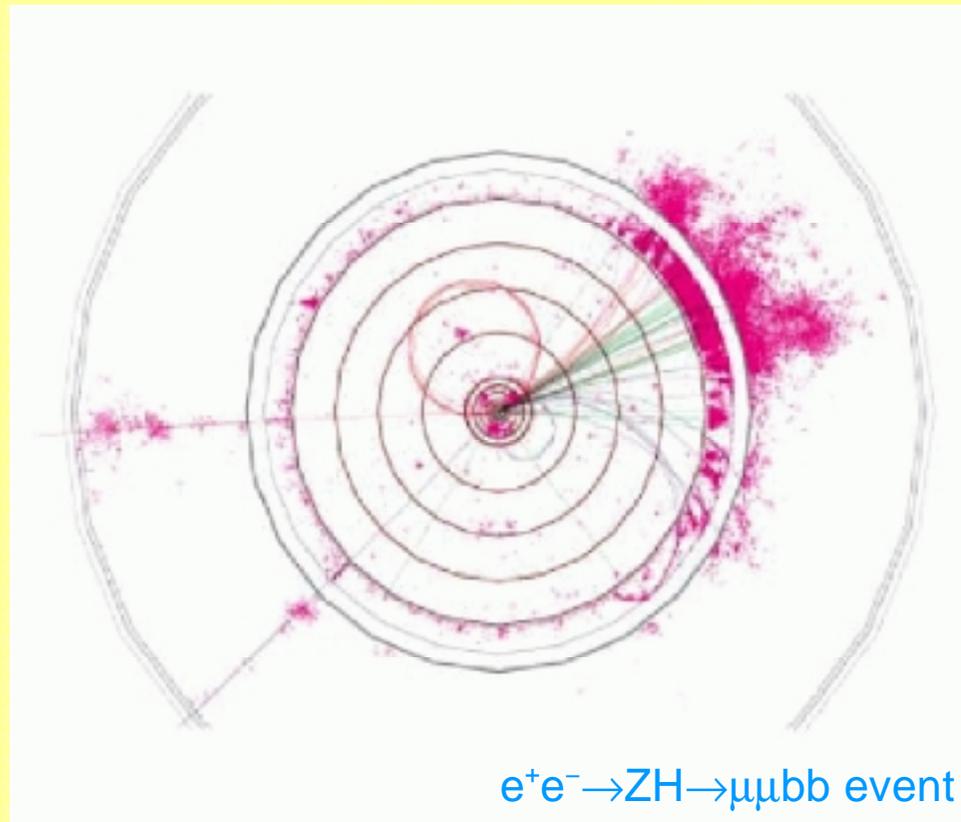
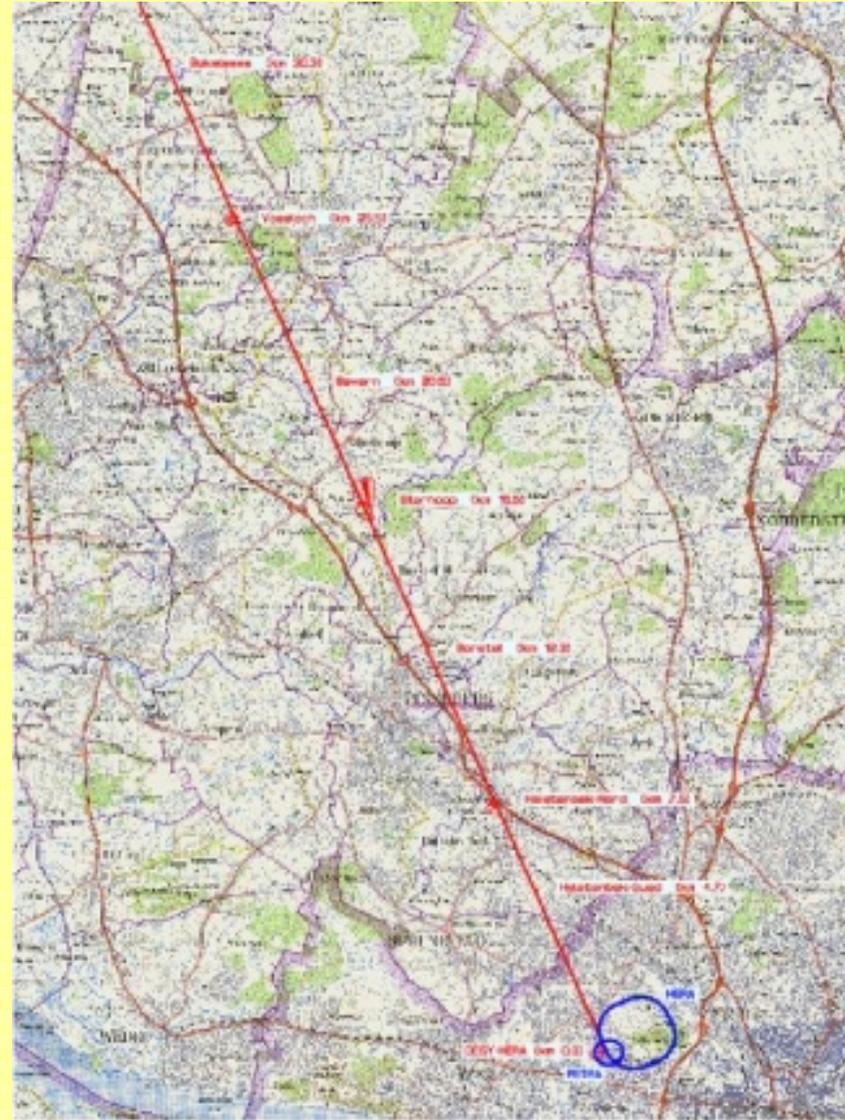
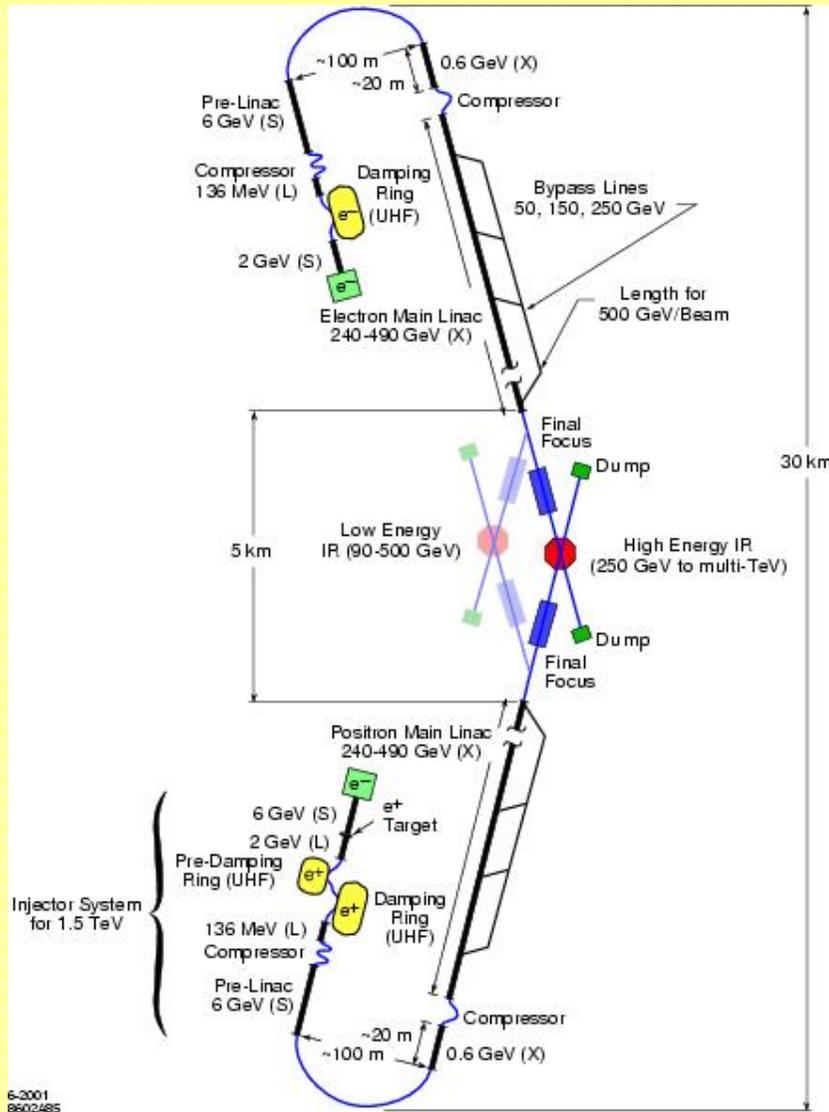


