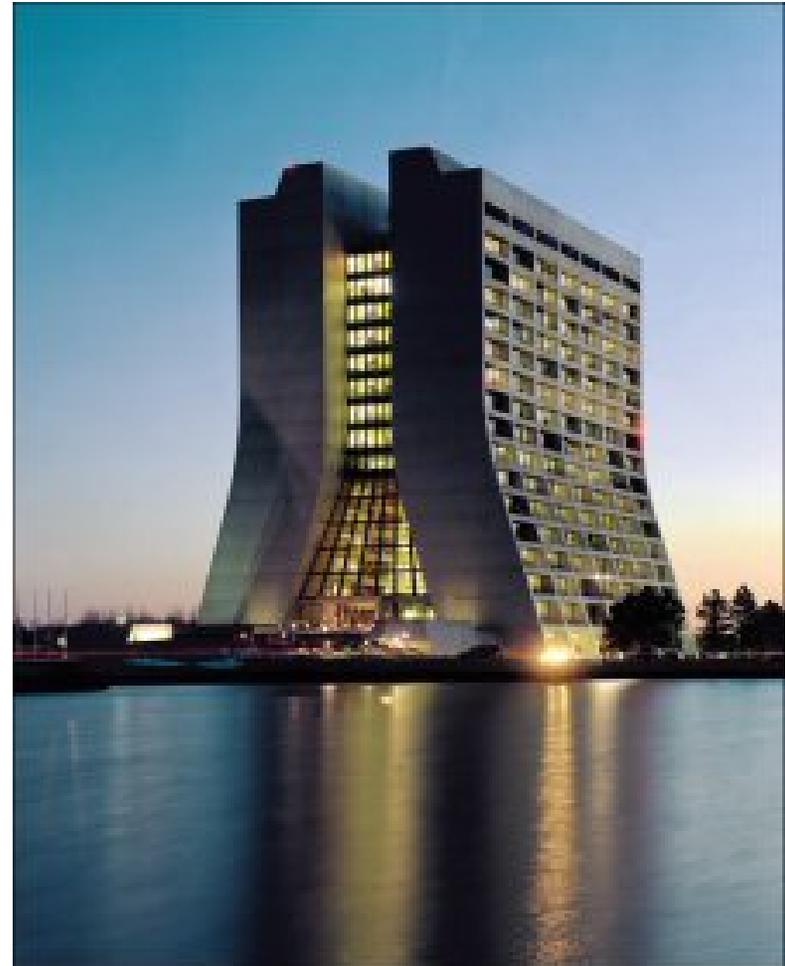


Tests of the Electroweak Theory

- History/introduction
- Weak charged current
- QED
- Weak neutral current
- Precision tests
- Rare processes
- CP violation and B decays
- Neutrino mass



Announcements

- e-mail questions to carena@fnal.gov or pgl@fnal.gov (between lectures)
- Slides and video at <http://theory.fnal.gov/AcademicLectures.html>
- Today and 12/6 in Curia II; others in One West

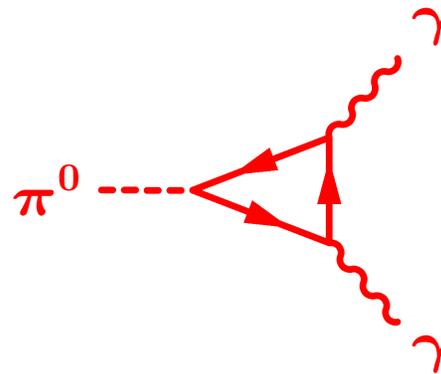
π and K Decays, and the Strong Interactions

- Charged pion decay

$\pi^+ \rightarrow \mu^+ \nu_\mu$	$(\pi_{\mu 2})$	99.99%	
$\rightarrow e^+ \nu_e$	$(\pi_{e 2})$	1.23×10^{-4}	
$\rightarrow e^+ \nu_e \pi^0$	$(\pi_{e 3})$	1.025×10^{-8}	π beta decay

- $\pi^0 \rightarrow 2\gamma$: 98.8%

- Electromagnetic
- Color counting via global anomaly



- **Space reflection (parity): violated maximally by WCC, conserved by strong interactions, since $J_\mu = V_\mu - A_\mu$ (also C ; $CP \sim$ conserved)**

$$\begin{aligned}
 P|\pi(\vec{q})\rangle &= -|\pi(-\vec{q})\rangle \\
 PV_\mu P^{-1} &= V^\mu \quad (P\vec{V}P^{-1} = -\vec{V}, \quad PV_0P^{-1} = V_0) \\
 PA_\mu P^{-1} &= -A^\mu \quad (P\vec{A}P^{-1} = \vec{A}, \quad PA_0P^{-1} = -A_0)
 \end{aligned}$$

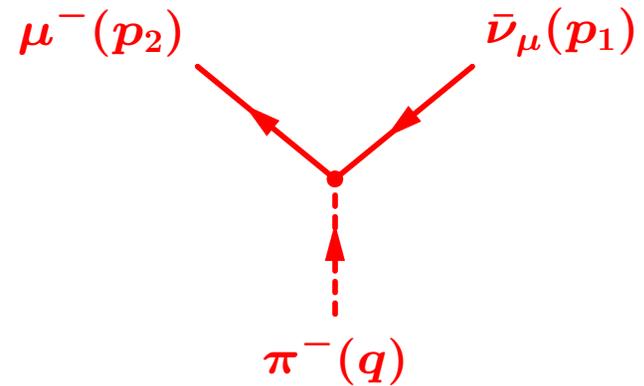
- **Semi-leptonic processes for light hadrons/leptons**

$$\begin{aligned}
 H &= \frac{G_F}{\sqrt{2}} \left[J^{\ell\mu} J_\mu^{h\dagger} + J^{\ell\mu\dagger} J_\mu^h \right] \\
 &= \frac{G_F}{\sqrt{2}} \left[\bar{\mu}\gamma^\mu(1 - \gamma^5)\nu_\mu + \bar{e}\gamma^\mu(1 - \gamma^5)\nu_e \right] \bar{u}\gamma_\mu(1 - \gamma^5)d' + HC \\
 d' &\equiv \sum_{i=d,s,b} V_{ui}d_i \simeq d \cos\theta_c + s \sin\theta_c
 \end{aligned}$$

