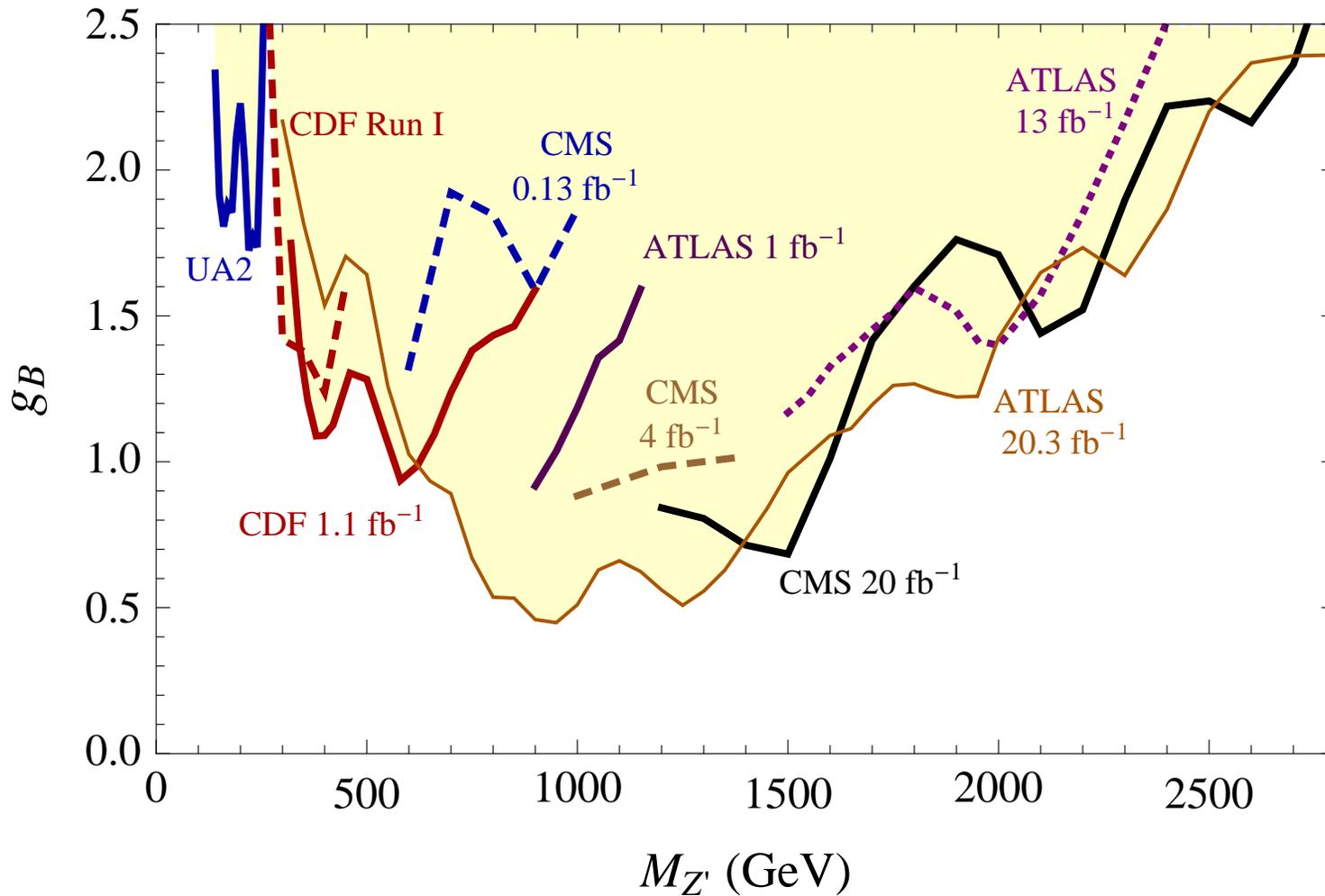


Limits on the production of a Gaussian resonance ...
(1407.1376)

Dijet resonances are not Gaussian (have a long tail due to final state radiation)

What are the limits on the Z' couplings to quarks?

