

Finding the Higgs boson

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FNAL LHC School, Lecture 2

✍ Properties of the Higgs boson

Theoretical uncertainties & motivations for
precision measurements

✍ Higgs production at the Tevatron and LHC

Who needs a Higgs?

- ✍ Gives masses to W , Z , and fermions in gauge invariant fashion
 - ✍ Unitarizes VV ? VV scattering
 - ✍ (More in Lecture 3)
 - ✍ Makes precision electroweak data consistent
-

But....

- ✍ Higgs mechanism doesn't explain masses or flavor structure
 - ✍ It accommodates them
 - ✍ Higgs mass is quadratically sensitive to physics at high scales
 - ✍ (More in Lecture 3)
 - ✍ Higgs potential stable only for certain Higgs masses
 - ✍ (More in Lecture 3)
-

Review of Higgs Couplings

- Higgs couples to fermion mass
 - Largest coupling is to heaviest fermion

$$L \bar{f} \frac{m_f}{v} f h \quad \bar{f}_L f_R \quad \bar{f}_R f_L h$$

$$v=246 \text{ GeV}$$

- Top-Higgs coupling plays special role?
- No Higgs coupling to neutrinos

Review of Higgs Couplings

- Higgs couples to gauge boson masses

$$L \supset g M_W W_\mu^a W_\mu^a h + \frac{g M_Z}{\cos \theta_W} Z_\mu Z_\mu h + \dots$$

- Only free parameter is Higgs mass!
- Everything is calculable.... **testable theory**

$$g^2 = \frac{G_F}{\sqrt{2}} 8 M_W^2 = \frac{e^2}{\sin^2 \theta_W} = \frac{4}{\sin^2 \theta_W}$$

Review of Higgs Boson Feynman Rules

- Couplings to EW gauge bosons ($V = W, Z$):

$$\begin{aligned}
 & \text{Diagram 1: } V^\mu \text{ and } V^\nu \text{ wavy lines meet at a vertex with a dashed line } H \text{ extending to the right.} \\
 & \text{Equation: } H = 2i \frac{M_V^2}{v} g^{\mu\nu} \\
 & \text{Diagram 2: } V^\mu \text{ and } V^\nu \text{ wavy lines meet at a vertex with two dashed lines } H \text{ extending to the right.} \\
 & \text{Equation: } = 2i \frac{M_V^2}{v^2} g^{\mu\nu}
 \end{aligned}$$

✍ Higgs couples to heavy particles

- Couplings to fermions ($f = l, q$):

$$\text{Diagram: } f \text{ and } \bar{f} \text{ lines meet at a vertex with a dashed line } H \text{ extending to the right.} \\
 \text{Equation: } H = -i \frac{m_f}{v}$$

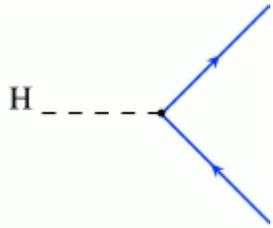
✍ No tree level coupling to gluons (g) or photons (?)

✍ $M_h^2 = 2v^2$? ? large M_h is strong coupling regime

- Self-couplings:

$$\begin{aligned}
 & \text{Diagram 1: } H \text{ and } H \text{ dashed lines meet at a vertex with a dashed line } H \text{ extending to the right.} \\
 & \text{Equation: } H = -3i \frac{M_H^2}{v} \\
 & \text{Diagram 2: } H \text{ and } H \text{ dashed lines meet at a vertex with two dashed lines } H \text{ extending to the right.} \\
 & \text{Equation: } = -3i \frac{M_H^2}{v^2}
 \end{aligned}$$

Higgs Decays



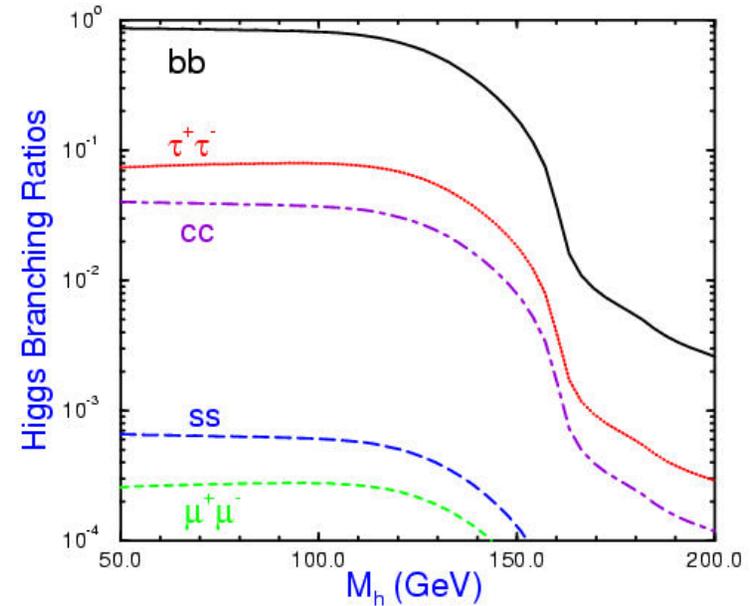
$$\Gamma(h \rightarrow f\bar{f}) \propto \frac{N_c G_F^2 m_f^2 M_h}{4\sqrt{2}}$$

$$\Gamma_f \propto \sqrt{1 - \frac{4m_f^2}{M_h^2}}$$

$\Gamma(h \rightarrow f\bar{f})$ proportional to m_f^2

$$\frac{BR(h \rightarrow b\bar{b})}{BR(h \rightarrow \tau^+\tau^-)} \propto N_c \frac{m_b^2}{m_\tau^2} \frac{\Gamma_b}{\Gamma_\tau}$$

Γ^3 typical of scalar
(pseudo-scalar decay ??)



For $M_h < 2M_W$, decays to $b\bar{b}$ most important

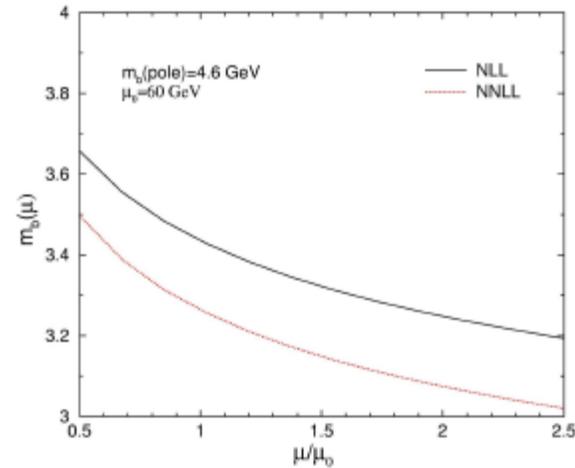
QCD Corrections to $h \rightarrow Q\bar{Q}$

✍ Tree level: $(h \rightarrow Q\bar{Q})_{tree} \approx \frac{3G_F M_h}{4\sqrt{2}} M_Q^2 \alpha_s^3$

✍ Add QCD: $(h \rightarrow Q\bar{Q})_{QCD} \approx \frac{3G_F M_h}{4\sqrt{2}} \bar{m}_Q^2(M_h) \alpha_s^3 \left[1 + 5.67 \frac{\alpha_s(M_h)}{\alpha_s} + \dots \right]$

✍ Large logs absorbed into running \overline{MS} mass:

$$m_b(\mu^2) \approx m_b(m_b^2) \left[\frac{\alpha_s(m_b^2)}{\alpha_s(\mu^2)} \right]^{12/23}$$



Higgs Decays to Gauge Bosons

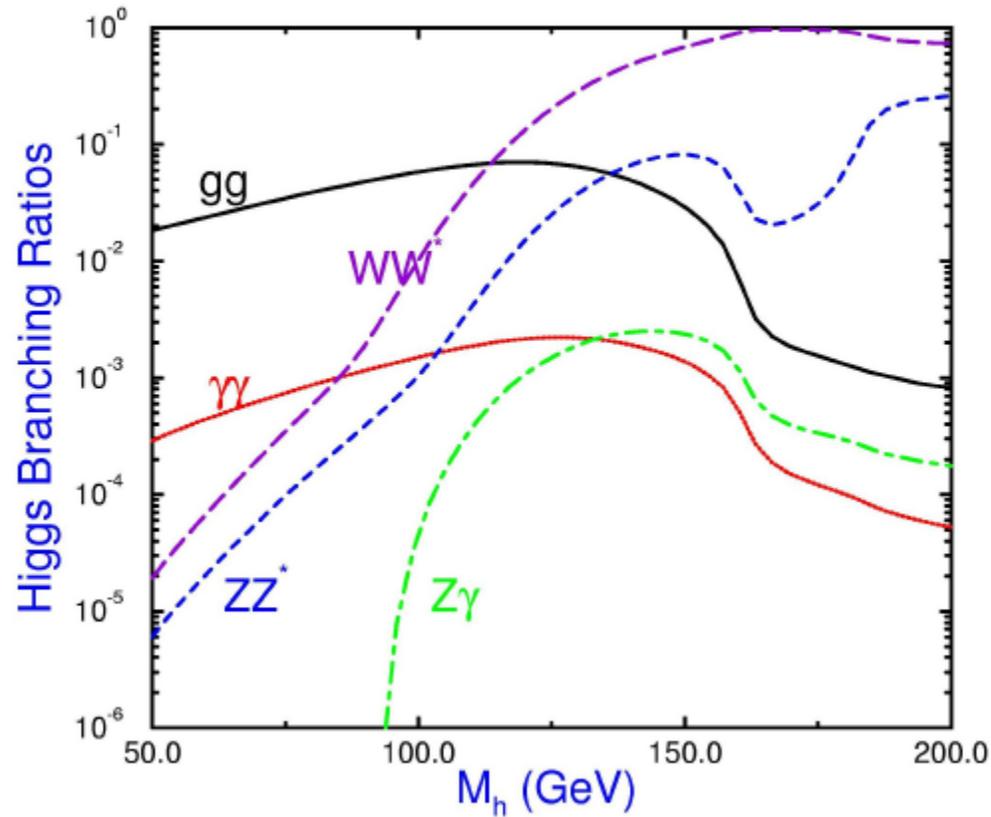
- ✍ h → gg sensitive to top loops
 - ✍ Remember no coupling at tree level
- ✍ h → ?? sensitive to W loops, only small contribution from top loops
- ✍ h → W⁺W⁻ → ffff has sharp threshold at 2 M_W, but large branching ratio even for M_h=130 GeV

$$\Gamma(h \rightarrow VV) \approx \frac{G_F M_h^3}{8\sqrt{2}} \sum_V \kappa_V^2 \dots \quad \kappa_{W,Z} \approx 2,1$$