($\theta_c =$ Cabibbo angle; $\sin\theta_c \simeq 0.22$)

- For $\pi^- \rightarrow \mu^- \bar{\nu}_\mu$

$$\begin{aligned}
 J_\mu^{h\dagger} &= \bar{u} \gamma_\mu (1 - \gamma^5) d \cos \theta_c \\
 &\equiv V_\mu^\dagger - A_\mu^\dagger
 \end{aligned}$$



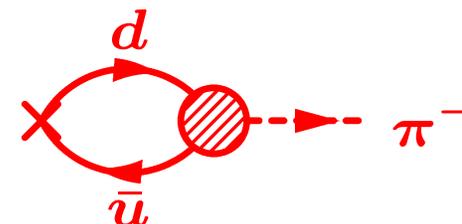
$$\begin{aligned}
 M &\equiv -i \langle \mu^-(p_2) \bar{\nu}_\mu(p_1) | H | \pi^-(q) \rangle \\
 &= -i \frac{G_F}{\sqrt{2}} \bar{u}_\mu \gamma^\mu (1 - \gamma^5) v_{\bar{\nu}_\mu} \langle 0 | J_\mu^{h\dagger} | \pi^-(q) \rangle
 \end{aligned}$$

- $\langle 0 | J_\mu^{h\dagger} | \pi^-(q) \rangle$ involves strong interaction bound state \rightarrow hard to calculate (recent: lattice QCD calculation)

- However, $J_\mu^{h\dagger}$ is Lorentz vector \rightarrow

$$\langle 0 | J_\mu^{h\dagger} | \pi^-(q) \rangle = \underbrace{-i \cos \theta_c}_{\text{convention}} f_\pi q_\mu$$

- $q = p_1 + p_2$ is only 4-vector available
- $f_\pi \equiv$ “pion decay constant”
- Related to pion $\bar{u}d$ “wave function”
(expect $f_\pi = O(\Lambda_{QCD}) = O(100 \text{ MeV})$)
- f_π could depend on q^2 , but $q^2 = m_\pi^2 =$
fixed



- Can show (exercise) that $\langle 0 | V_\mu^\dagger | \pi^-(q) \rangle = 0$ using parity invariance
(weak P violation already included explicitly to 1st order \rightarrow strong interaction calculation)

- Rate (exercise)

$$\Gamma(\pi \rightarrow \mu \nu) = \underbrace{\frac{G_F^2 \cos^2 \theta_c}{8\pi} f_\pi^2 m_\mu^2}_{\text{matrix element}} \underbrace{m_\pi \left(1 - \frac{m_\mu^2}{m_\pi^2}\right)}_{\text{phase space}}$$

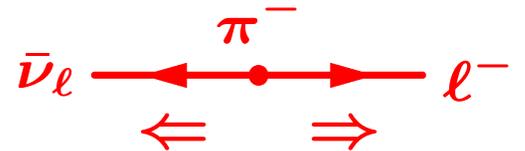
- Experiment: $\tau_{\pi^\pm} \sim 2.60 \times 10^{-8} \text{ s}$
 - $G_F \sim 1.17 \times 10^{-5} \text{ GeV}^{-2}$ from μ lifetime
 - $\cos \theta_c \sim 0.975$ from superallowed β decay
- $\Rightarrow f_\pi \sim 132 \text{ MeV} \sim 0.95 m_\pi$

- e/μ ratio

$$\frac{\Gamma(\pi \rightarrow e\nu)}{\Gamma(\pi \rightarrow \mu\nu)} = \frac{m_e^2}{m_\mu^2} \left(\frac{m_\pi^2 - m_e^2}{m_\pi^2 - m_\mu^2} \right) (1 + \underbrace{O(\alpha)}_{\text{rad. corr.}})$$

$$= 1.28 \times 10^{-4} (1 + O(\alpha)) = 1.23 \times 10^{-4}$$

- Experiment: $(1.23 \pm 0.02) \times 10^{-4}$
- $V - A$ favors $h_{\ell^-} = -1$
- Angular momentum forces “wrong helicity”
- Amplitude suppressed by m_ℓ/E_ℓ



- Kaon decays

$K^+ \rightarrow \mu^+ \nu_\mu$	$(K_{\mu 2})$	63.5%	f_K / f_π
$\rightarrow e^+ \nu_e$	$(K_{e 2})$	1.6×10^{-5}	
$\rightarrow \mu^+ \nu_\mu \pi^0$	$(K_{\mu 3})$	3.2%	universality test
$\rightarrow e^+ \nu_e \pi^0$	$(K_{e 3})$	4.8%	universality test
$\rightarrow \pi^+ \pi^0$	$(K_{2\pi})$	21.2%	nonleptonic
$\rightarrow \pi^+ \pi^+ \pi^-$	$(K_{3\pi})$	5.6%	
$\rightarrow \pi^+ \pi^0 \pi^0$	$(K_{3\pi})$	1.9%	

- K is also pseudoscalar, related to π by flavor $SU(3)$

$$\langle 0 | J_\mu^{h\dagger} | K^-(q) \rangle = \underbrace{-i \sin \theta_c}_{\text{convention}} f_K q_\mu$$