Detectors at the Linear Collider



David Gerdes
University of Michigan
March 26, 2002

Linear Collider Overview

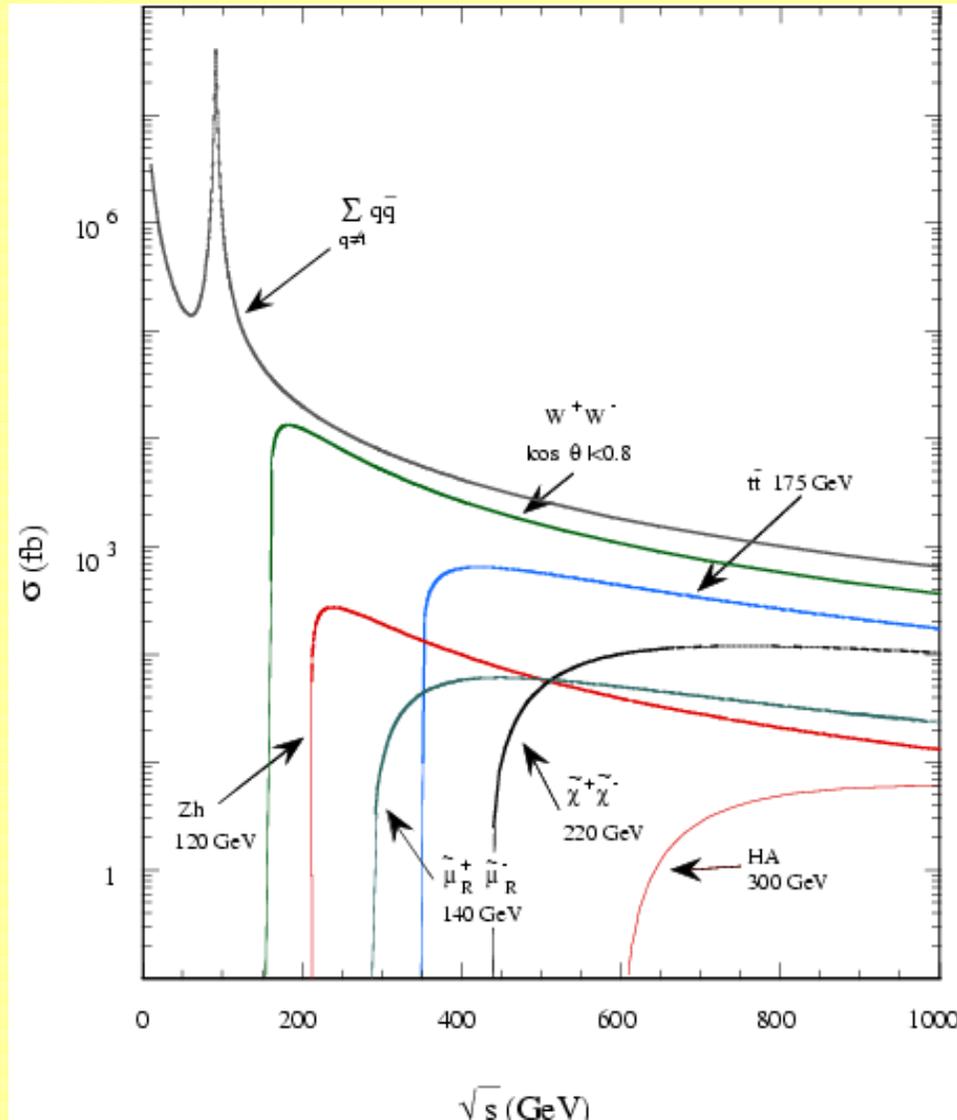


Machine Parameters

	TESLA(500)	TESLA(800)	NLC(500)	NLC(1000)	Tevatron
E (GeV)	500	800	500	1000	2000
Lum. x 1E33	31	5	20	34	0.1
Rep rate (Hz)	5	3	120	120	--
Bunches/pulse	2820	4500	190	190	--
Bunch sep (ns)	337	189	1.4	1.4	396
$\sigma(x)$ at i.p.	553 nm	391 nm	245 nm	190 nm	30 μ m
$\sigma(y)$ at i.p.	5 nm	2 nm	2.7 nm	2.1	30 μ m
$\sigma(z)$ at i.p.	0.4 mm	0.3 mm	110 nm	110 nm	30 cm
δB(%)	3.3	4.7	4.7	10.2	0
P(e⁻) (%)	80-90	80-90	80-90	80-90	--
P(e⁺)(%)	60	60	--	--	--

NB: Total cross section:
 ~5 pb (e⁺e⁻, 500 GeV)
 ~50 mb (p-pbar, 2 TeV)

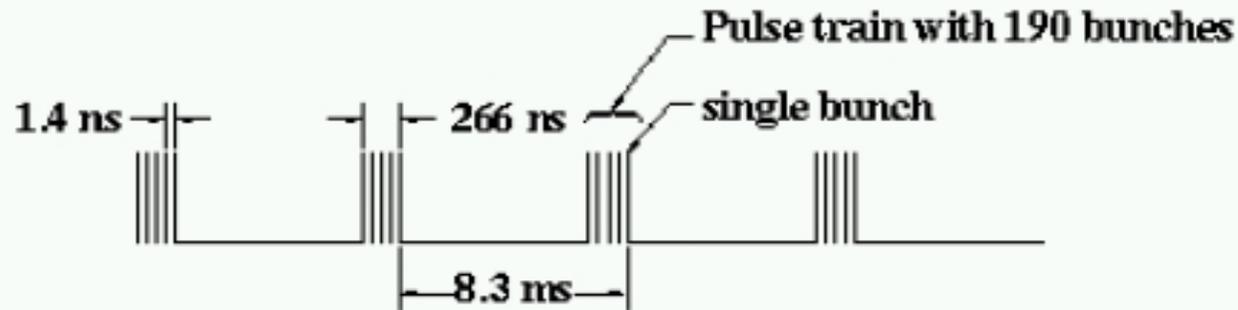
e^+e^- Cross Sections



Objects with electroweak couplings are produced "democratically"

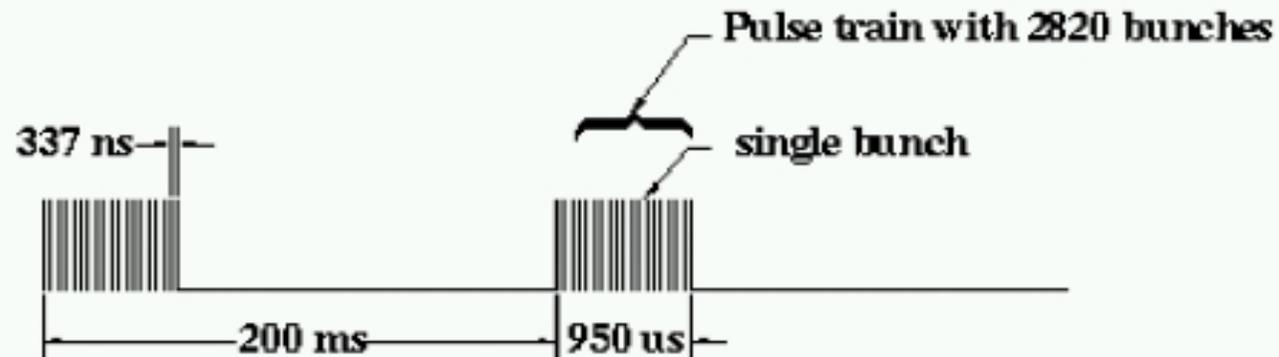
Bunch structure

NLC



a. NLC/JLC 120 pulse trains/sec

TESLA



b. TESLA 5 pulse trains/sec

Low rates \Rightarrow Read out between crossings
and "trigger" in software

The Challenge of LC detectors?

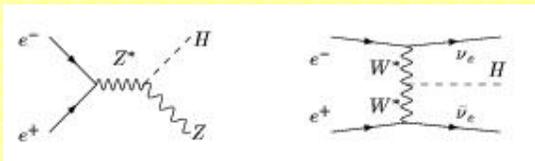
- Typical Tevatron experimentalist's view:
 - "It's trivial!"
 - "Just copy a LEP detector or SLD!"
- Sure, the issues of trigger/DAQ appear straightforward.
- But the challenge is to build a high-precision instrument capable of producing the "textbook data on the next energy scale" (M. Peskin)

Physics Unique to the LC

- Measurement of Higgs branching ratios, e.g. via $e^+e^- \rightarrow Zh$ followed by $Z \rightarrow ll$, $h \rightarrow X$
- Study of the top quark threshold, measurements of m_t , Γ_t , α_s , g_{tth}
- Precision SUSY spectroscopy
- ttH production; top–Higgs Yukawa coupling
- Trilinear Higgs coupling; reconstruction of the Higgs potential.