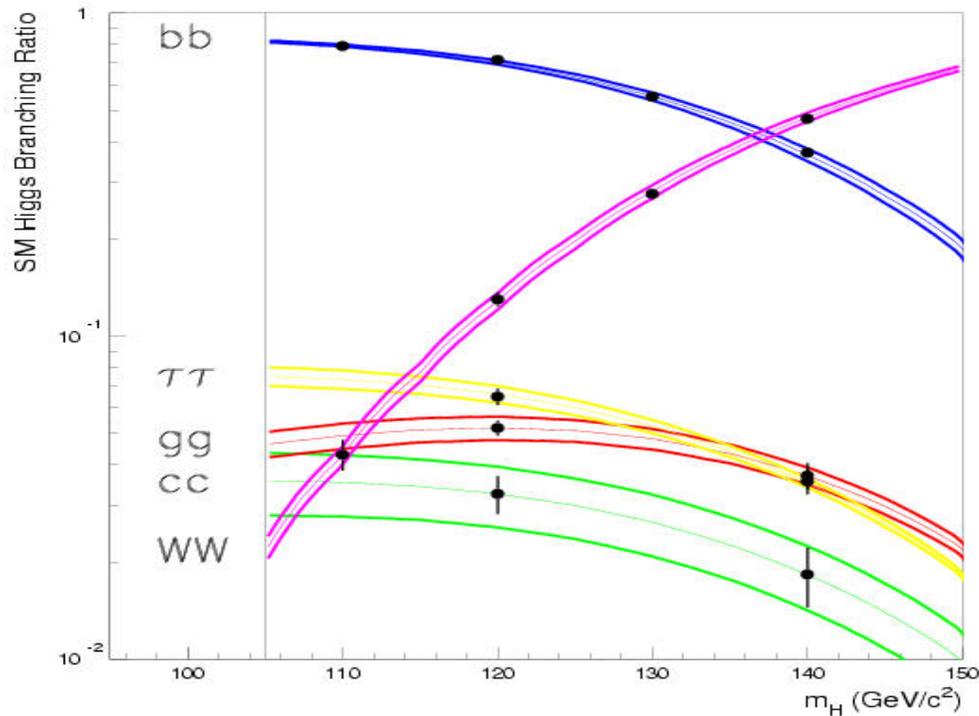
Cubic in M_h, so for heavy Higgs, decays to vector boson dominate

Decays to Gauge Bosons

Higgs Branching Ratios to Gauge Boson Pairs



Status of Theory for Higgs BRs



Bands show theory errors

Largest source of uncertainty is b quark mass

Data points are e^+e^- ILC at $\sqrt{s}=350$ GeV with $L=500 \text{ fb}^{-1}$

Total Higgs Width

- Total width sensitive function of M_h
- Small M_h , Higgs is narrower than detector resolution
- As M_h becomes large, width also increases

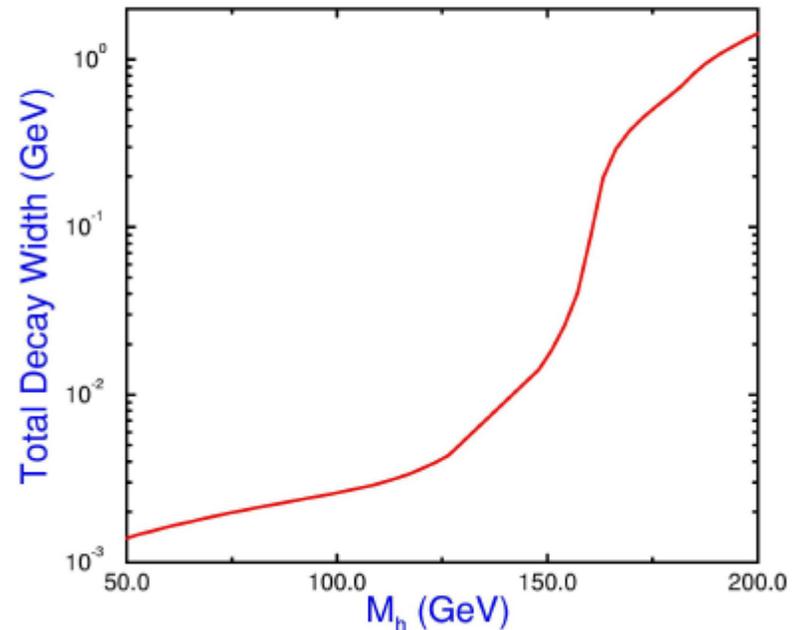
No clear resonance

For $M_h \approx 1.4 \text{ TeV}$, $\Gamma_{\text{tot}} \approx M_h$

$$\Gamma(h \rightarrow W^+W^-) \approx \frac{1}{16 \sin^2 \theta_w} \frac{M_h^3}{M_W^2}$$

$$\approx 330 \text{ GeV} \left(\frac{M_h}{1 \text{ TeV}} \right)^3$$

Higgs Boson Decay Width



• Higgs branching ratios easily computed with HDECAY program to NLO

• <http://mspira.home.cern.ch/mspira/proglist.html>

Higgs Searches at LEP2

- LEP2 searched for $e^+e^- \rightarrow Z h$
- Rate turns on rapidly after threshold, peaks just above threshold, $\sim s^{3/2}$
- Measure recoil mass of Higgs; **result independent of Higgs decay pattern**

$$P_{e^-} = \sqrt{s}/2(1, 0, 0, 1)$$

$$P_{e^+} = \sqrt{s}/2(1, 0, 0, -1)$$

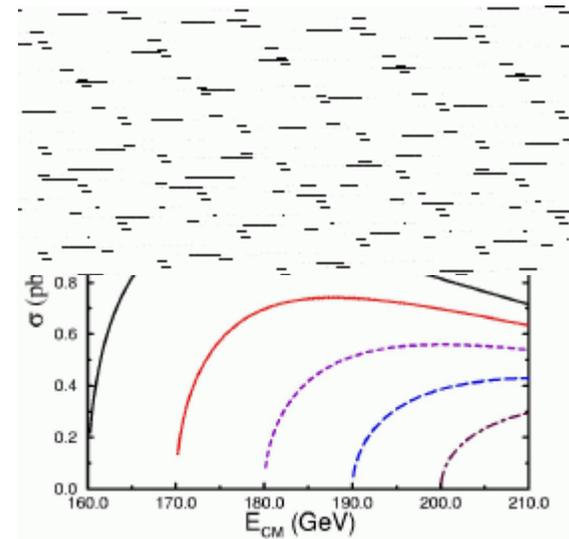
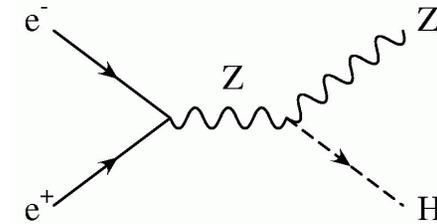
$$P_Z = (E_Z, \vec{p}_Z)$$

- Momentum conservation:

$$(P_{e^-} + P_{e^+} - P_Z)^2 = P_h^2 = M_h^2$$

$$s - 2 E_Z \sqrt{s} + M_Z^2 = M_h^2$$

- LEP2 limit, $M_h > 114.1 \text{ GeV}$



Higgs production at Hadron Colliders

- ✍ Many possible production mechanisms; Importance depends on:
 - ✍ Size of production cross section
 - ✍ Size of branching ratios to observable channels
 - ✍ Size of background
 - ✍ Importance varies with Higgs mass
 - ✍ Need to see more than one channel to establish Higgs properties and verify that it is a Higgs boson
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Production in Hadron Colliders

✎ Gluon fusion

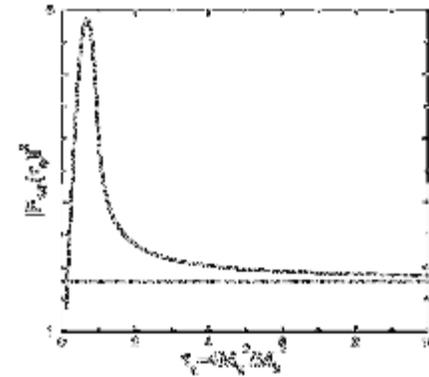
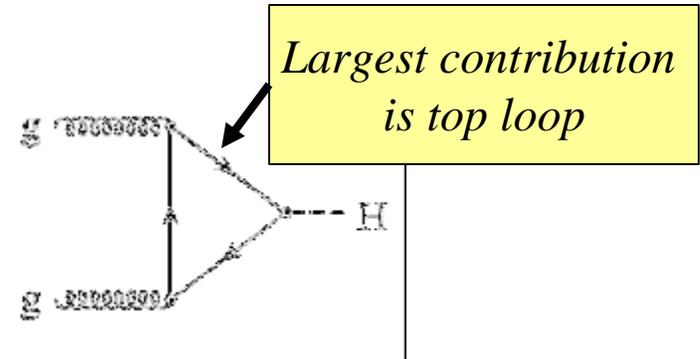
- ✎ Largest rate for all M_h at LHC
- ✎ Gluon-gluon initial state
- ✎ Sensitive to top quark Yukawa κ_t

✎ Lowest order cross section:

$$\sigma_0(gg \rightarrow h) = \frac{\kappa_s(\kappa_R)^2}{1024v^2} \left| \sum_q F_{1/2}(\tau_q) \right|^2 (M_h^2 \hat{s})$$

- ✎ $\tau_q = 4M_q^2/M_h^2$
- ✎ Light Quarks: $F_{1/2} \approx (M_b/M_h)^2 \log(M_b/M_h)$
- ✎ Heavy Quarks: $F_{1/2} \approx -4/3$