“Baryonic” Z'_B : same coupling (g_B) to all six quark flavors.

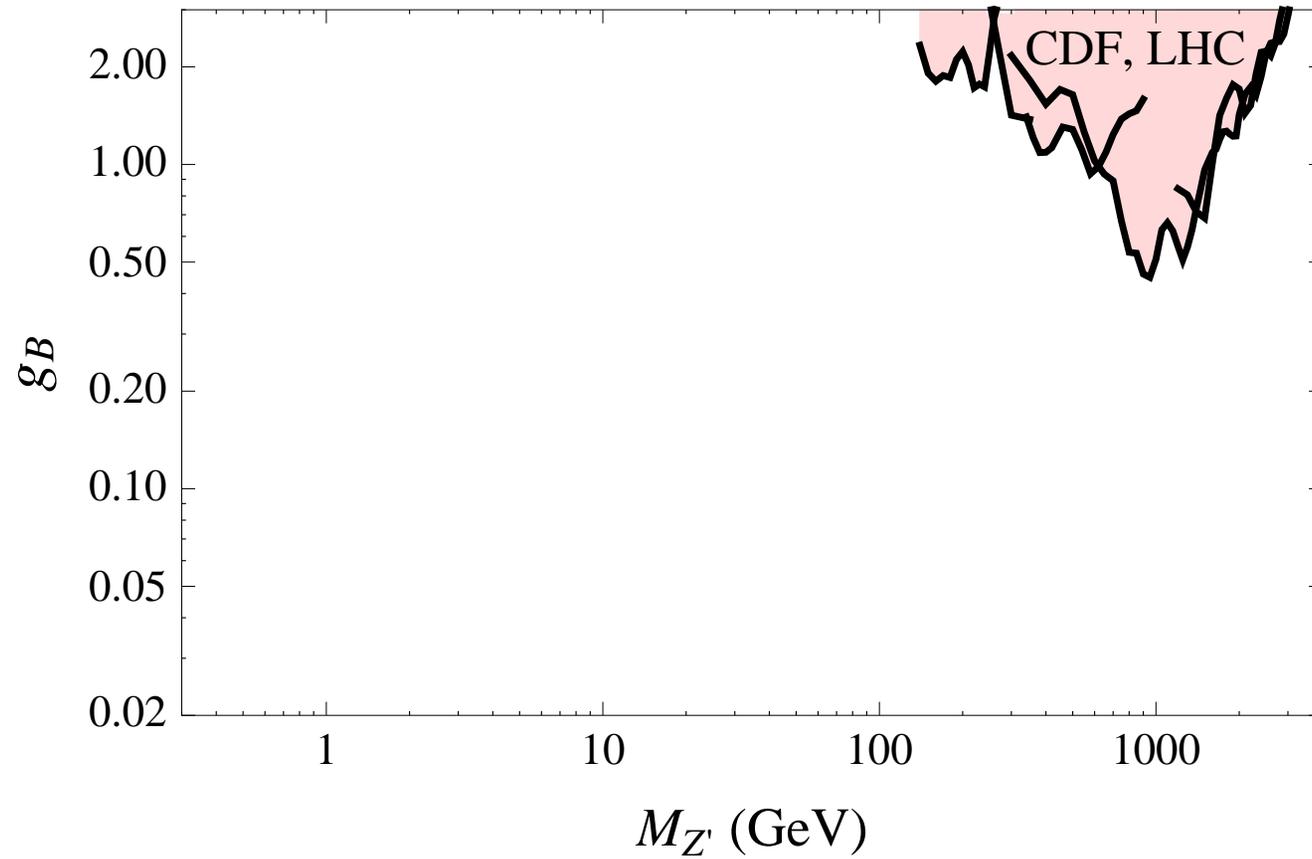


$$\mathcal{L}_q = \frac{g_B}{2} Z'_\mu \sum_q \left(\frac{1}{3} \bar{q}_L \gamma^\mu q_L + \frac{1}{3} \bar{q}_R \gamma^\mu q_R \right)$$