- $f_K =$ kaon decay constant
- In $SU(3)$ limit $f_K = f_\pi$
- However, $SU(3)$ typically broken by 20-30%

$$\Gamma(K \rightarrow \mu \nu) = \frac{G_F^2 \sin^2 \theta_c}{8\pi} f_K^2 m_\mu^2 m_K \left(1 - \frac{m_\mu^2}{m_K^2} \right)$$

- Observed rate $+\sin \theta_c \simeq 0.221$ from $K_{\ell 3}$ or hyperon decay:

$$f_K \sim 160 \text{ MeV} \sim (1.22 \pm 0.01) f_\pi$$

(20% $SU(3)$ breaking)

- Pion beta decay (π_{e3}), $\pi^\pm \rightarrow \pi^0 e^\pm \bar{\nu}_e^{(-)}$

$$M = -i \frac{G_F}{\sqrt{2}} \bar{u}_e \gamma^\mu (1 - \gamma^5) v_{\bar{\nu}_e} \langle \pi^0(p_2) | J_\mu^{h\dagger} | \pi^-(p_1) \rangle$$

- Parity: only vector current $V_\mu^\dagger \equiv \cos \theta_c \bar{u} \gamma_\mu d$ contributes
- Lorentz invariance: only two momenta, $p_{1,2}$ in hadronic matrix element \rightarrow

$$\langle \pi^0(p_2) | V_\mu^\dagger | \pi^-(p_1) \rangle = \cos \theta_c [f_+(q^2)(p_{1\mu} + p_{2\mu}) + f_-(q^2)(p_{1\mu} - p_{2\mu})]$$

- $f_\pm(q^2)$ are form factors, which can depend on $q^2 \equiv (p_2 - p_1)^2$

- Are $f_{\pm}(q^2)$ totally unknown because of strong interactions? No!
- Symmetry rescues us. Strong interactions almost invariant under global $SU(2)$ isospin symmetry (broken at 1% level by electromagnetism and by $m_d - m_u \sim \text{few MeV} \neq 0$)
- $V_{\mu}^{\dagger} / \cos \theta_c \equiv \bar{u} \gamma_{\mu} d$ is generator of isospin

$$\langle \pi^0(p_2) | V_{\mu}^{\dagger} | \pi^{-}(p_1) \rangle = \cos \theta_c [f_{+}(q^2)(p_{1\mu} + p_{2\mu}) + f_{-}(q^2)(p_{1\mu} - p_{2\mu})]$$

- If isospin were exact, $|f_{+}(0)| = \sqrt{2}$. Also, $\partial^{\mu} V_{\mu}^{\dagger} = 0 \Rightarrow f_{-}(q^2) = 0$ (cf non-renormalization of electric charge)

- **Conserved vector current (CVC):** $\bar{u}\gamma_\mu d$, $\bar{d}\gamma_\mu u$ and $\frac{1}{2}(\bar{u}\gamma_\mu u - \bar{d}\gamma_\mu d)$ (isovector part of $J_\mu^{elm} = \frac{1}{2}(\bar{u}\gamma_\mu u - \bar{d}\gamma_\mu d) + \frac{1}{6}(\bar{u}\gamma_\mu u + \bar{d}\gamma_\mu d)$) are related by isospin and have same form factors (Was hypothesis. Natural in quark model.)
- **Ademollo-Gatto theorem:** corrections to $f_+(0)$ from isospin breaking are *second order*, $\sim 10^{-3}$ (negligible)
- For π_{e3} , $\Delta = m_{\pi^+} - m_{\pi^0} \sim 4.6 \text{ MeV} \rightarrow$ can neglect $p_2 - p_1$ term and take $f_+(q^2) \sim f_+(0)$. Neglecting m_e ,

$$\Gamma(\pi^\pm \rightarrow \pi^0 e^\pm \nu_e^{(-)}) = \frac{G_F^2 \cos^2 \theta_c |f_+(0)|^2 \Delta^5}{60\pi^3} \Rightarrow |f_+(0)| = 1.37 \pm 0.02$$

Isospin: $|f_+(0)| = \sqrt{2}$ + second order

- $K_{\ell 3}$ decays: $K^+ \rightarrow \ell^+ \nu_\ell \pi^0$, $K^0 \rightarrow \ell^+ \nu_\ell \pi^-$, etc

$$\langle \pi^0(p_2) | V_\mu | K^+(p_1) \rangle = |V_{us}| [f_+^{K^+}(q^2)(p_{1\mu} + p_{2\mu}) + f_-^{K^+}(q^2)(p_{1\mu} - p_{2\mu})]$$