In SM, b-quark loops unimportant



Rapid approach to heavy quark limit

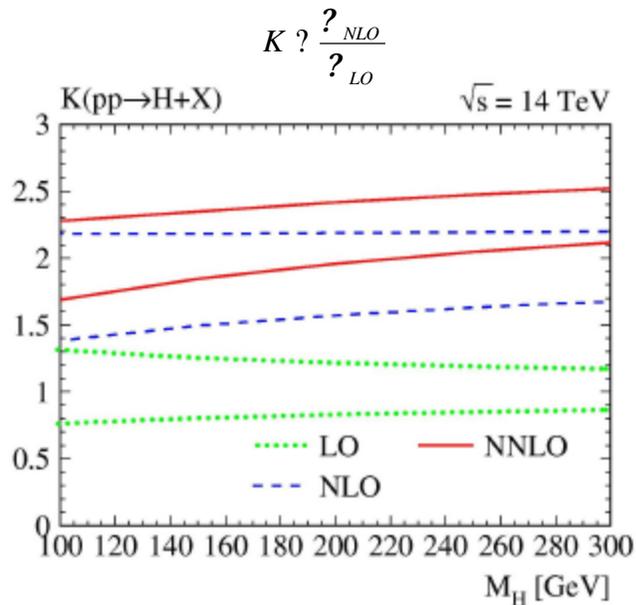
Gluon fusion, continued

- ✍ Integrate parton level cross section with gluon parton distribution functions

$$\sigma_0(pp \rightarrow h) = \sum_{i,j} \int_0^1 \int_0^1 \frac{dx}{x} g(x, \mu_F) g\left(\frac{z}{x}, \mu_F\right)$$

- ✍ $z = M_h^2/S$, S is hadronic center of mass energy
- ✍ Rate depends on μ_R, μ_F
- ✍ Rate for gluon fusion independent of M_t for $M_t \gg M_h$
- ✍ Counts number of heavy fermions

NNLO, gg? h



Rates depend on renormalization scale, $\mu_s(\mu_R)$, and factorization scale, $g(\mu_F)$

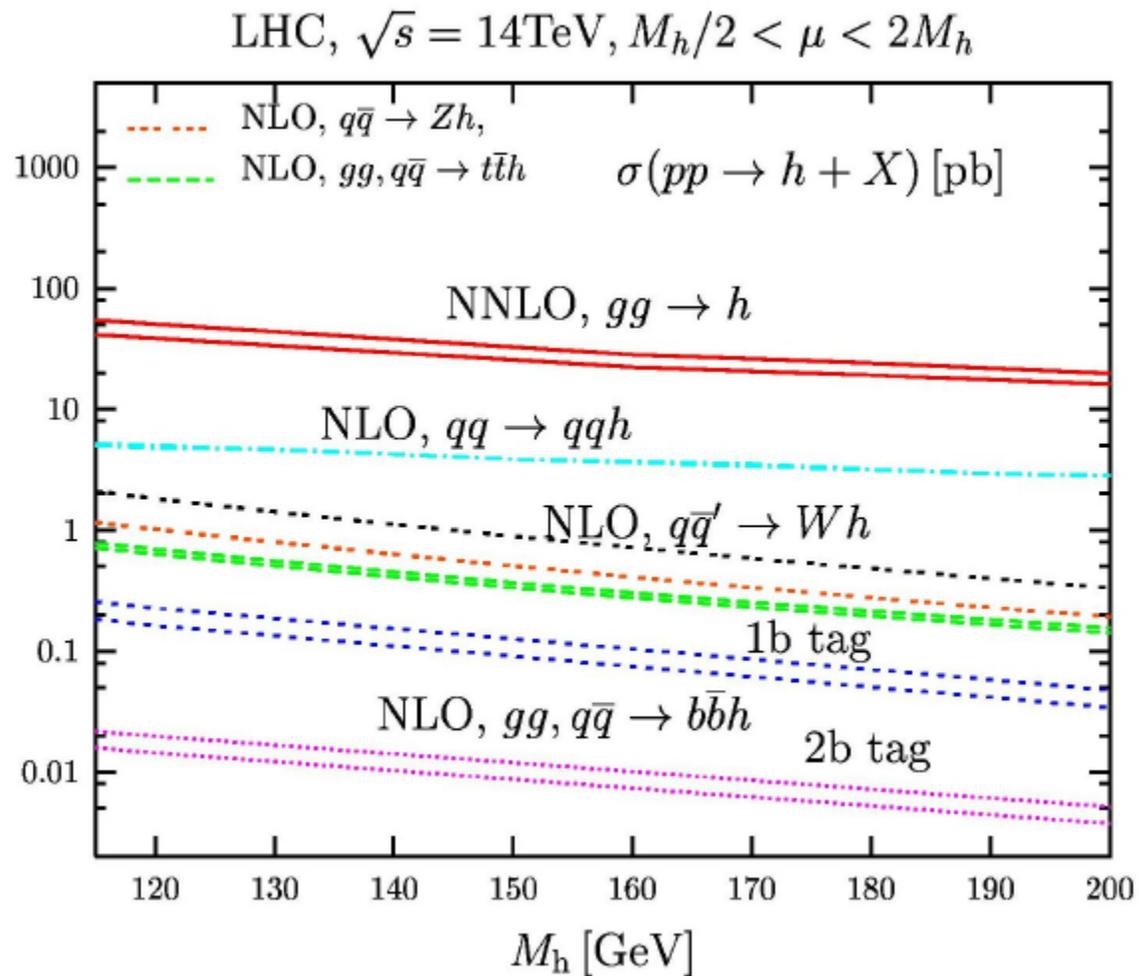
Bands show $.5M_h < \mu < 2 M_h$

LO and NLO μ dependence bands don't overlap

μ Dependence used as estimate of theoretical uncertainty

NLO&NNLO results allow improved estimates of theoretical uncertainties

Higgs production at the LHC



Vector Boson Fusion

- W+W- → X is a real process:

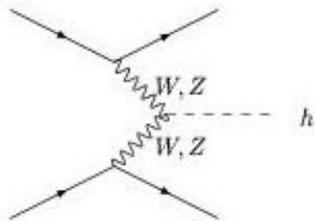
$$\sigma_{pp \rightarrow WW \rightarrow X}(s) \approx \int dz \frac{dL}{dz} \bigg|_{pp/WW} \sigma_{WW \rightarrow X}(zs)$$

- Rate increases at large s: $\propto (1/M_W^2) \log(s/M_W^2)$
- Integral of cross section over final state phase space has contribution from W boson propagator:

$$\frac{d\sigma}{d^3k} \propto \frac{d^3k}{(k^2 - M_W^2)^2} \propto \frac{d^3k}{(2EE'(1 - \cos\theta) - M_W^2)^2}$$

Peaks at small θ

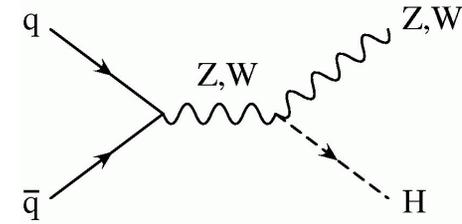
- Outgoing jets are mostly forward and can be tagged



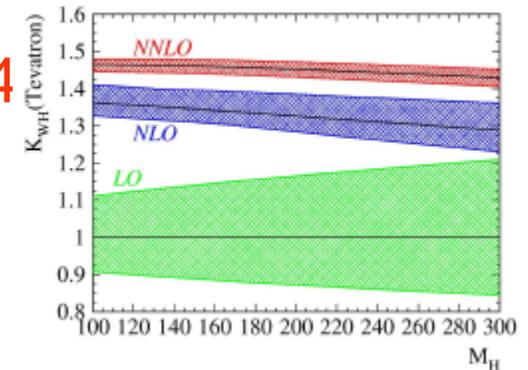
Idea: Look for h decaying to several different channels

Ratio of decay rates will have smaller systematic errors

$W(Z)$ -strahlung



- ✎ $W(Z)$ -strahlung ($q\bar{q}$? Wh , Zh) important at Tevatron
 - ✎ Same couplings as vector boson fusion
 - ✎ Rate proportional to **weak** coupling
 - ✎ Below 130-140 GeV, look for $q\bar{q}$? Vh, h ? $b\bar{b}$
 - ✎ For $M_h > 140$ GeV, look for h ? W^+W^-
- ✎ Theoretically very clean channel
 - ✎ NNLO QCD corrections: $K_{\text{QCD}} \approx 1.3-1.4$
 - ✎ Electroweak corrections known (-5%)
 - ✎ Small scale dependence (3-5%)
 - ✎ Small PDF uncertainties



$t\bar{t}h$ Production

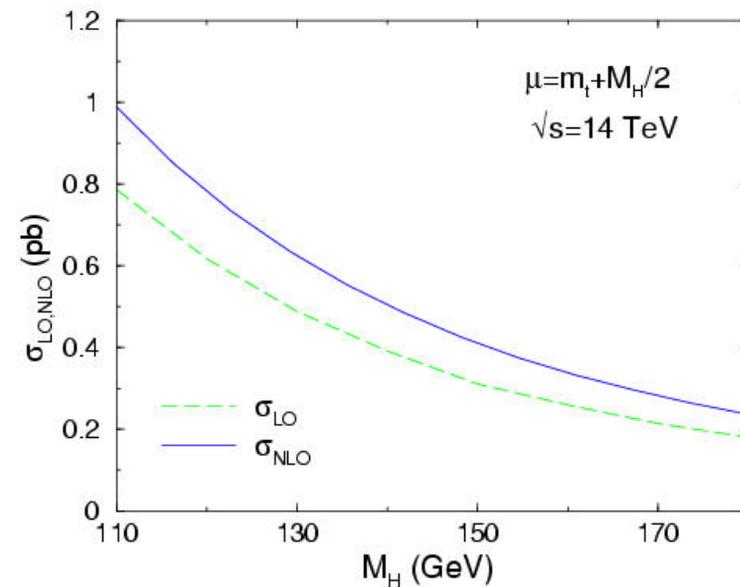
- $t\bar{t}h$ production unique channel to measure top quark Yukawa coupling

–h? $t\bar{t}$ never important

- $b\bar{b}h$ small in SM, but can be enhanced in SUSY models with large $\tan\beta$?