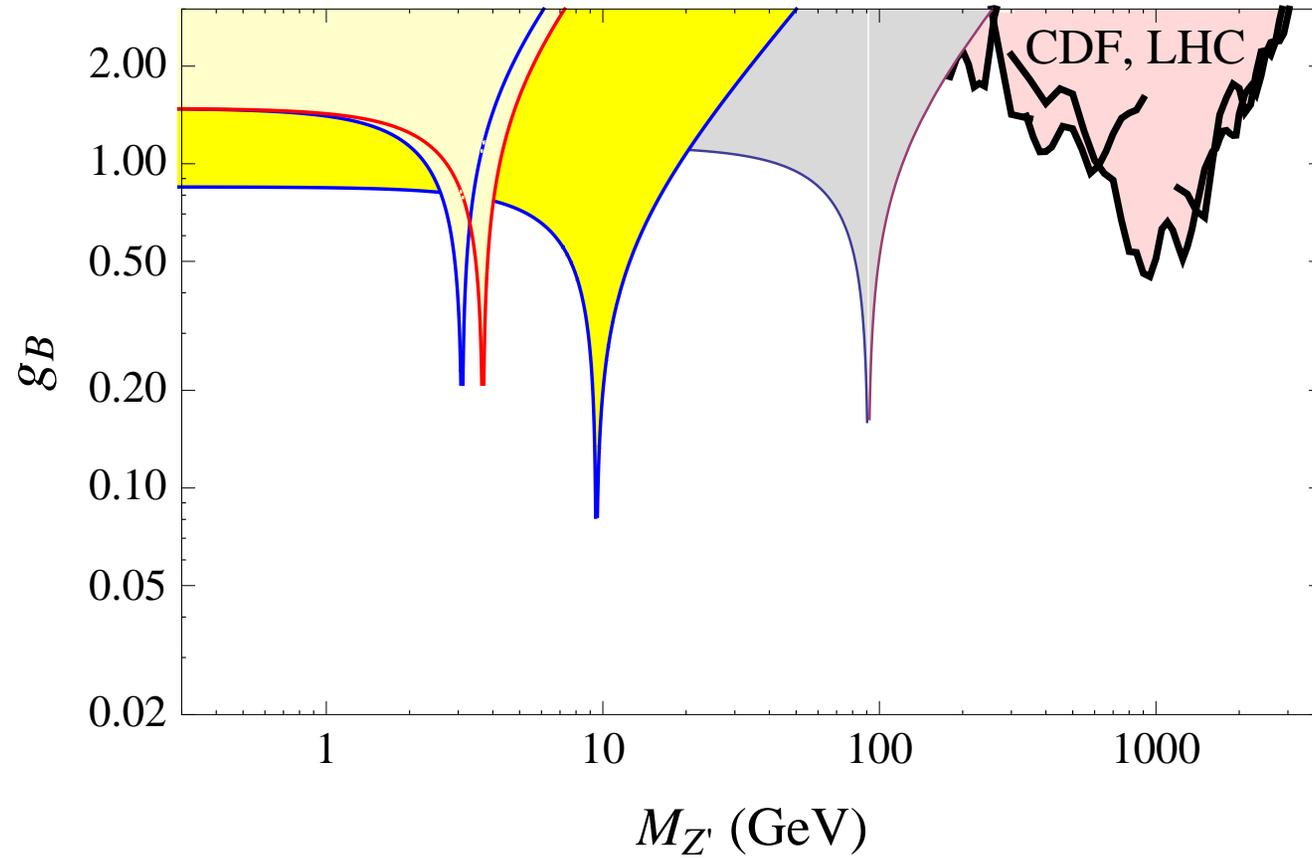
with Felix Yu:
1306.2629



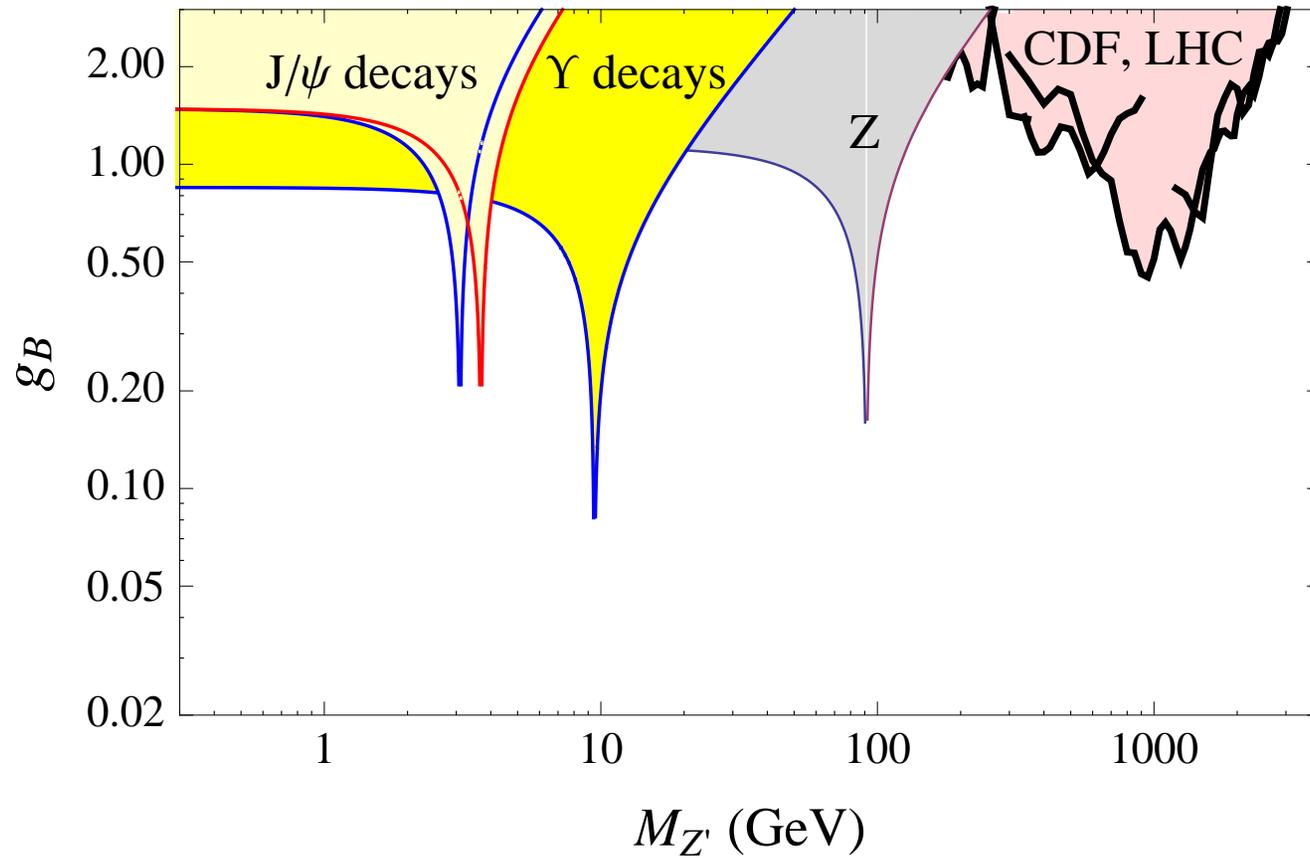
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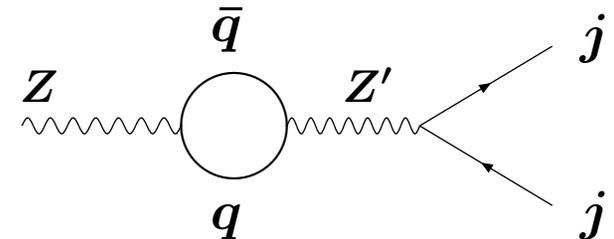
What are the limits on g_B for $M_{Z'} < 140$ GeV?