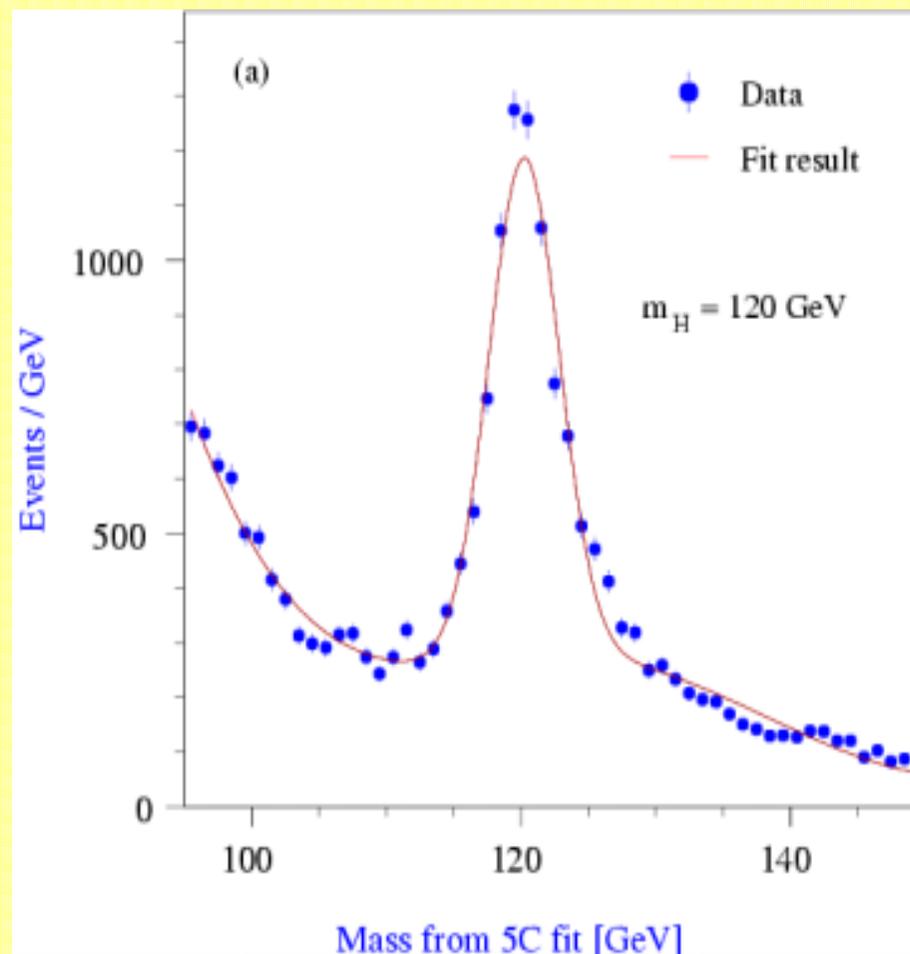
Higgs Reconstruction

- Current EWK fits favor a light Higgs, $m_h < 200$ GeV
- Produced at the LC in recoil off the Z



Detector requirement:

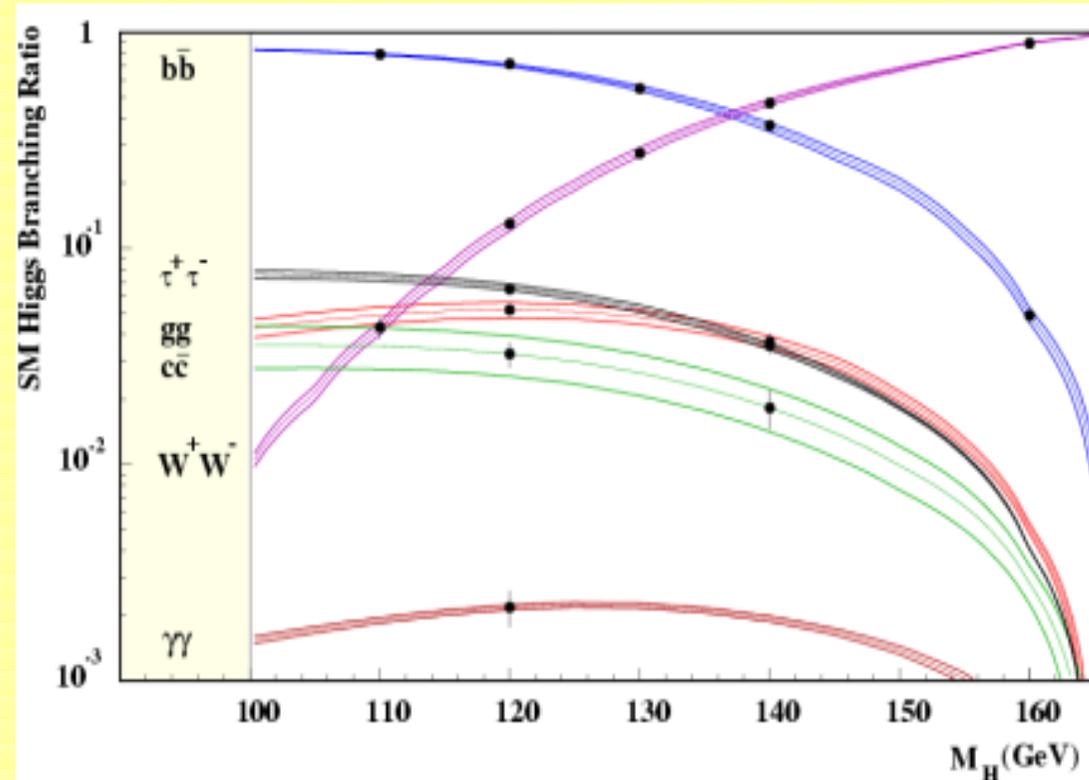
Excellent tracking resolution,
 $\sigma(1/p_T) \sim 5 \times 10^{-5} \text{ GeV}^{-1}$



Measurement of Higgs BR's

- Event selection is based on the Z, not the H, so the Higgs decays constitute an unbiased, inclusive sample.
- Key to determining that "the Higgs is the Higgs"—or something else!