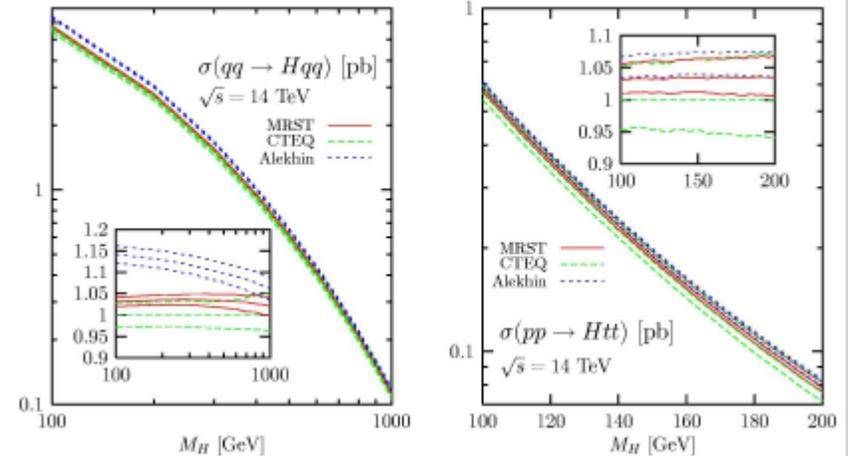
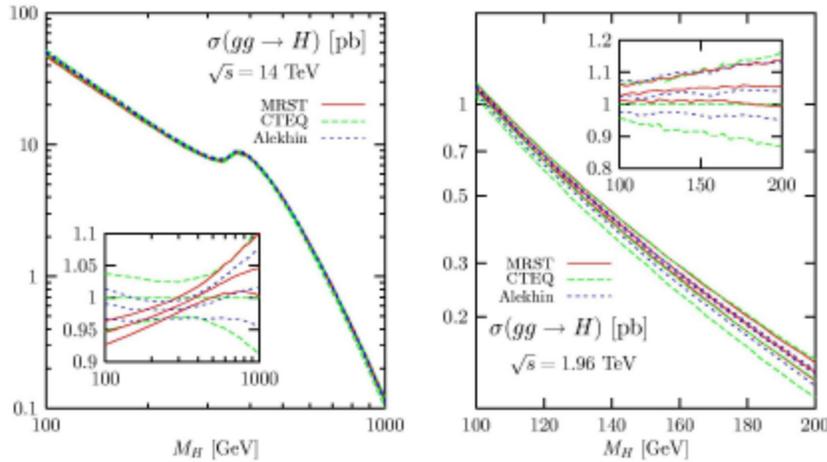
Large QCD effects



Higher order corrections

- ✍ QCD effects can be large
 - ✍ Leading order cross sections have large uncertainties due to:
 - ✍ Renormalization/factorization scale dependence
 - ✍ Uncertainties from parton distribution functions (PDFs)
 - ✍ Important modes have large QCD backgrounds
 - ✍ Often backgrounds only known to leading order
-

PDF uncertainties



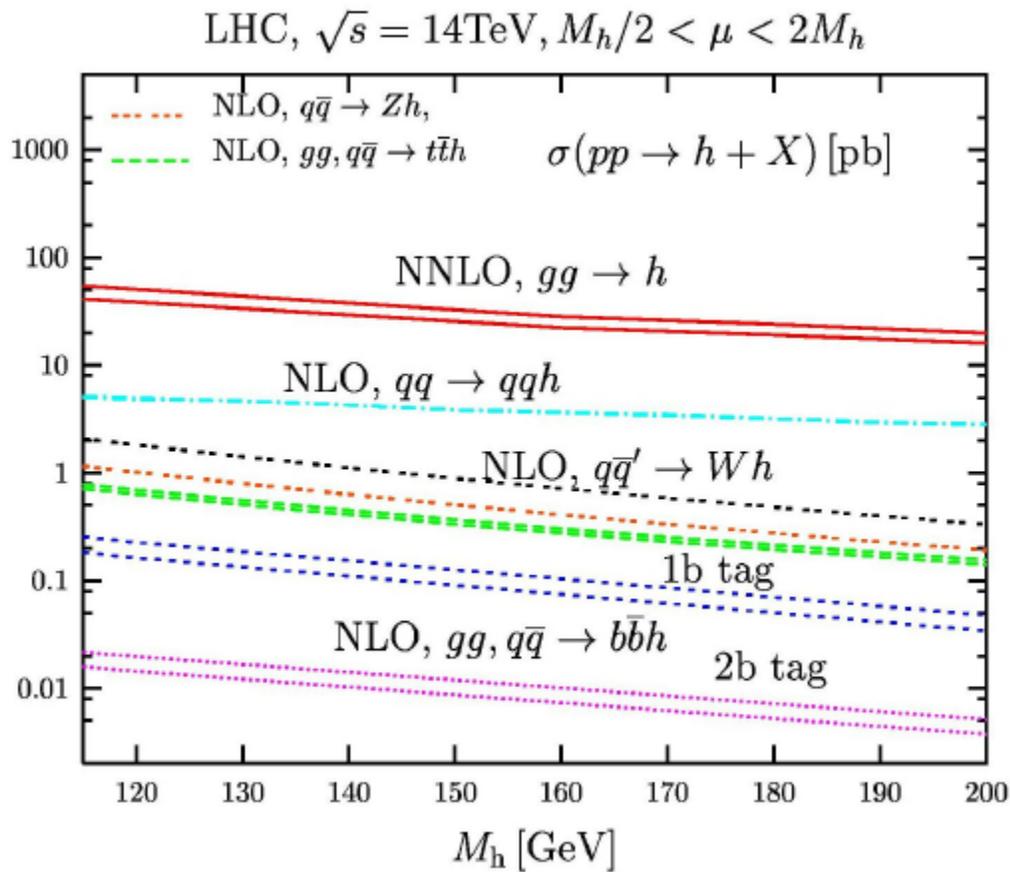
***NLO PDFs with NLO
cross sections!***

CTEQ6m: 40 PDFs for uncertainty studies

<http://user.pa.msu.edu/wkt/cteq/cteq6pdf.html>

Smaller PDF uncertainties in vector boson fusion ($\bar{q}q$ initial channel)

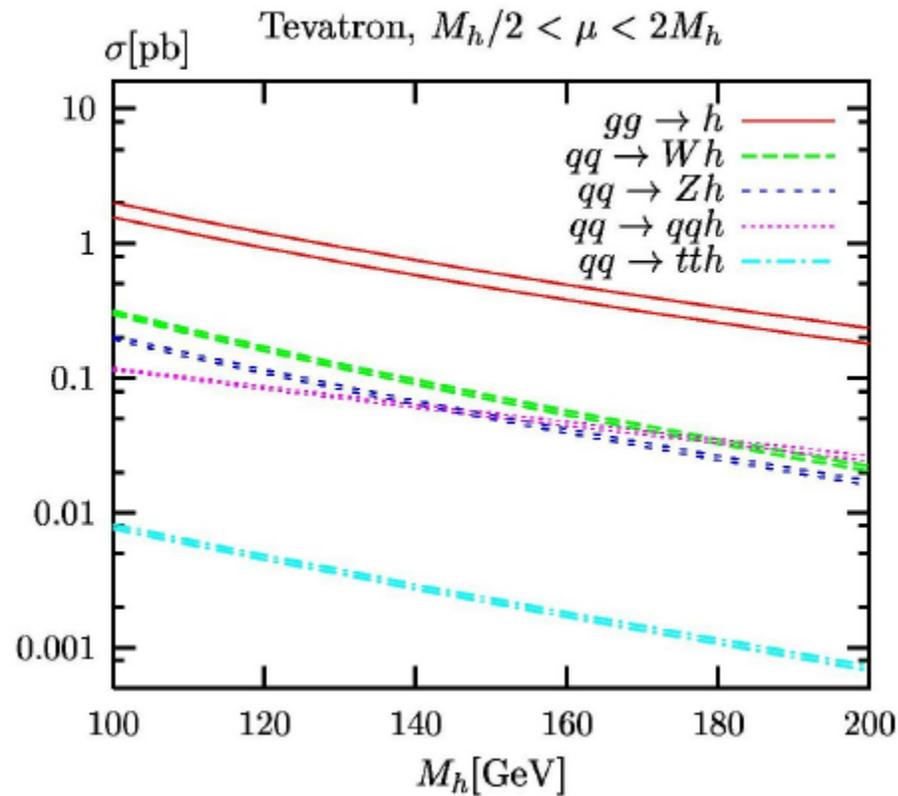
Production mechanisms at LHC



Bands show scale dependence

All important channels
calculated to NLO or NNLO

Comparison of rates at Tevatron

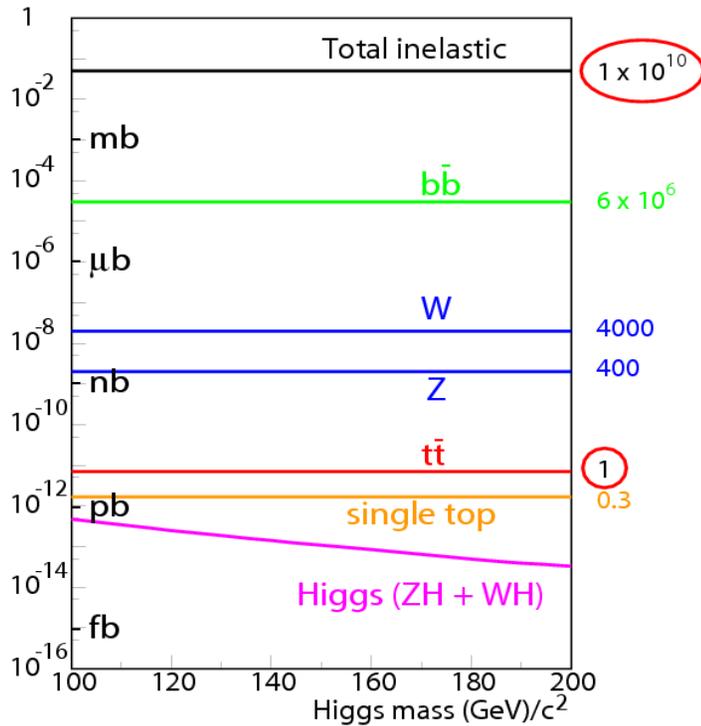


✍ Luminosity goals for Tevatron: 6-8 fb^{-1}

✍ Higgs very, very hard at Tevatron

Higgs at the Tevatron

- Largest rate, $gg \rightarrow h, h \rightarrow b\bar{b}$, is overwhelmed by background