Υ are $b\bar{b}$ mesons, masses $\sim O(10)$ GeV.

ARGUS experiment, 1986: limit on $\Upsilon \rightarrow Z'^* \rightarrow jj$

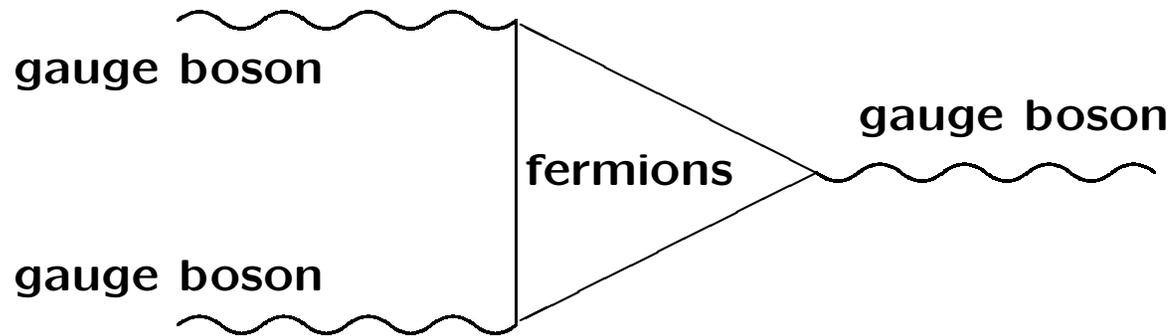
Small $Z - Z'$ mixing induced at one loop:



Gauge anomaly cancellation

Gauge symmetries may be broken by quantum effects.

Cure: sums over fermion triangle diagrams must vanish.



Standard Model – anomalies cancel within each fermion generation:

$$[SU(3)_c]^2 U(1)_Y: \quad 2(1/6) + (-2/3) + (1/3) = 0$$

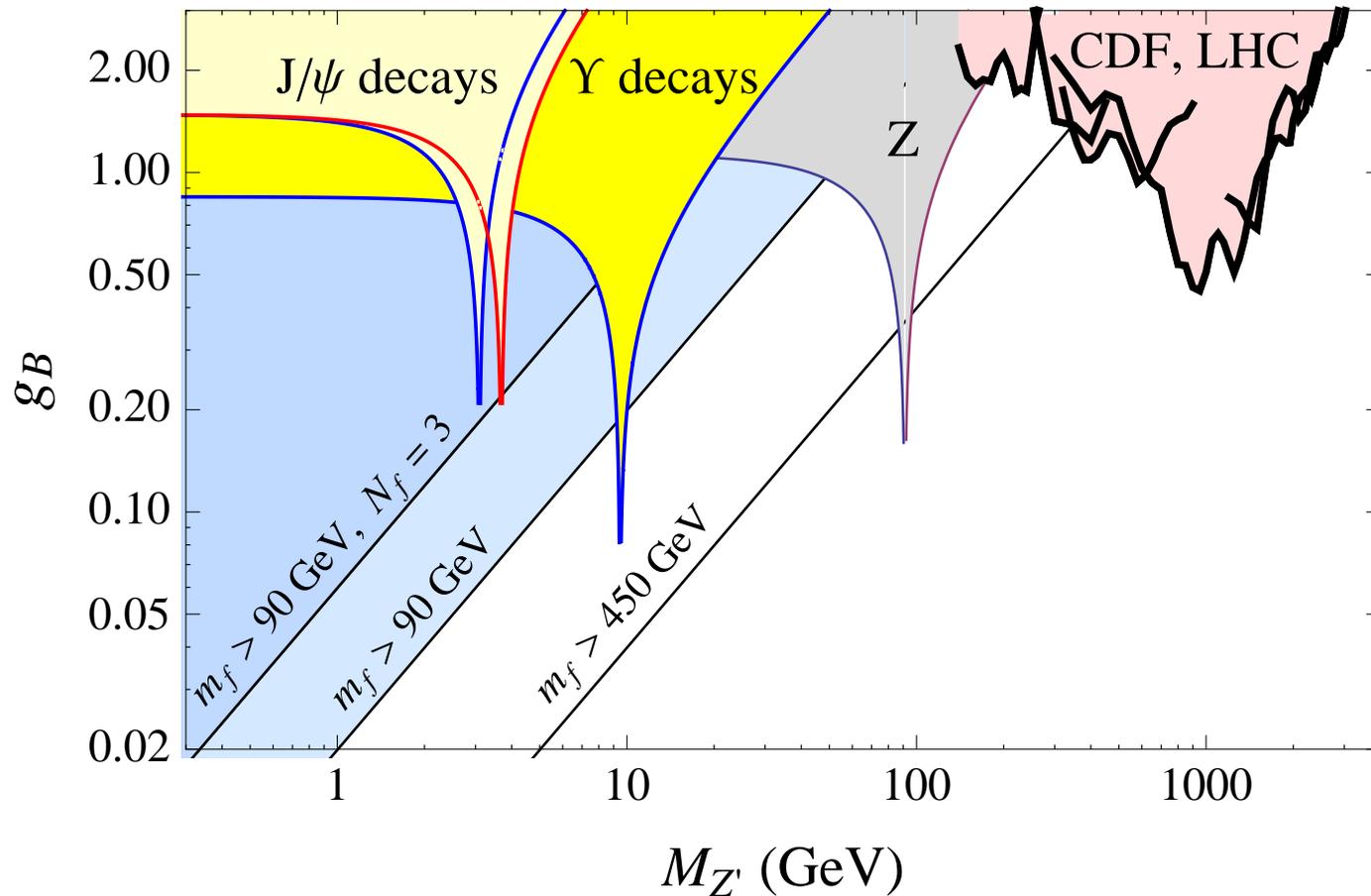
$$[SU(2)_W]^2 U(1)_Y: \quad 3(1/6) + (-1/2) = 0$$

$$[U(1)_Y]^3: \quad 3 \left[2(1/6)^3 + (-2/3)^3 + (1/3)^3 \right] + 2(-1/2)^3 + (-1)^3 = 0$$

$$\dots \quad (u_L, d_L) \quad u_R \quad d_R \quad (\nu_L, e_L) \quad e_R$$

Z' requires a new gauge symmetry: $SU(3)_c \times SU(2)_W \times U(1)_Y \times U(1)_B$

$U(1)_Y [U(1)_B]^2$ anomaly cancelation \Rightarrow charged vector-like fermions
 \Rightarrow colliders limits on fermion masses lead to a limit on g_B :



$U(1)_B$ is spontaneously broken
 by the VEV of a new scalar field φ .

with Claudia Frugiuele:
PRL 113, 061801 (2014)



A leptophobic Z' model with relaxed constraints:

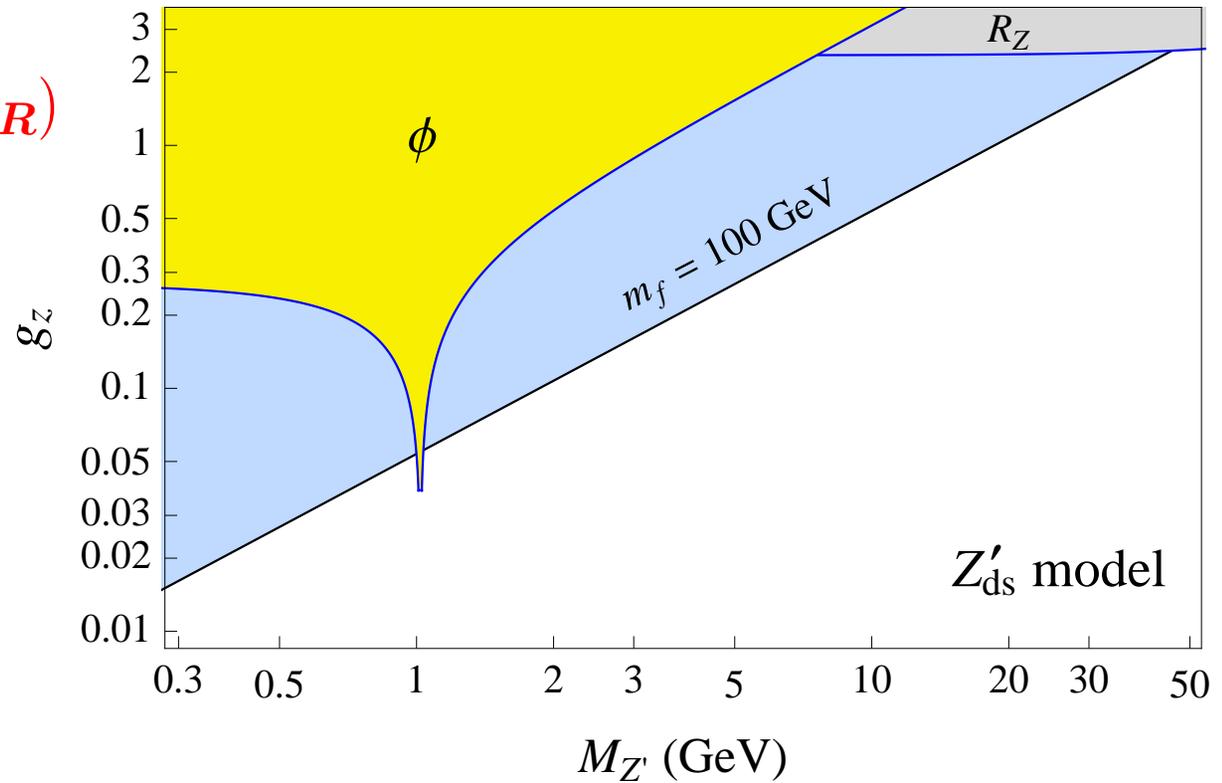
| field | $SU(3)_C$ | $SU(2)_W$ | $U(1)_Y$ | $U(1)_{ds}$ |
|-------------|-----------|-----------|----------|--------------|
| d_R | 3 | 1 | $+1/3$ | $+1/3$ |
| s_R | 3 | 1 | $-1/3$ | $-1/3$ |
| f_L, f'_L | 1 | 1 | $+1$ | $+2/3, -2/3$ |
| f_R, f'_R | 1 | 1 | $+1$ | $+1/3, -1/3$ |
| φ | 1 | 1 | 0 | 1 |

$$\mathcal{L}_q = \frac{g_z}{2} Z'_\mu (\bar{d}_R \gamma^\mu d_R - \bar{s}_R \gamma^\mu s_R)$$

no coupling to q_L

→ no flavor-changing Z'
couplings.

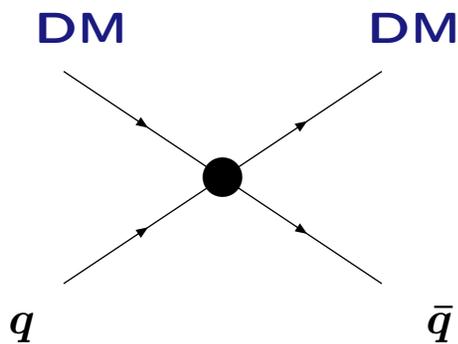
s, d masses from dimension-5
operators.



Dark matter requires particle(s) beyond the SM

DM particle may be part of a large hidden sector.