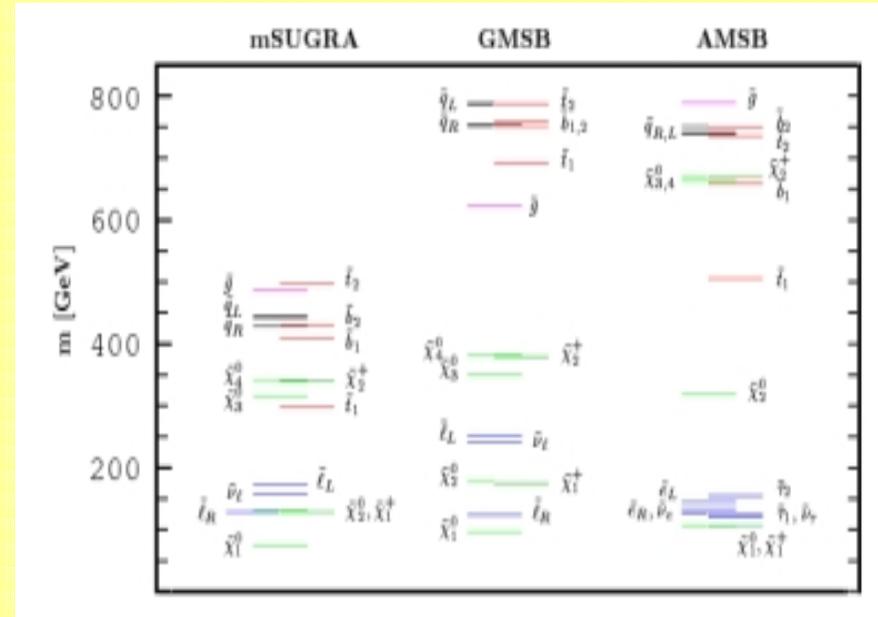
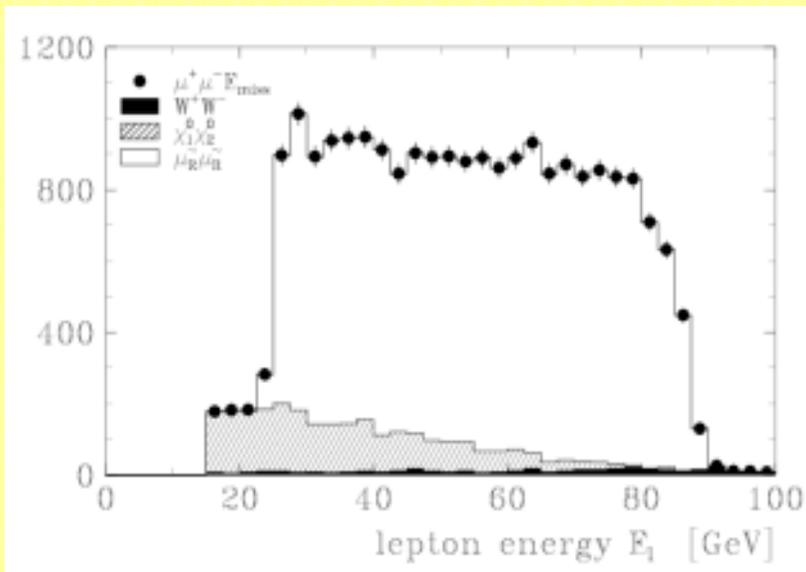
- Detector requirements: precision vertexing, flavor tagging.



M. Battaglia

SUSY Spectroscopy

- Pattern of sparticle masses encodes info about SUSY-breaking.
- Example: endpoints of muon energy spectrum in $e^+e^- \rightarrow \mu_L \mu_L \rightarrow \mu^+ \mu^- \chi_L^0 \chi_L^0$

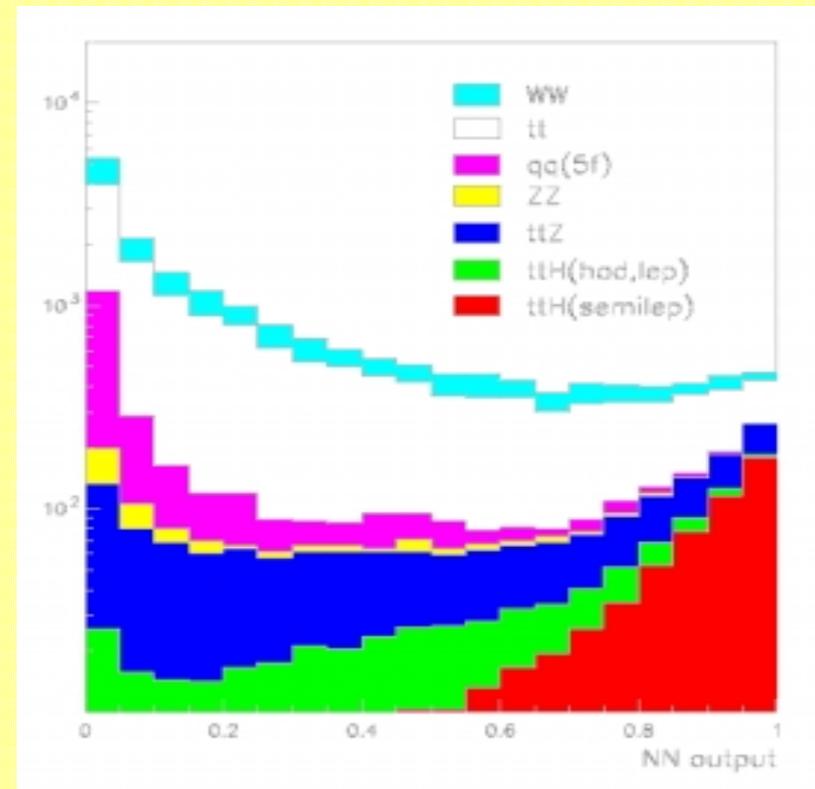


- Detector requirements:
- Excellent resolution
 - Hermeticity
- Accelerator requirements:
- Ability to vary energy over wide dynamic range
 - Adjustable beam polarization

ttH production and the Top Yukawa Coupling

- $e^+e^- \rightarrow ttH \rightarrow WbWb bb$
- Very complicated final state:
 - Up to 8 jets
 - 4 b's
 - Many kinematic constraints
- Tiny cross section (~ 2 fb), with backgrounds ~ 3 orders of magnitude higher.