? ($gg \rightarrow h$)? $1 \text{ pb} \ll ? (b\bar{b})$

Higgs at the Tevatron

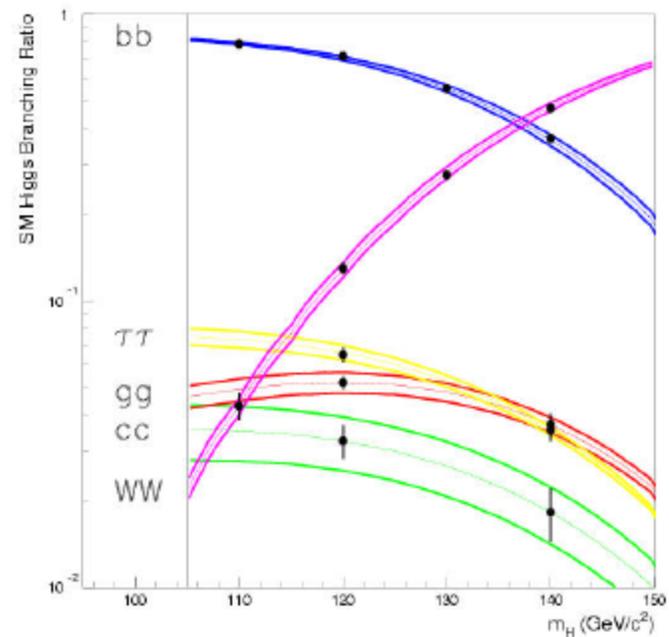
- ✍ Wh, Zh production important for $M_h < 140$ GeV, h? bb
- ✍ Background from $Wb\bar{b}$, $Zb\bar{b}$
- ✍ One of the few examples where both signal and background known to NLO

Wh, Zh and background in MCFM Monte Carlo to NLO

<http://mcfm.fnal.gov>

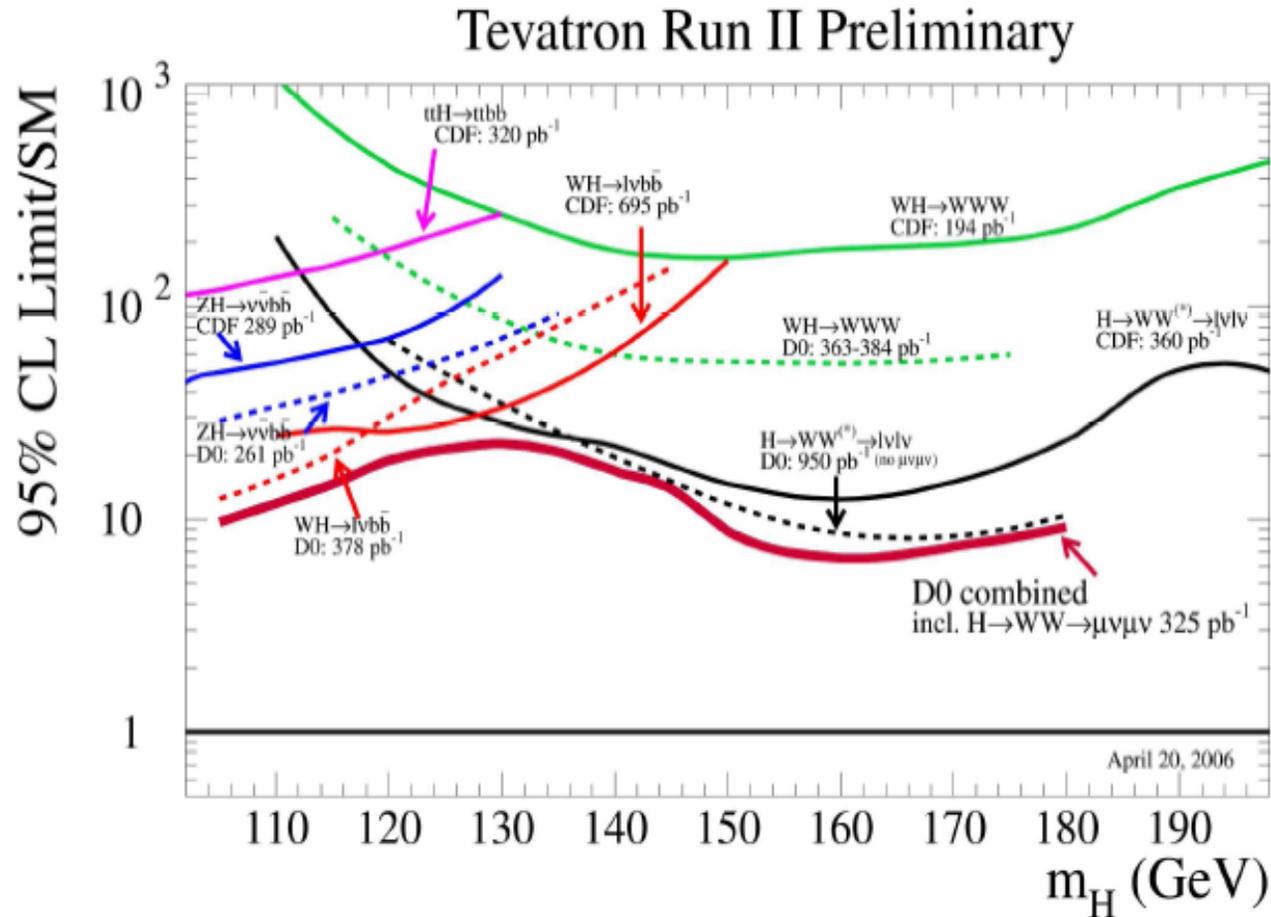
Search channels at Tevatron

- ✍ For heavier Higgs, look for $h? W^+W^-$
- ✍ Searches for $gg? h? W^+W^-$ (dileptons)
- ✍ And $Wh? W^?W^+W^-$ (2 and 3 leptons)

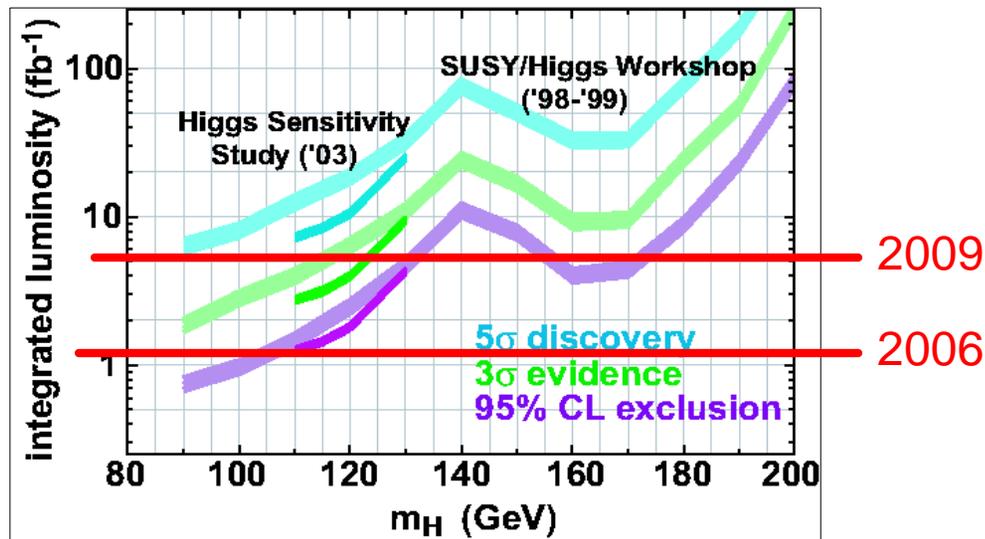


Requiring leptons reduces backgrounds

Tevatron Higgs Searches



Can the Tevatron discover the Higgs?



This relies on statistical combination of multiple weak channels

Search Channels at the LHC

gg? h? $b\bar{b}$ has huge QCD bkd:
Must use rare decay modes of h

✎ **gg? h? ??**

✎ Small BR ($10^{-3} - 10^{-4}$)

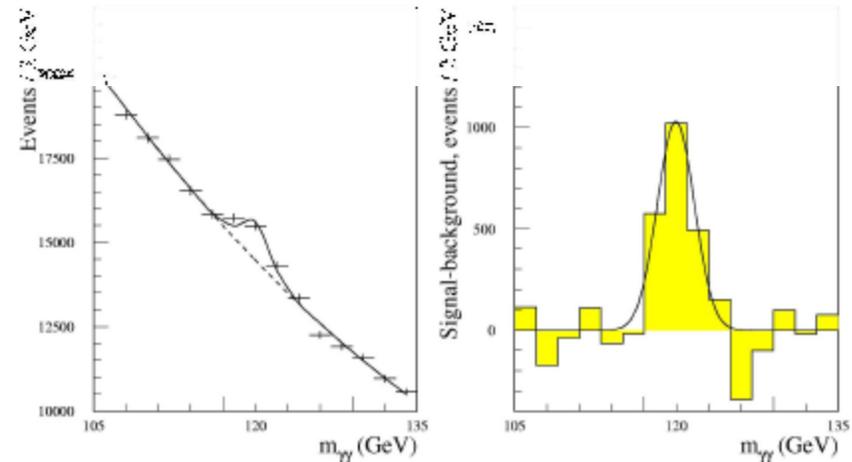
✎ Only measurable for $M_h < 140$ GeV

✎ Largest Background: QCD continuum production of $\gamma\gamma$

✎ Also from γ -jet production, with jet faking γ , or fragmenting to γ^0

✎ Fit background from sidebands of data

$M_h = 120$ GeV; $L = 100$ fb $^{-1}$



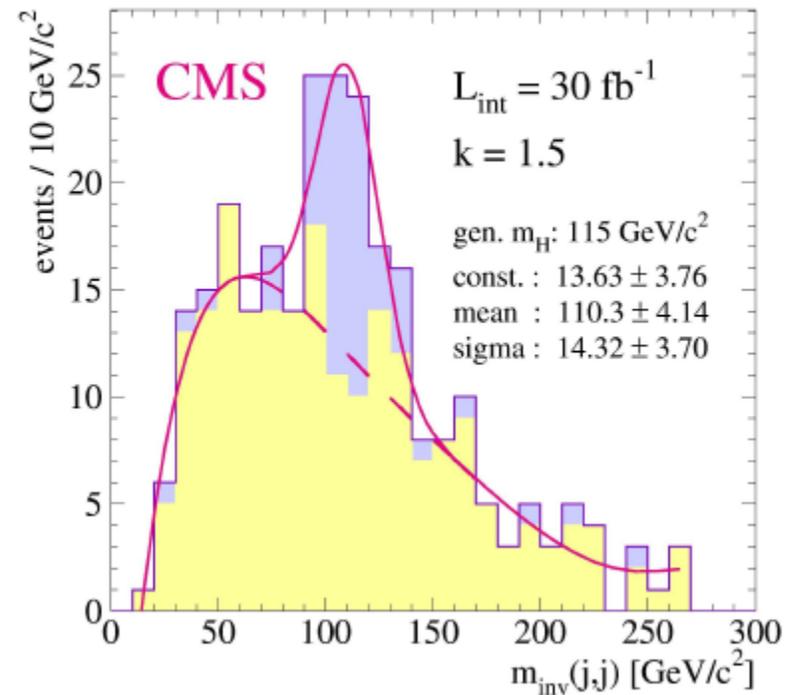
$S/\gamma B = 2.8$ to 4.3 ?