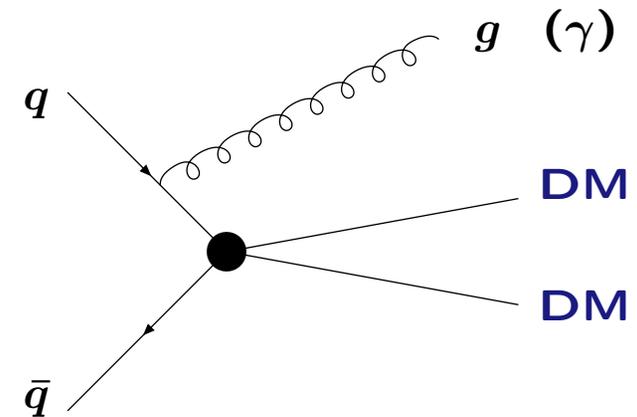
There can be several particles contributing to the DM density.



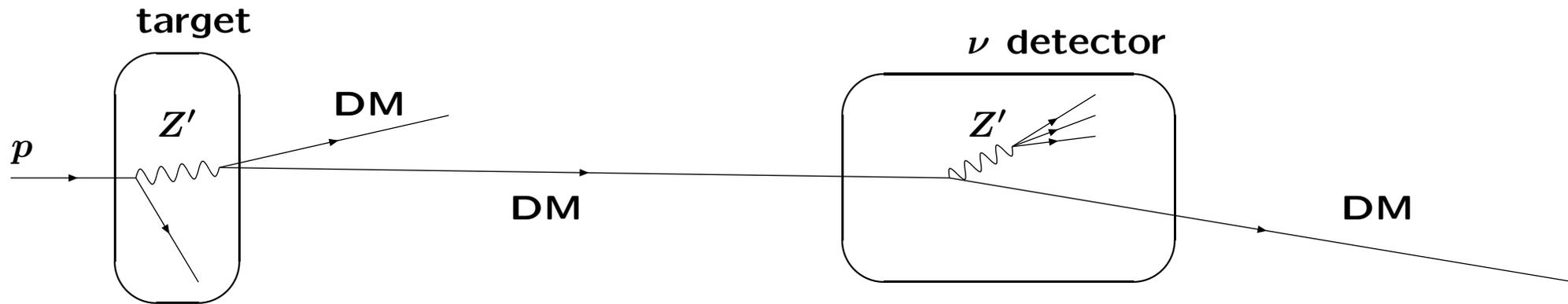
DM direct detection relies on DM interaction with nuclei (*i.e.*, quarks or gluons), and $m_{\text{DM}} \gtrsim 5 \text{ GeV}$

If DM interacts with quarks, it may be produced at the LHC. (*Bai, Fox, Harnik, Tait, ...*)

Current limit: $g_B \lesssim 0.2$



Quark – DM interaction: fixed target experiments → DM beam



Neutral-current events → background: ν beam!

DM events have higher energy, can separate the signal from background.

Main Injector: $> 10^{21}$ protons (of 120 GeV) on target.

Presence of DM particles can be probed with the NO ν A, MINOS, LBNF, ..., detectors.

Similar search for lighter Z' at MiniBoone – B. Batell et al, 1405.7049

Even when quarks are neutral under the new gauge group, they can interact with Z' :

$$\frac{1}{M^2} Z'_{\mu\nu} (C_u \bar{Q}_L \sigma^{\mu\nu} u_R H + \text{h.c.})$$

$$Z'_{\mu\nu} = \partial_\mu Z'_\nu - \partial_\nu Z'_\mu \quad (\text{field strength})$$

C_u : dimensionless parameters

→ magnetic- and electric-like dipole moments

Interaction strength increases with the energy of the production process!

Conclusions

- 5 known interactions of quarks:
 - electromagnetic; gravitational
 - strong; weak; Higgs
- Additional interactions may exist.
Leptophobic Z' — mass:
 - $\gtrsim 100$ GeV: probed at colliders
 - $\gtrsim 10$ GeV: probed in meson decays
 - $\gtrsim 50$ GeV: “theoretical” limits
- Many interesting possibilities:
 - interactions with DM
(use ν detectors ...)
 - momentum-dependent couplings
 - ...