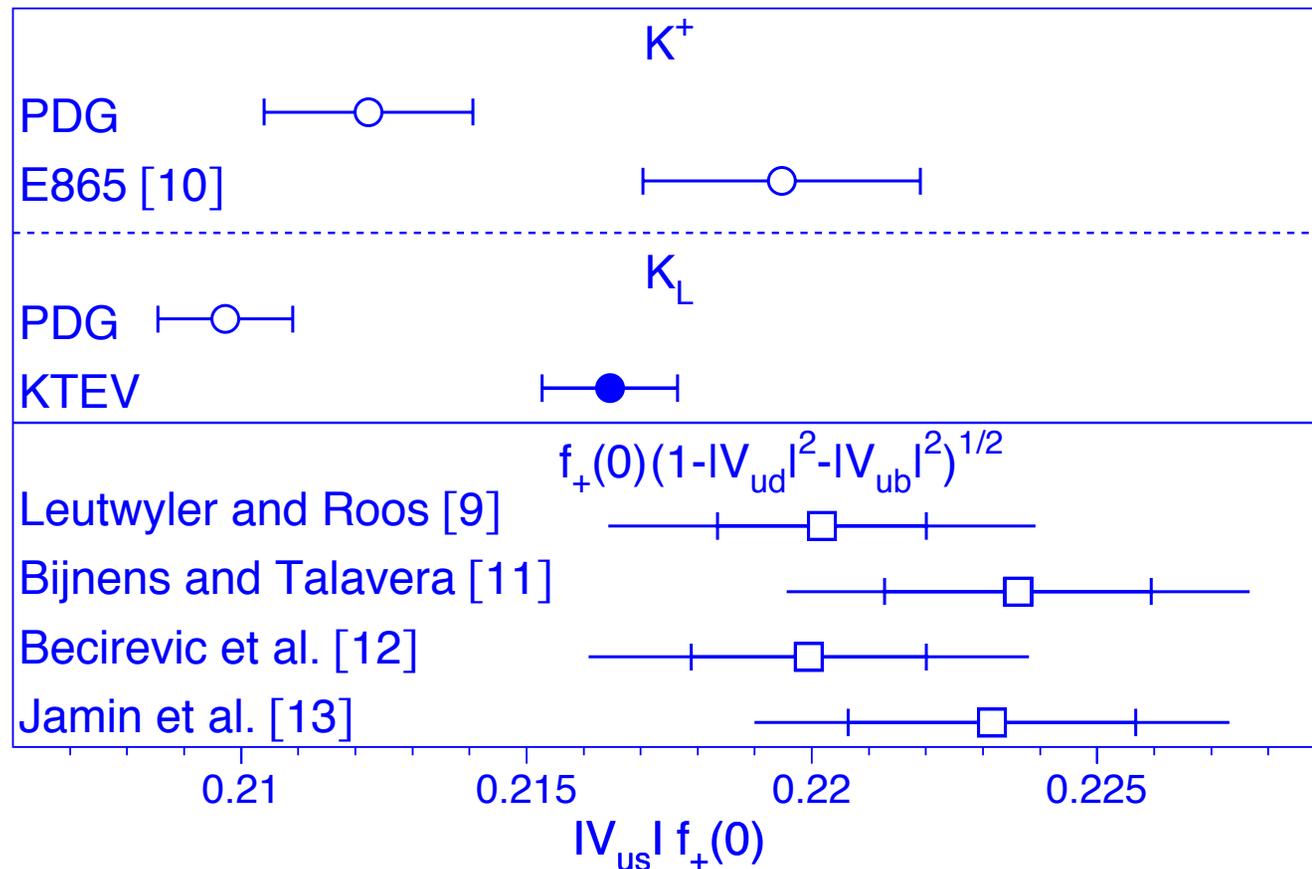
- Write $|V_{us}|$ rather than $\sin \theta_c$ (most precise determination)
- Large energy release. Cannot neglect q^2 dependence of form factors or f_- term
- Can measure them (linear or quadratic approximation) from decay distributions, but need $|f_+^{K^+}(0)|$ and $|f_+^{K^0}(0)|$
- V_μ are $SU(3)$ generators \rightarrow

$$\sqrt{2}|f_+^{K^+}(0)| \sim |f_+^{K^0}(0)| \sim 1 + O(\epsilon^2)$$

where $\epsilon \sim 20\%$ is typical $SU(3)$ breaking (Ademollo-Gatto)

- Detailed estimates (chiral perturbation theory): $|f_+^{K^0}(0)| \sim 0.961(8)$, $\sqrt{2}|f_+^{K^+}(0)| \sim 1.022|f_+^{K^0}(0)|$

- V_{us} from K_{l3} extremely important for weak universality test
- Recent significant shift, KTEV (FNAL) and E865 (BNL)



(from KTEV, PRL 93, 181802)

β Decay and Related Processes

$n \rightarrow p e^- \bar{\nu}_e$	neutron
$(N, Z) \rightarrow (N - 1, Z + 1) e^- \bar{\nu}_e$	nuclear (heavy)
$(N, Z) \rightarrow (N + 1, Z - 1) e^+ \nu_e$	nuclear (light, e.g., Sun)
$\ell^- p \rightarrow n \nu_\ell$	atomic e^- or μ^- capture
$\nu_e n \leftrightarrow e^- p, \quad \bar{\nu}_e p \leftrightarrow e^+ n$	inverse β decay

$$H = \frac{G_F}{\sqrt{2}} V_{ud} \bar{e} \gamma^\mu (1 - \gamma^5) \nu_e \bar{u} \gamma_\mu (1 - \gamma^5) d + HC$$

- Radiative corrections divergent in Fermi theory, finite in $SU(2) \times U(1)$

- Nuclear filter: V_μ (Fermi transition), A_μ (Gamow-Teller transition), or both relevant
- Superaligned: $0_i^+ \rightarrow 0_f^+$ in same isomultiplet
 - Pure Fermi transition
 - $\langle 0_f^+ | \bar{u} \gamma_0 d | 0_i^+ \rangle = 1 + O(\delta^2)$, where $\delta =$ isospin breaking (Ademollo-Gatto). Corrections tiny but critical.
 - Best determination of $|V_{ud}| = 0.97377(11)(15)(19)$ (various theoretical uncertainties) (Marciano, Sirlin, hep-ph/0510099)