• Gives 1% mass measurement

$t\bar{t}h$ at the LHC

- ✍ $gg?$ $t\bar{t}h?$ $t\bar{t}b\bar{b}$
- ✍ Spectacular signal
- ✍ $t?$ Wb
- ✍ Look for 4 b jets, 2 jets, 1 lepton

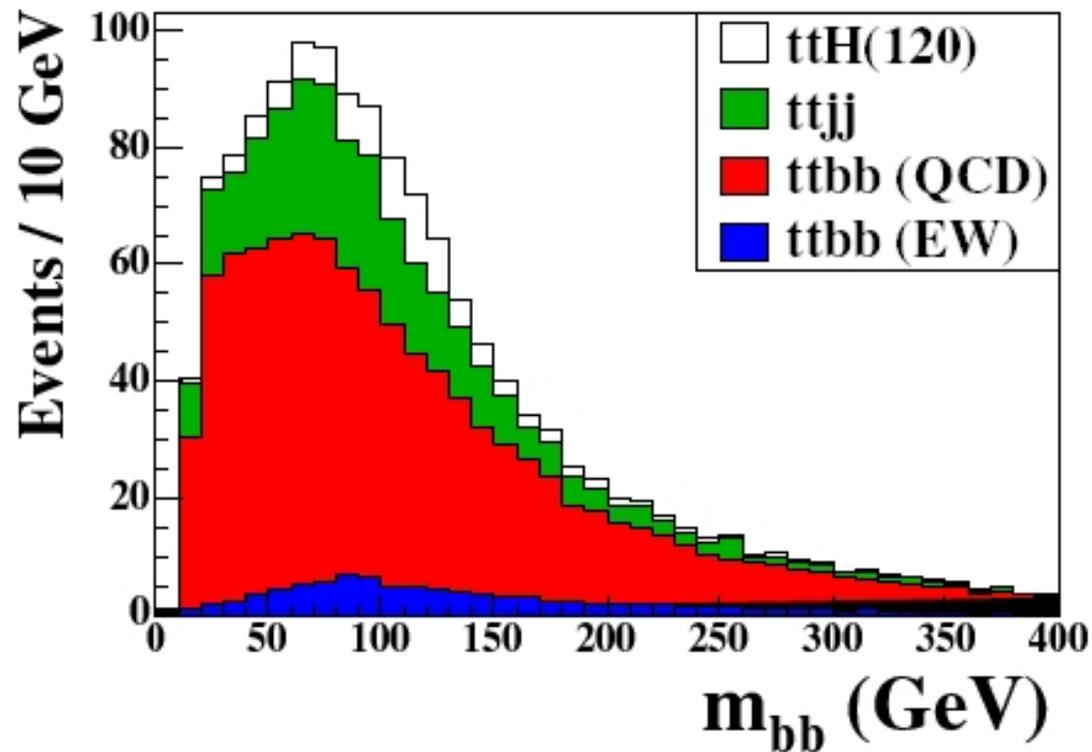
Unique way to measure top quark Yukawa coupling



Early studies looked promising

BUT...Large QCD background to $t\bar{t}h$

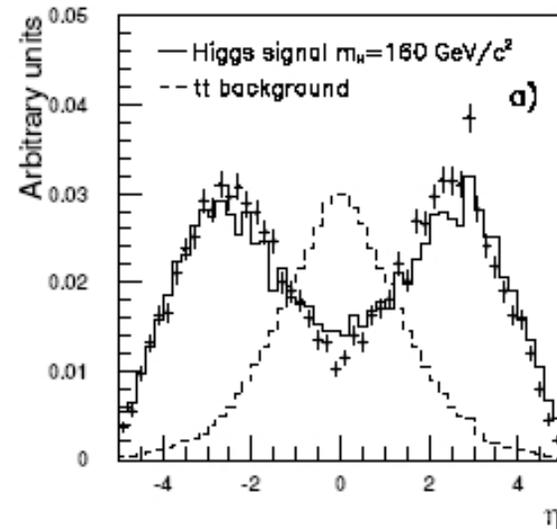
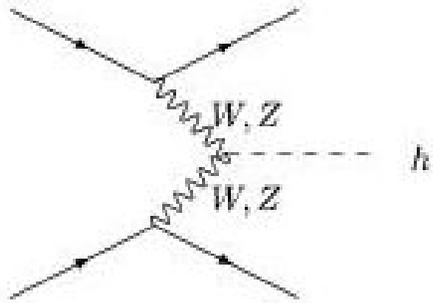
Current $t\bar{t}H, H \rightarrow b\bar{b}$ outlook: (30 fb^{-1})



$S/B=1/6$ for $M_h=120 \text{ GeV}$

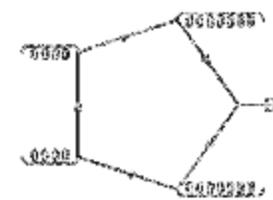
Vector Boson Fusion

- ✎ Outgoing jets are mostly forward and can be tagged
- ✎ Vector boson fusion and QCD background look different

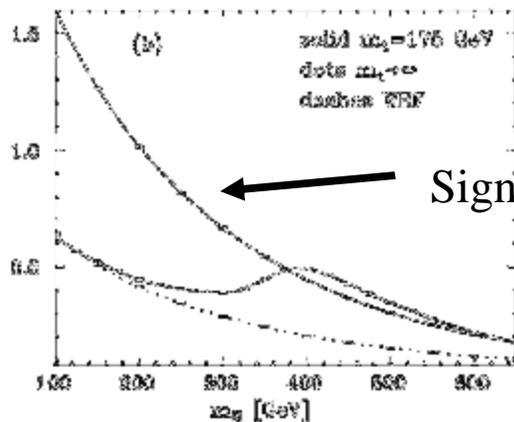


Vector Boson Fusion

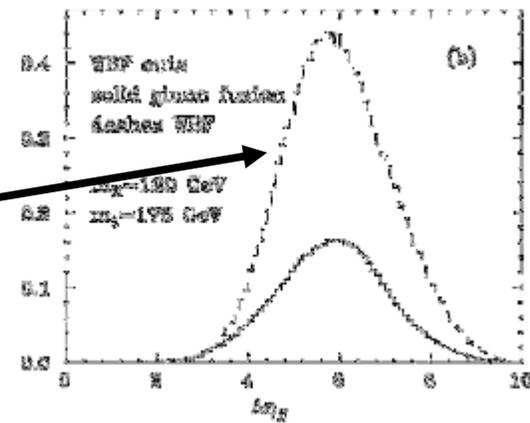
- Identify signal with forward jet tagging and central jet veto
- Large Higgs + 2 jet background from $gg \rightarrow ggh$
- Kinematic cuts effective at identifying signal



Higgs + 2 jet Production



Rapidity between outgoing jets



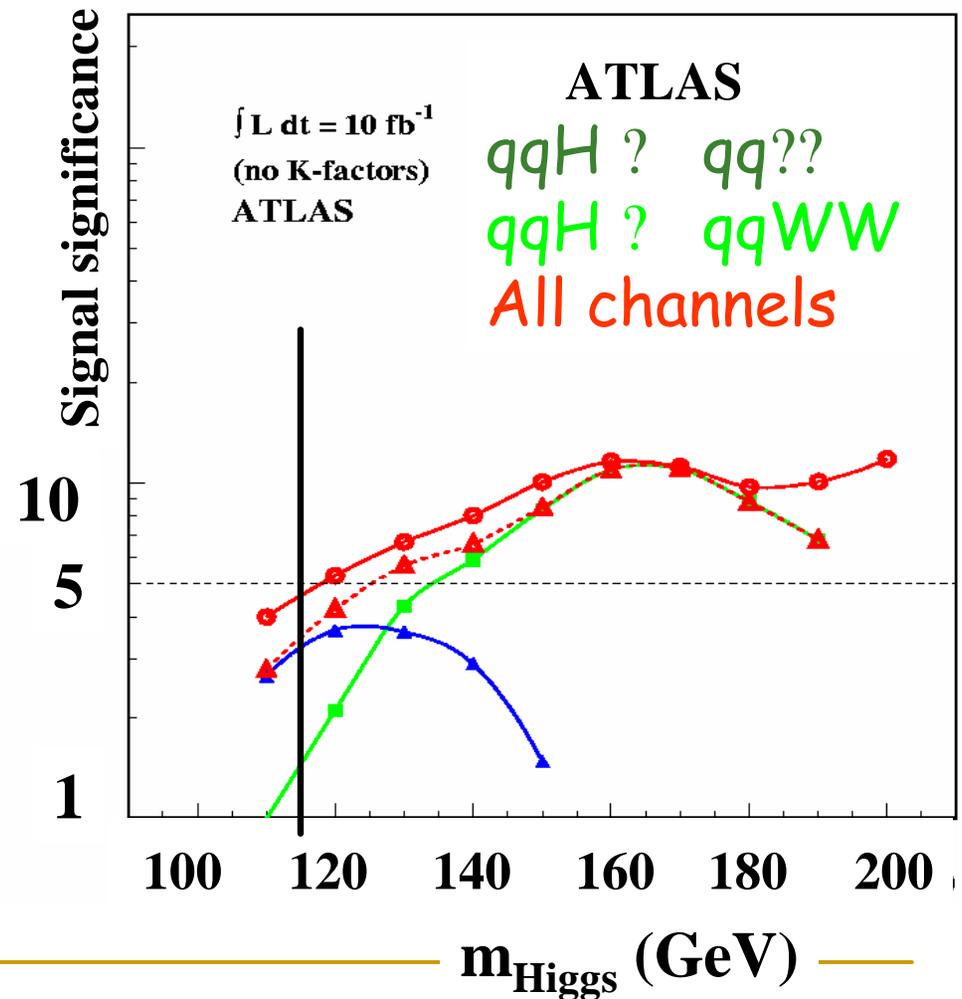
?