- Detector requirements:
 - Excellent jet resolution
 - Flavor tagging



Juste, Merino: hep-ph/9910301

Note: ttH cross section maximized at $\sqrt{s} \sim 800$ GeV

Reconstructing the Higgs Potential

Physical basis of EWSB: nonzero minimum of the Higgs potential,

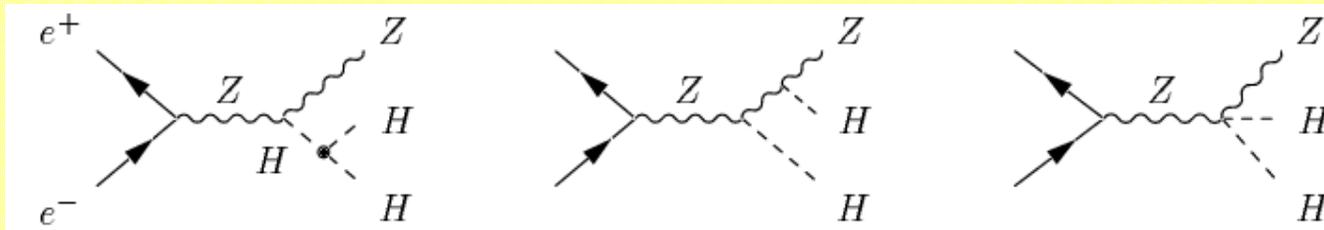
$$V = \lambda(|\phi|^2 - \frac{1}{2}v^2)^2$$
$$\rightarrow \lambda v^2 H^2 + \lambda v H^3 + \frac{1}{4}\lambda H^4,$$

where $v=246$ GeV.

- Measure coefficient of the quadratic term by measuring the mass,

$$M_h^2 = 2\lambda v^2$$

- Measure the trilinear self-coupling via Higgs pair production, $e^+e^- \rightarrow ZHH$

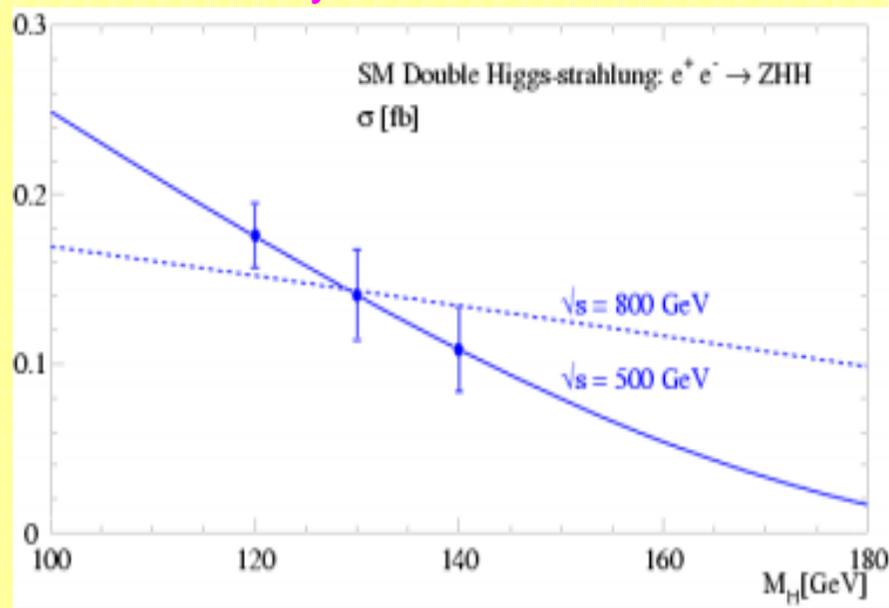


Reconstructing the Higgs Potential

Striking final state: 4 b jets, plus dileptons, jets, or missing energy from the Z decay.

But: Tiny cross section (~ 0.2 fb), huge 4-, 6-fermion backgrounds make this among the most challenging processes to study at the LC.

Sensitivity with 1 ab^{-1} :



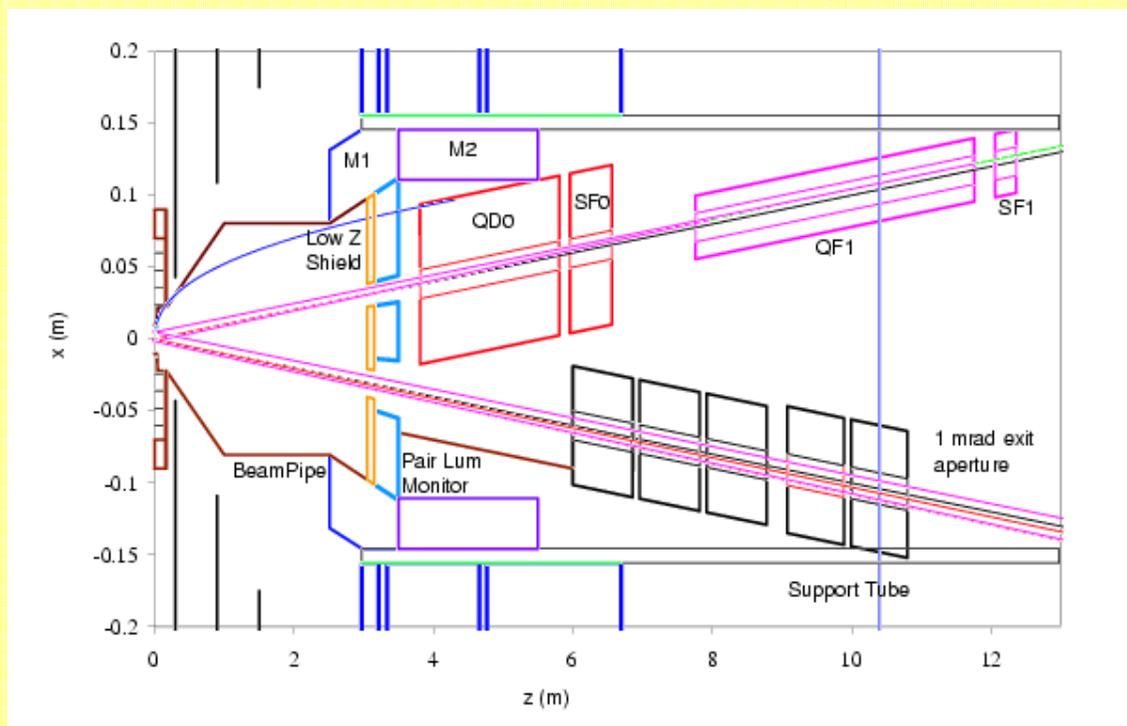
Detector requirements:

Excellent jet resolution
Flavor tagging

Accelerator requirements:

Lots and lots of luminosity

The Interaction Region: Where the Rubber Meets the Road

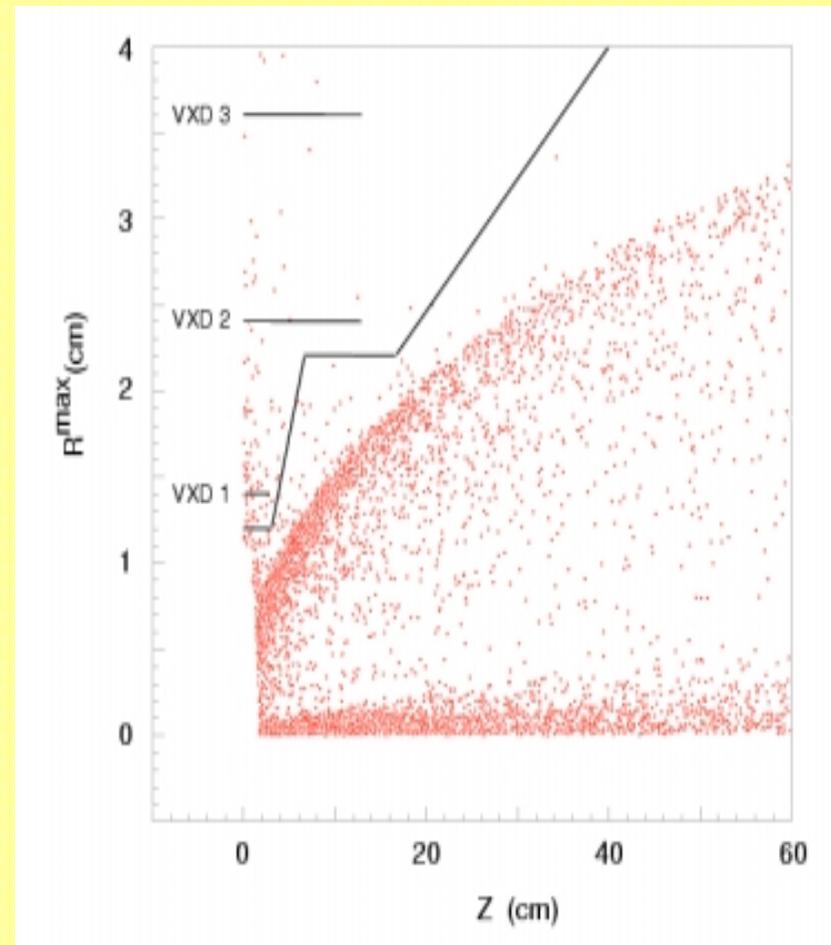


Beams collide at a 20 mrad crossing angle to avoid parasitic collisions.

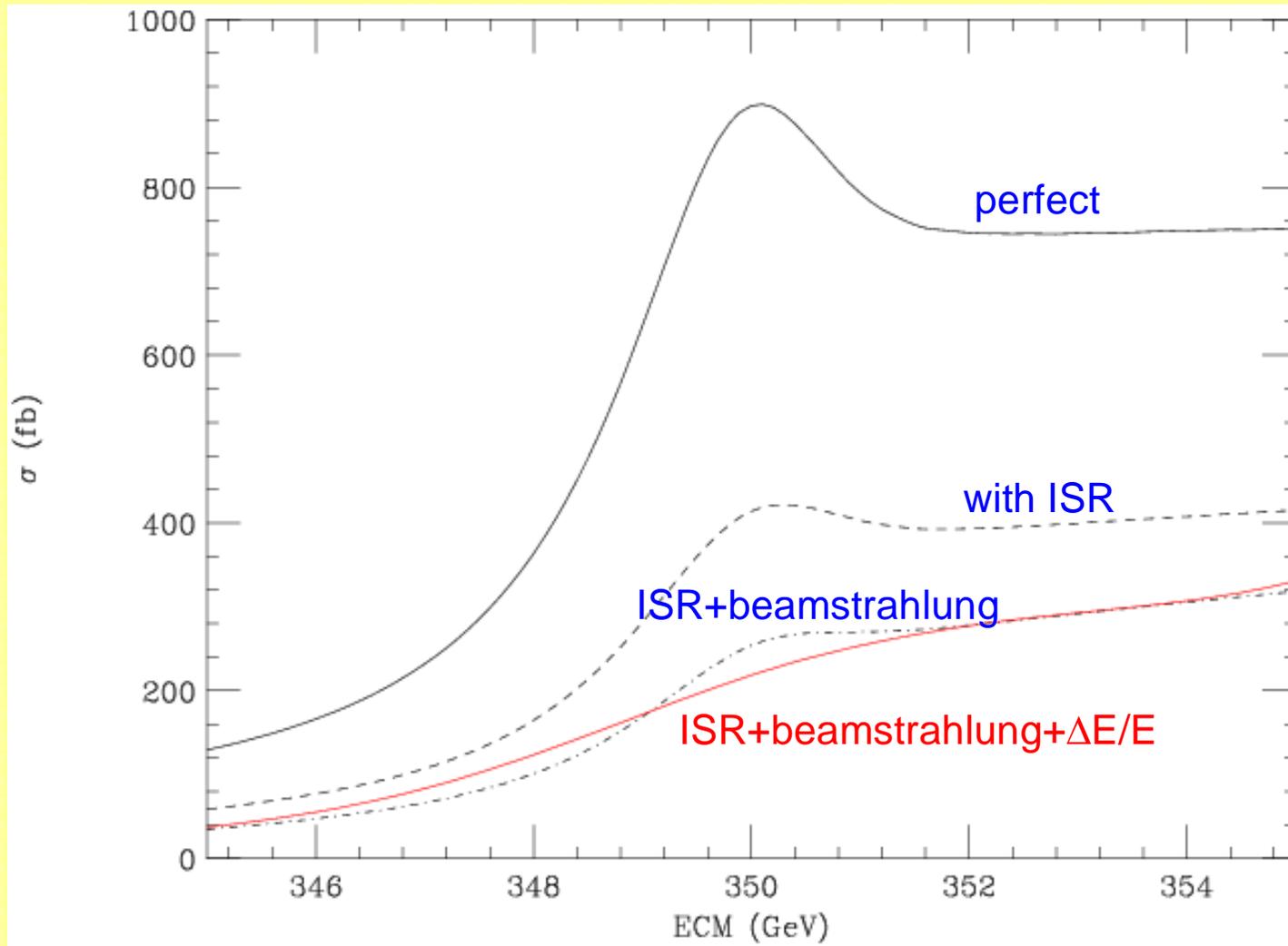
Leads to effects from the solenoidal field of the detector \rightarrow compensate by offsetting the final quadrupole (QD0)