- Neutron β decay

$$\begin{aligned}
 & \langle p | \bar{u} \gamma_\mu (1 - \gamma^5) d | n \rangle \\
 &= \bar{u}_p \left[\gamma_\mu f_1(q^2) + \underbrace{\frac{i\sigma_{\mu\nu} q^\nu}{2m} f_2(q^2)}_{\text{weak magnetism}} + \underbrace{q_\mu f_3(q^2)}_{\text{vanishes by CVC}} \right. \\
 & \quad \left. - \gamma_\mu \gamma^5 g_1(q^2) - \underbrace{\frac{\sigma_{\mu\nu} \gamma^5 q^\nu}{2m} g_2(q^2)}_{\text{T-odd}} - \underbrace{q_\mu \gamma^5 g_3(q^2)}_{\text{induced pseudoscalar}} \right] u_n \\
 & \sim \bar{u}_p \gamma_\mu (g_V - g_A \gamma^5) u_n
 \end{aligned}$$

- $q = p_p - p_n$, $g_V \sim 1 + O(\delta^2)$
- $g_A \neq 1$, since no symmetry prevents large strong interaction effect (Adler-Weisberger estimate: $g_A = O(1.25)$)

- Lifetime $\propto G_F^2 |V_{ud}|^2 m_e^5 (g_V^2 + 3g_A^2)$
- Asymmetries w.r.t. n polarization, $e - \bar{\nu}$ correlation
(e.g., ILL (Grenoble)) $\rightarrow g_A/g_V$

$$|V_{ud}| = 0.9725(13)$$

$$g_A/g_V = 1.2695(29)$$

CKM Universality

- Universality of quark and lepton couplings (required by gauge theory) and unitarity of CKM relate μ , β , $K_{\ell 3}$, and $b \rightarrow u\ell^- \bar{\nu}_\ell$
- Measure $G_F |V_{ui}|$. Divide by G_F from μ decay $\rightarrow |V_{ui}|$
- Expect

$$\underbrace{|V_{ud}|^2}_{\beta} + \underbrace{|V_{us}|^2}_{K_{\ell 3}} + \underbrace{|V_{ub}|^2}_{\text{negligible}} = 1$$

- Severe constraint on new interactions, W_R , $W_L - W_R$ mixing, fermion-exotic mixing, neutrino-exotic mixing (excludes some explanations of NuTeV)
 - For many years, appeared to be 2σ discrepancy, $\sum |V_{ui}|^2 \sim 0.9969(15)$
 - Particle physicists doubted nuclear physics ($|V_{ud}|$)
 - But problem was really in V_{us} from $K_{\ell 3}$!
- Current: $\sum |V_{ui}|^2 \sim 0.9992(5)_{ud}(4)_{us}(8)_{f_+(0)}$
(M.S., hep-ph/0510099)
- Without radiative corrections, would find ~ 1.04 ! These are only finite and meaningful in full $SU(2) \times U(1)$ gauge theory

Quantum Electrodynamics (QED)

Incorporated into standard model

Lagrangian:

$$\mathcal{L} = -\frac{gg'}{\sqrt{g^2 + g'^2}} J_Q^\mu (\cos \theta_W B_\mu + \sin \theta_W W_\mu^3)$$

Photon field ($A = \gamma$):

$$A_\mu = \cos \theta_W B_\mu + \sin \theta_W W_\mu^3$$

Positron electric charge: $e = g \sin \theta_W$, where $\tan \theta_W \equiv g'/g$

Electromagnetic current:

$$\begin{aligned} J_Q^\mu &= \sum_{m=1}^F \left[\frac{2}{3} \bar{u}_m^0 \gamma^\mu u_m^0 - \frac{1}{3} \bar{d}_m^0 \gamma^\mu d_m^0 - \bar{e}_m^0 \gamma^\mu e_m^0 \right] \\ &= \sum_{m=1}^F \left[\frac{2}{3} \bar{u}_m \gamma^\mu u_m - \frac{1}{3} \bar{d}_m \gamma^\mu d_m - \bar{e}_m \gamma^\mu e_m \right] \end{aligned}$$

Flavor diagonal: Same form in weak and mass bases because fields which mix have same charge

Purely vector (parity conserving): L and R fields have same charge

Quantum Electrodynamics

Experiment	Value of α^{-1}		Difference from $\alpha^{-1}(a_e)$
Deviation from gyromagnetic ratio, $a_e = (g - 2)/2$ for e^-	137.035 999 58 (52)	$[3.8 \times 10^{-9}]$	–
ac Josephson effect	137.035 988 0 (51)	$[3.7 \times 10^{-8}]$	$(0.116 \pm 0.051) \times 10^{-4}$
h/m_n (m_n is the neutron mass) from n beam	137.036 011 9 (51)	$[3.7 \times 10^{-8}]$	$(-0.123 \pm 0.051) \times 10^{-4}$
Hyperfine structure in muonium, μ^+e^-	137.035 993 2 (83)	$[6.0 \times 10^{-8}]$	$(0.064 \pm 0.083) \times 10^{-4}$
Cesium D_1 line	137.035 992 4 (41)	$[3.0 \times 10^{-8}]$	$(0.072 \pm 0.041) \times 10^{-4}$

Spectacularly successful:

Most precise: e anomalous magnetic moment $\rightarrow \alpha$

Many low energy tests to few $\times 10^{-8}$

$$m_\gamma < 6 \times 10^{-17} \text{ eV}$$

$$q_\gamma < 5 \times 10^{-30} |e|$$

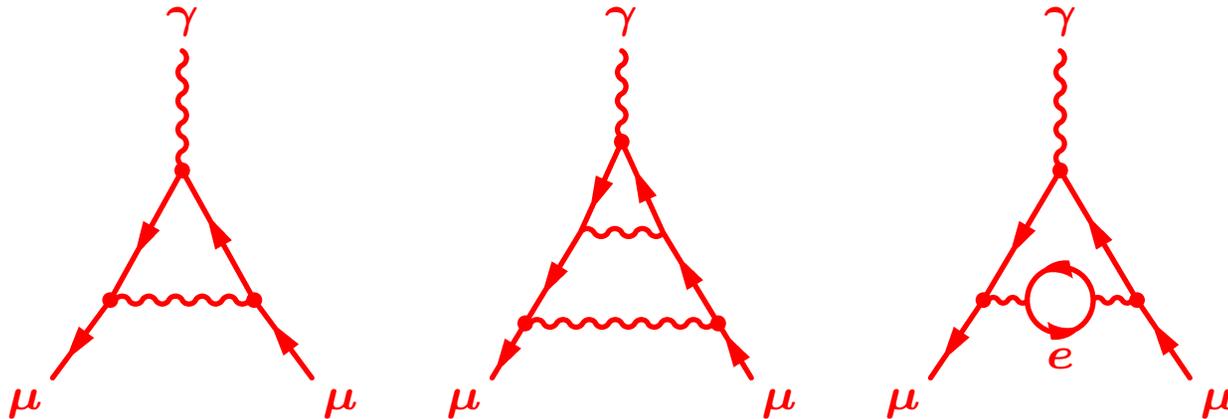
High energy well-measured (PEP, PETRA, TRISTAN, LEP)

The Anomalous Magnetic Moment of the Muon

- Muon $a_\mu \equiv \frac{g_\mu - 2}{2}$ sensitive to new physics (usually $\sim (m_\mu/M_X)^2$)

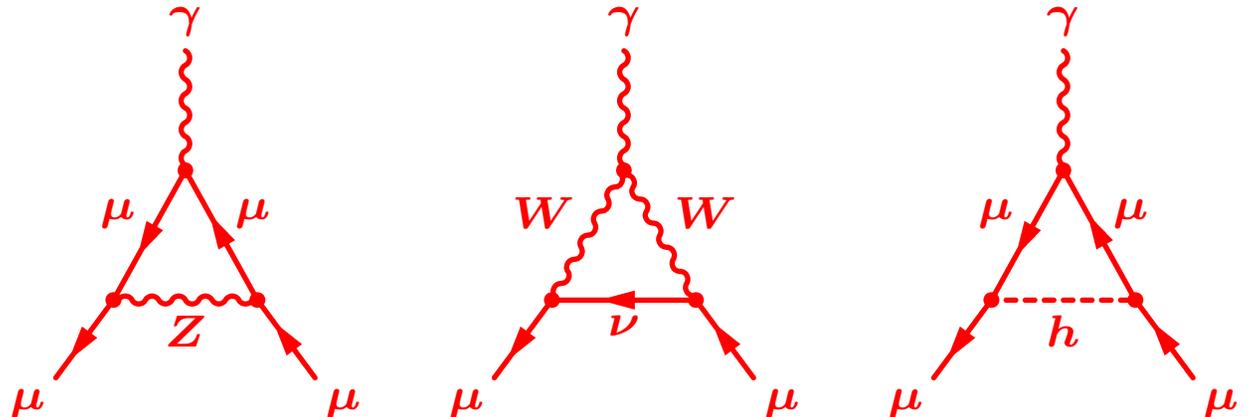
$$a_\mu^{\text{SM}} = a_\mu^{\text{QED}} + a_\mu^{\text{Had}} + a_\mu^{\text{EW}}$$

- a_μ^{QED} known to four loops (3 analytic); leading logs to five



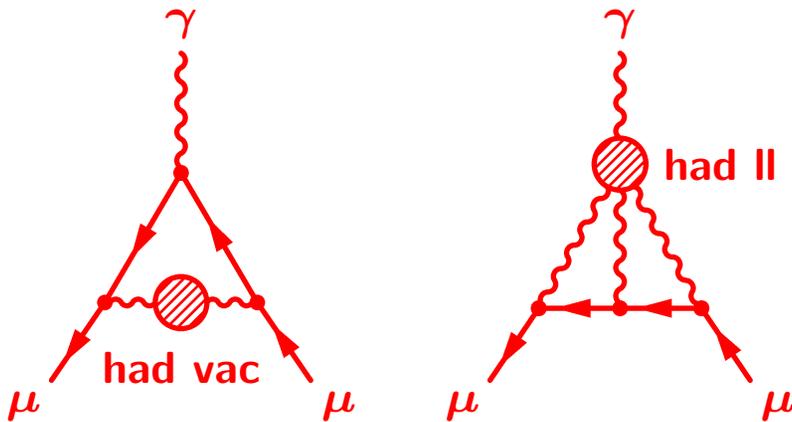
$$\begin{aligned}
a_{\mu}^{\text{QED}} = & \frac{\alpha}{2\pi} + 0.765857376(27) \left(\frac{\alpha}{\pi}\right)^2 \\
& + 24.05050898(44) \left(\frac{\alpha}{\pi}\right)^3 + 126.07(41) \left(\frac{\alpha}{\pi}\right)^4 \\
& + 930(170) \left(\frac{\alpha}{\pi}\right)^5 = 1165847.06(3) \times 10^{-9}
\end{aligned}$$

- $a_{\mu}^{\text{EW}} = 1.52(3) \times 10^{-9}$ (goal of experiments) **includes leading 2 and 3 loops** (cancellation)