Vector Boson Fusion for light Higgs

✍ For $M_h = 115$ GeV
combined significance ~ 5 ?

*Vector boson fusion
effective for measuring
Higgs couplings*

✍ Proportional to g_{WWH} and g_{ZZH}
✍ Often assume they are in SU(2)
ratio: $g_{WWH}/g_{ZZH} = \cos^2 \theta_w$



Vector Boson Fusion for Heavy Higgs

200 GeV < M_h < 600 GeV:

- **discovery in h ? ZZ ? l+l+l+l**
- **Background smaller than signal**
- **Higgs width larger than experimental resolution (M_h > 300 GeV)**
- **confirmation in h ? ZZ ? l+l+jj channel**

M_h > 600 GeV:

4 lepton channel statistically limited

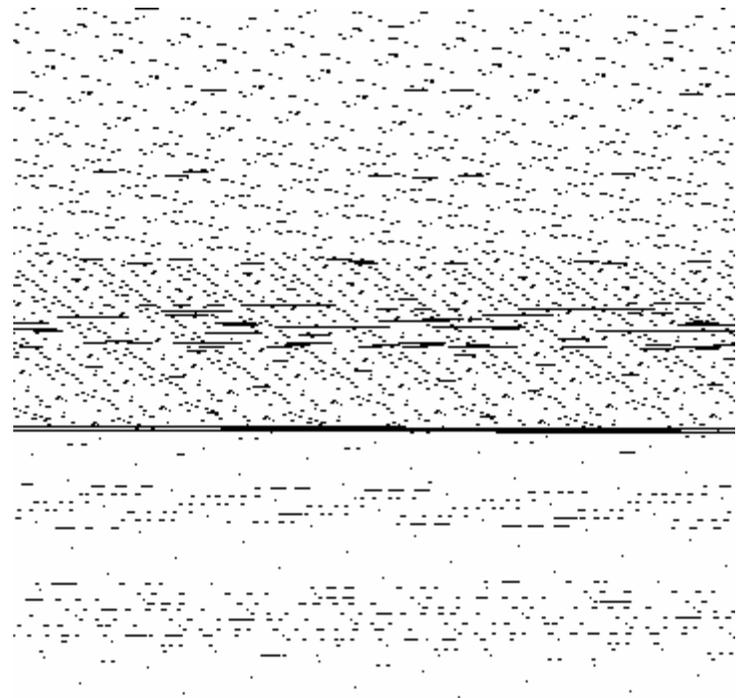
h ? ZZ ? l+l+??

h ? ZZ ? l+l+jj, h ? WW ? l?jj

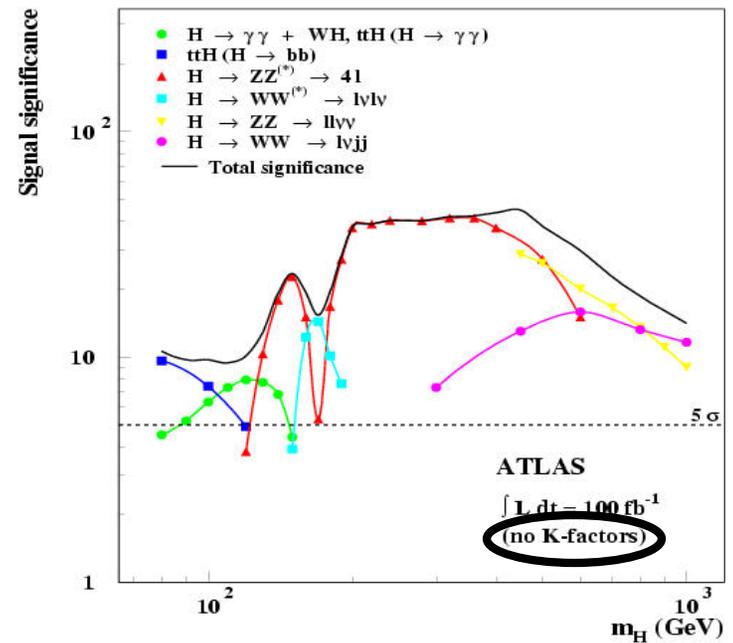
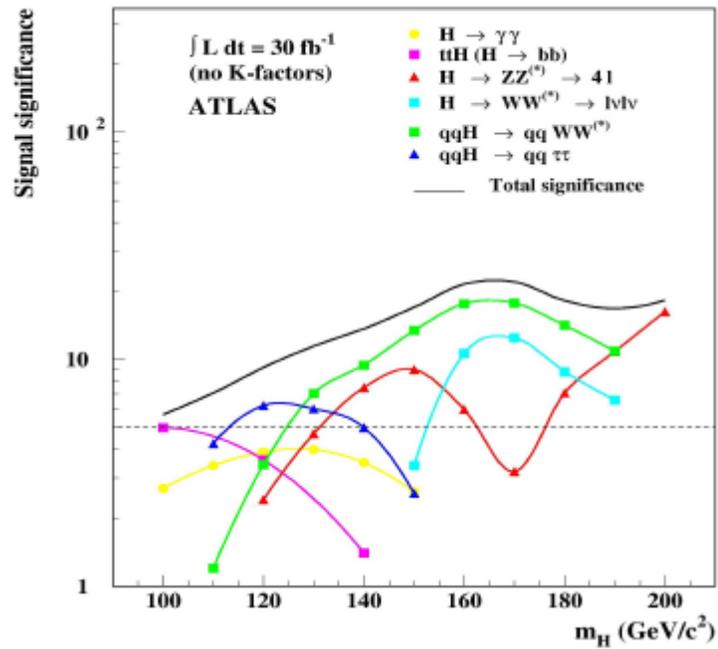
-150 times larger BR than 4l channel

Gold-plated

h ? ZZ ? l+l+l+l



If there is a light SM Higgs, we'll find it at the LHC



No holes in M_h coverage

If we find a “Higgs-like” object, what then?

 **We need to:**

-  Measure Higgs couplings to fermions & gauge bosons
-  Measure Higgs spin/parity
-  Reconstruct Higgs potential
-  Is it the SM Higgs?

 **Reminder: Many models have other signatures:**

-  New gauge bosons (little Higgs)
 -  Other new resonances (Extra D)
 -  Scalar triplets (little Higgs, NMSSM)
 -  Colored scalars (MSSM)
 -  etc
-

Is it a Higgs?

- How do we know what we've found?
- Measure couplings to fermions & gauge bosons

$$\frac{?(h ? b\bar{b})}{?(h ? ???)} ? 3 \frac{m_b^2}{m_\gamma^2}$$

- Measure spin/parity

$$J^{PC} ? 0^{??}$$

- Measure self interactions

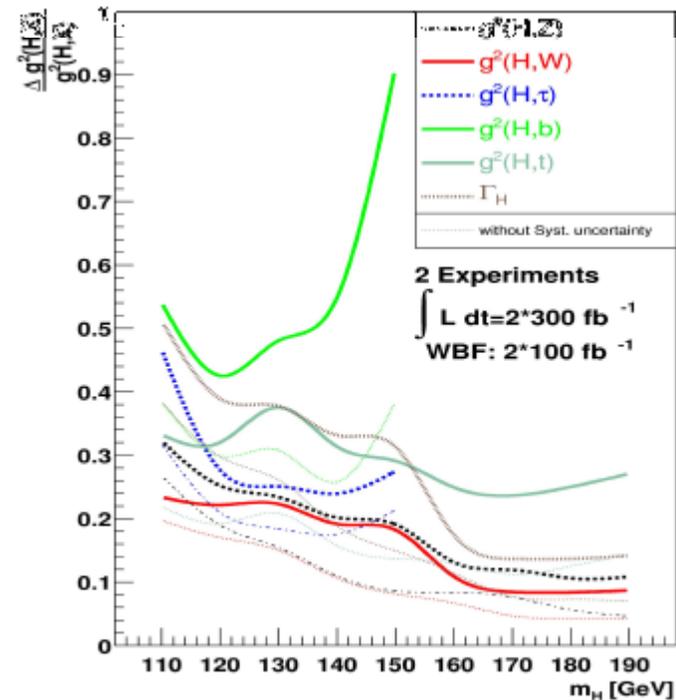
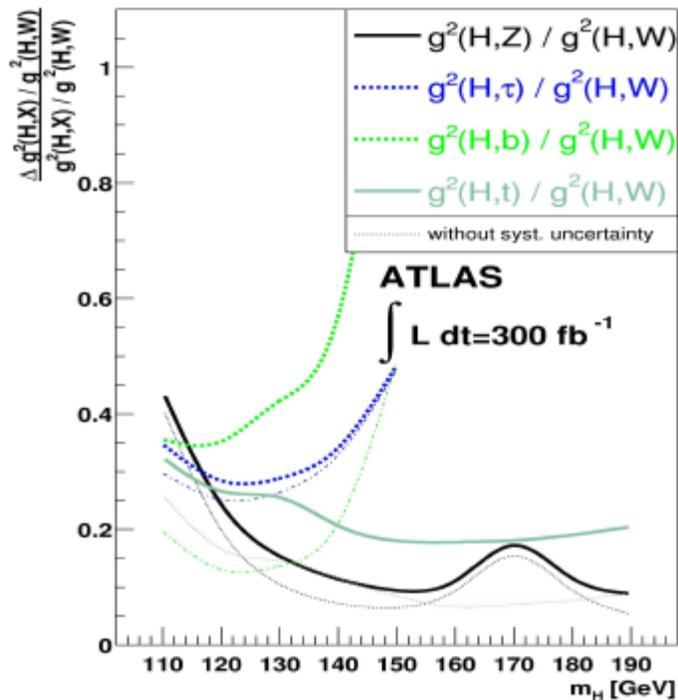
$$V ? \frac{M_h^2}{2} h^2 ? \frac{M_h^2}{2v} h^3 ? \frac{M_h^2}{8v^2} h^4$$