Beam-Beam Interactions

- **Beamstrahlung:** radiation of photons from one beam due to the field of the other beam
 - Leads to energy spread
- **Production of e^+e^- pairs**
 - About 10^5 per bunch, mean energy 13 GeV (~few W of power)
 - Most go into "dead cone" shown at right.
 - Drives B-field magnitude
- **Other backgrounds:** hadrons from $\gamma\gamma$ interactions; neutrons, muons.
- **R&D question:** time-stamp calorimeter and track info?



Machine Effects on Top Threshold Lineshape



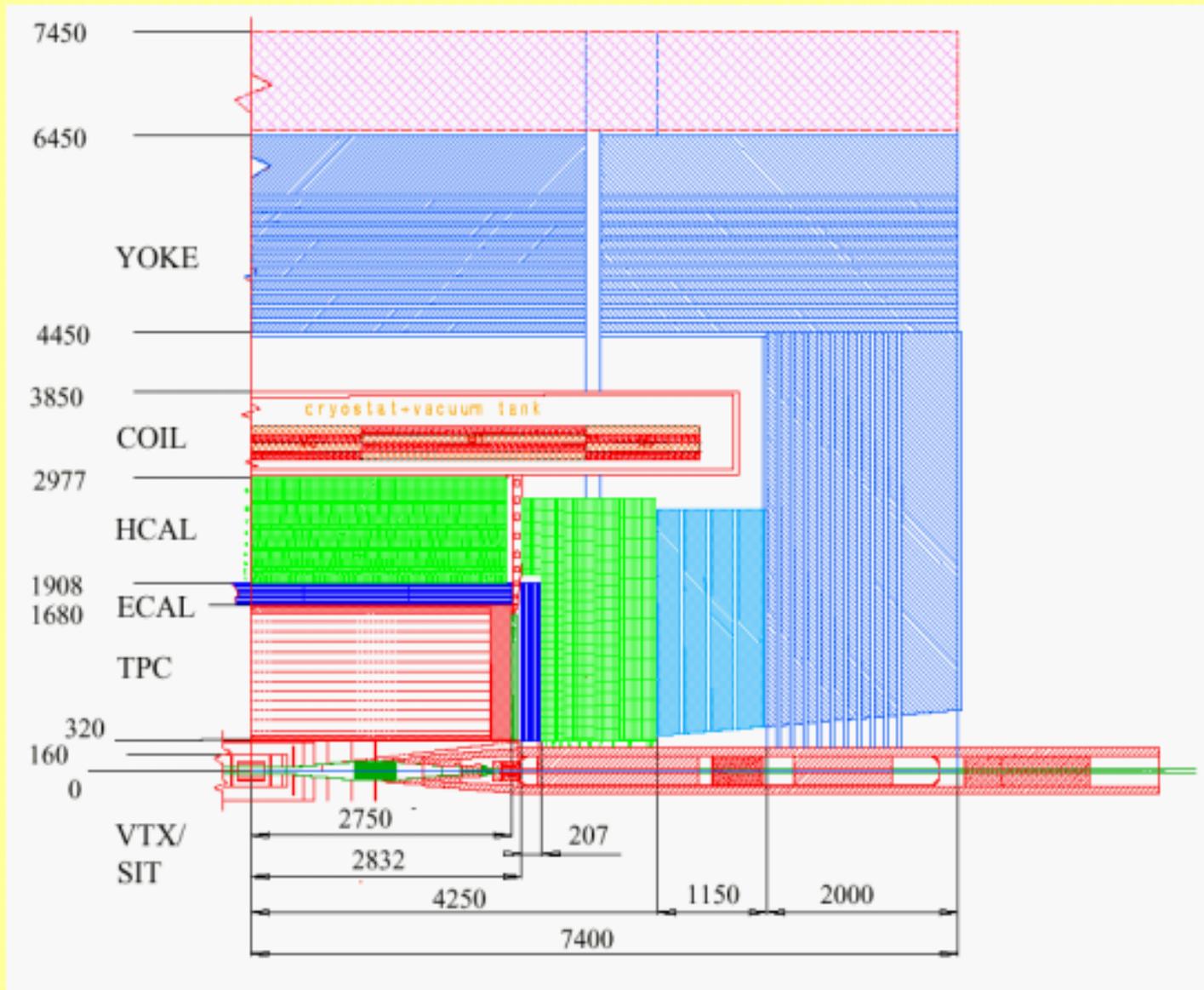
So to sum up the detector challenges, we need:

- Excellent hermeticity
 - Excellent calorimetry
 - Excellent tracking resolution
 - Excellent flavor-tagging
- multijet final states,
forward-peaked cross sections
W/Z separation
Precision mass measurements
b/c/ τ identification

• As opposed to a Tevatron or LHC detector where major design drivers include:

- Ability to handle very high rates
- Radiation tolerance

TESLA Detector



LC-LHC Detector Comparisons

	CMS	ATLAS	TESLA design
Tracker thickness (X0)	30%	28%	5%
Vertex layer thickness	1.7%	1.7%	0.06%
Vertex channel count (Mpixels)	39	100	800
ECAL granularity (det. elements)	76,000	120,000	32,000,000

- 6x less material for photon conversions
- Vertex detector 3–6x closer to interaction point
- 35x smaller vertex detector pixel size
- 30x thinner vertex detector layers
- >200x higher ECAL granularity (*expensive*)

"A Ferarri, as opposed to a state-of-the art Abrams battle tank." (C. Damerell)

North American "Large" Detector