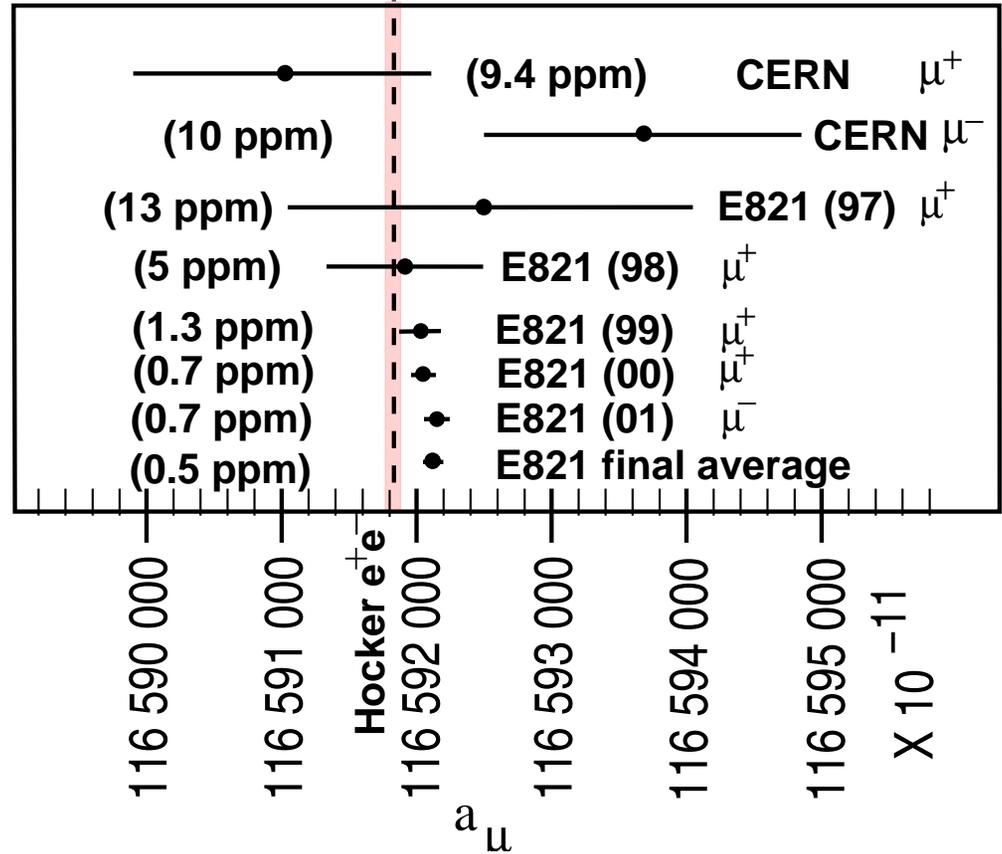
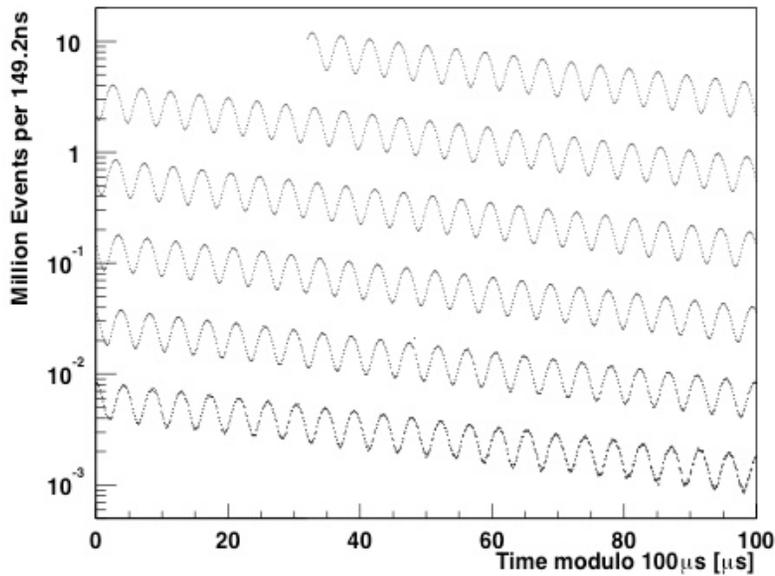
- Biggest uncertainty: $\alpha_{\mu}^{\text{Had}}$ = hadronic vacuum polarization (2 loop) and hadronic light by light (3 loop)

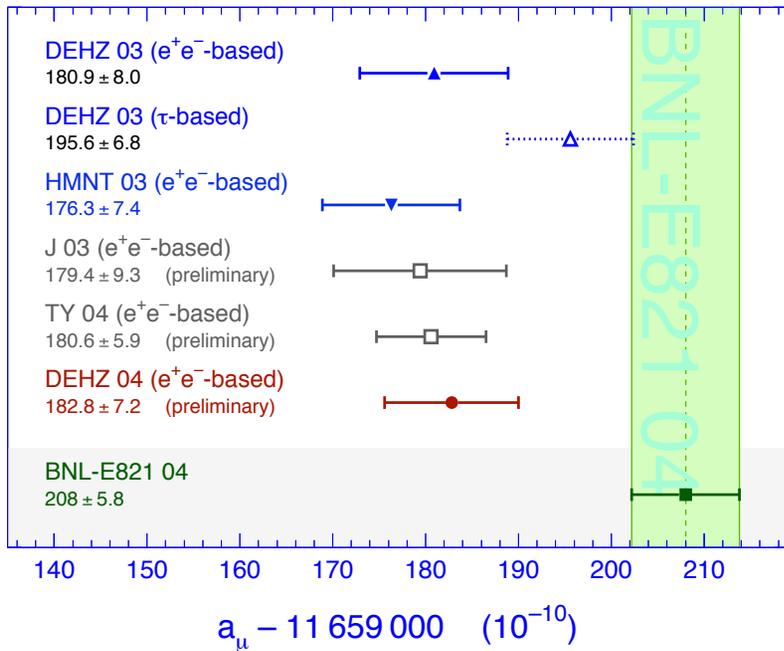
$$\alpha_{\mu}^{\text{Had vac}} = \frac{1}{3} \left(\frac{\alpha}{\pi} \right)^2 \int_{4m_{\pi}^2}^{\infty} \frac{ds}{s} \underbrace{K(s)}_{\text{fnc of } m_{\mu}^2/s} \frac{\sigma(e^{+}e^{-} \rightarrow \text{had})}{\sigma(e^{+}e^{-} \rightarrow \mu^{+}\mu^{-})}$$



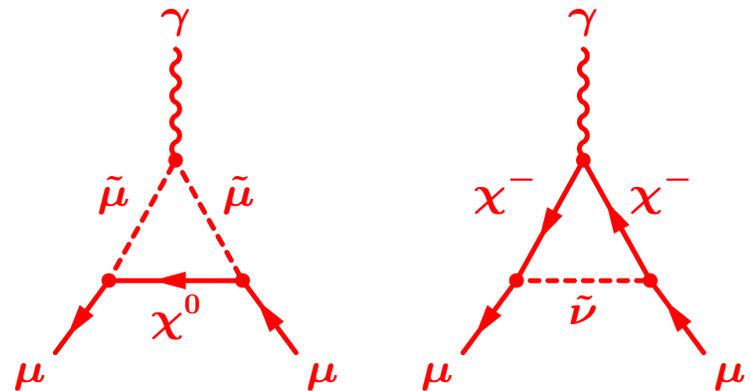
- $\alpha_{\mu}^{\text{Had vac}}$: discrepancy between $e^{+}e^{-}$ (new KLOE, SND) and τ decay (isospin violation?)
- $\alpha_{\mu}^{\text{Had l.l.}}$ sign now settled down. Small but non-negligible

● $a_{\mu}^{\text{exp}} = 1165920.80(63) \times 10^{-9}$ (dominated by BNL 821)





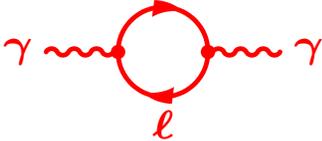
- e^+e^- data: $a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = 2.5(9) \times 10^{-9} \quad (2.7\sigma)$
- τ decay data: no discrepancy (0.7σ)
- Supersymmetry: central value (e^+e^-) for $m_{\text{SUSY}} \sim 72\sqrt{\tan\beta}$ GeV



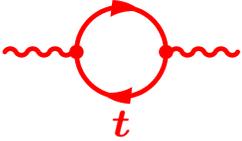
Running of α

- Largest theory uncertainty in $M_Z - \sin^2 \theta_W$ (cf. a_μ^{had})

$$\alpha(M_Z^2) = \frac{\alpha}{1 - \Delta\alpha}$$



ℓ



t



had

$$\begin{aligned} \Delta\alpha &= \Delta\alpha_\ell + \Delta\alpha_t + \Delta\alpha_{\text{had}}^{(5)} \\ &\sim 0.03149769 - 0.000070(5) + \Delta\alpha_{\text{had}}^{(5)} \end{aligned}$$

$$\alpha^{-1} \sim 137.036$$

- **Calculation of $\Delta\alpha_{\text{had}}^{(5)}$**
 - **Data driven: R_{had} up to ~ 40 GeV; PQCD above**
 - **Theory driven: PQCD + NPQCD (OPE, sum rules) above ~ 2 GeV \rightarrow smaller uncertainties**
 - **e^+e^- data vs τ decays**
 - **Table for $\alpha_s = 0.120$**
 - **Correlation with $g_\mu - 2$**
- **$\alpha(M_Z)^{-1} = 128.954 \pm 0.031$**

Reference	Result	Comment
Martin & Zeppenfeld [21]	0.02744 ± 0.00036	PQCD for $\sqrt{s} > 3$ GeV
Eidelman & Jegerlehner [22]	0.02803 ± 0.00065	PQCD for $\sqrt{s} > 40$ GeV
Geshkenbein & Morgunov [13]	0.02780 ± 0.00006	$\mathcal{O}(\alpha_s)$ resonance model
Burkhardt & Pietrzyk [24]	0.0280 ± 0.0007	PQCD for $\sqrt{s} > 40$ GeV
Swartz [25]	0.02754 ± 0.00046	use of fitting function
Alemany, Davier, Höcker [26]	0.02816 ± 0.00062	includes τ decay data
Krasnikov & Rodenbergh [27]	0.02737 ± 0.00039	PQCD for $\sqrt{s} > 2.3$ GeV
Davier & Höcker [28]	0.02784 ± 0.00022	PQCD for $\sqrt{s} > 1.8$ GeV
Kühn & Steinhauser [29]	0.02778 ± 0.00016	complete $\mathcal{O}(\alpha_s^2)$
Erlar [16]	0.02779 ± 0.00020	converted from $\overline{\text{MS}}$ scheme
Davier & Höcker [30]	0.02770 ± 0.00015	use of QCD sum rules
Groote <i>et al.</i> [31]	0.02787 ± 0.00032	use of QCD sum rules
Martin, Outhwaite, Ryskin [32]	0.02741 ± 0.00019	includes new BES data
Burkhardt & Pietrzyk [33]	0.02763 ± 0.00036	PQCD for $\sqrt{s} > 12$ GeV
de Troconiz & Yndurain [34]	0.02754 ± 0.00010	PQCD for $s > 2$ GeV ²
Jegerlehner [35]	0.02765 ± 0.00013	converted from MOM scheme
Hagiwara <i>et al.</i> [36]	0.02757 ± 0.00023	PQCD for $\sqrt{s} > 11.09$ GeV
Burkhardt & Pietrzyk [37]	0.02760 ± 0.00035	includes KLOE data

- Measurements: DORIS, PEP, PETRA, TRISTAN, LEP II