Very hard at
hadron collider

Absolute measurements of Higgs couplings

✎ Ratios of couplings more precisely measured than absolute couplings

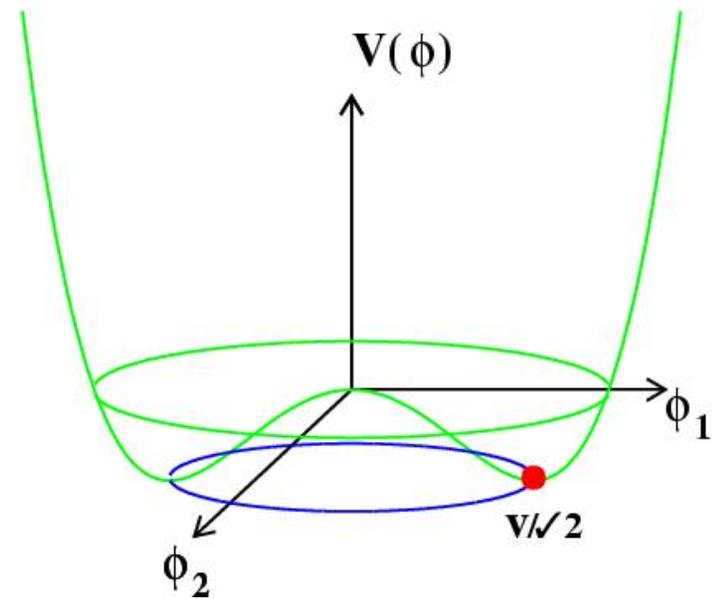
✎ 10-40% measurements of most couplings



Can we reconstruct the Higgs potential?

$$V \sim \frac{M_h^2}{2} h^2 + \lambda_3 v h^3 + \frac{\lambda_4}{4} h^4$$

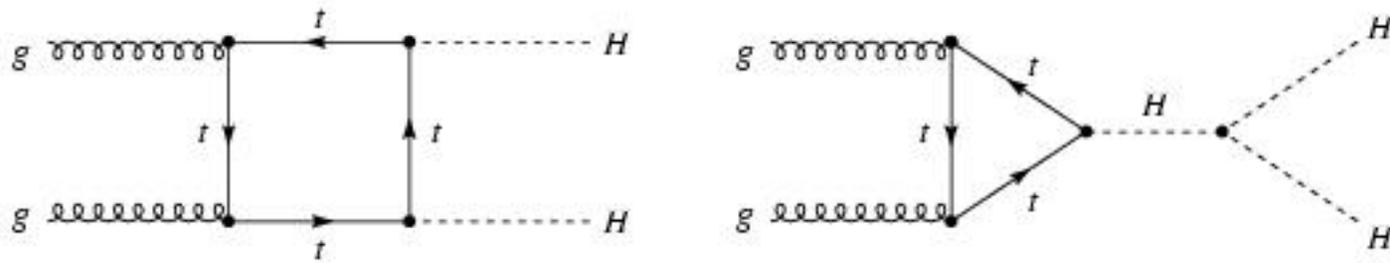
$$SM : \lambda_3 \sim \lambda_4 \sim \frac{M_h^2}{2v^2}$$



✍ Fundamental test of model!

✍ We have no idea how to measure λ_4

Reconstructing the Higgs potential



- ✍ α_3 requires 2 Higgs production
- ✍ $M_h < 140$ GeV, $h \rightarrow b\bar{b}b\bar{b}$
- ✍ Overwhelming QCD background
- ✍ Easier at higher M_h

Can determine whether $\alpha_3=0$ at 95% cl with
 300 fb^{-1} for $150 < M_h < 200$ GeV

Higgs measurements test model!

✍ Supersymmetric models are our favorite comparison

✍ SUSY Higgs sector

✍ At least 2 Higgs doublets

✍ SM masses from

$$L \ ? \ ? \ g_d \bar{Q}_L \ ? \ d_R \ ? \ g_u \bar{Q}_L \ ? \ ^c u_R \ ? \ h.c.$$

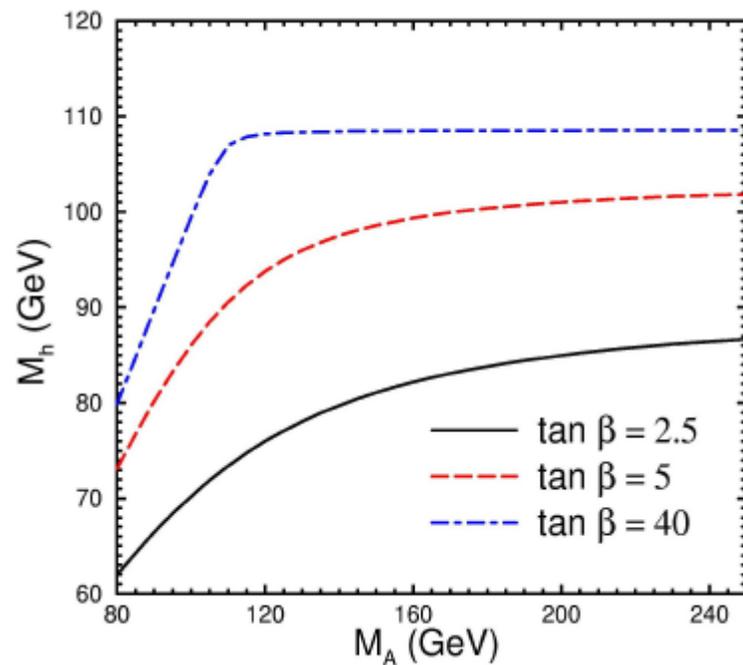
✍ $\ ? \ ^c$ term not allowed in SUSY models: Need second Higgs doublet with opposite hypercharge

✍ 5 physical Higgs: h^0, H^0, A^0, H^\pm

SUSY Higgs

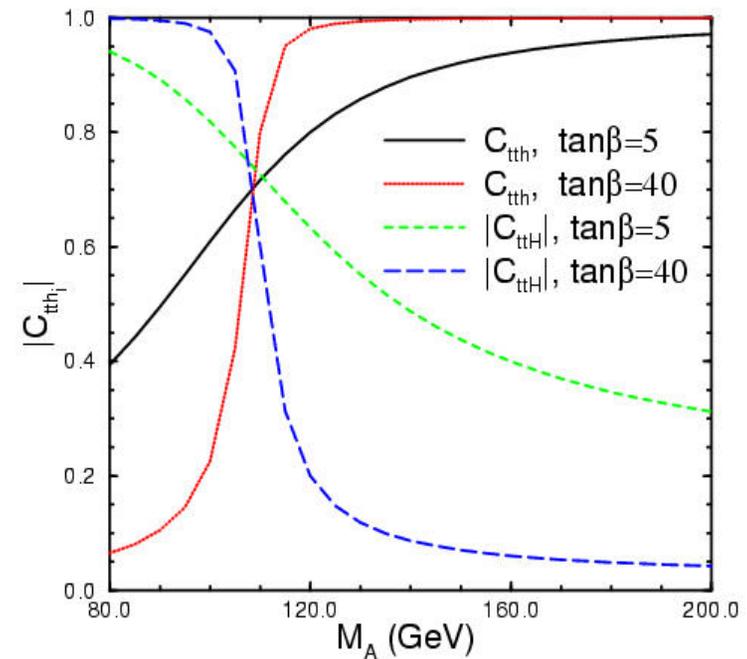
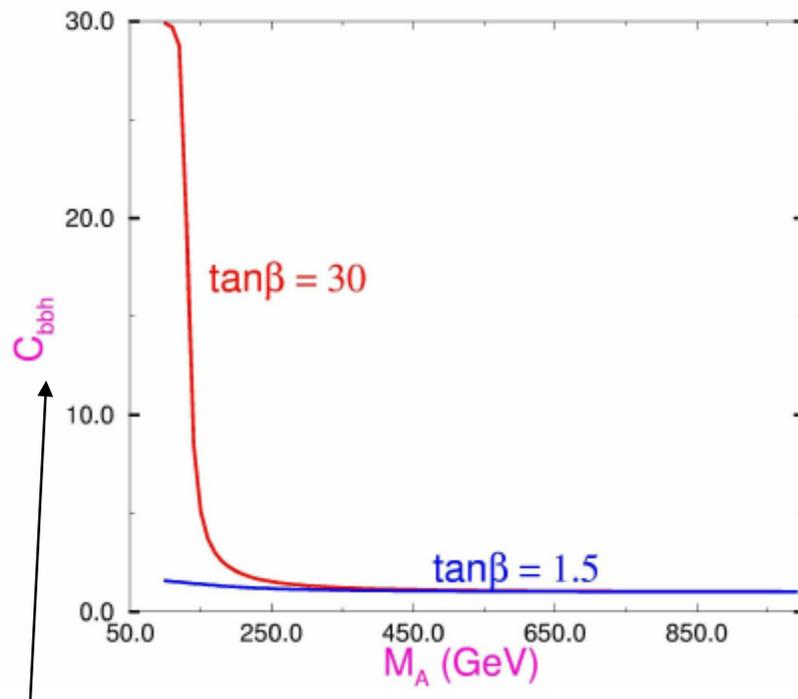
- ✍ General 2 Higgs doublet potential has 6 couplings and a phase
 - ✍ SUSY Higgs potential has only 2 couplings
 - ✍ Take these to be **M_A and $\tan\beta$** ?
 - ✍ At tree level Higgs couplings, neutral and charged Higgs masses are predicted
 - ✍ Lightest Higgs mass has upper limit
-

Upper Limit on Higgs Mass in SUSY Models



Can tune parameters,
but always have
upper limit below
 $M_h \approx 130$ GeV

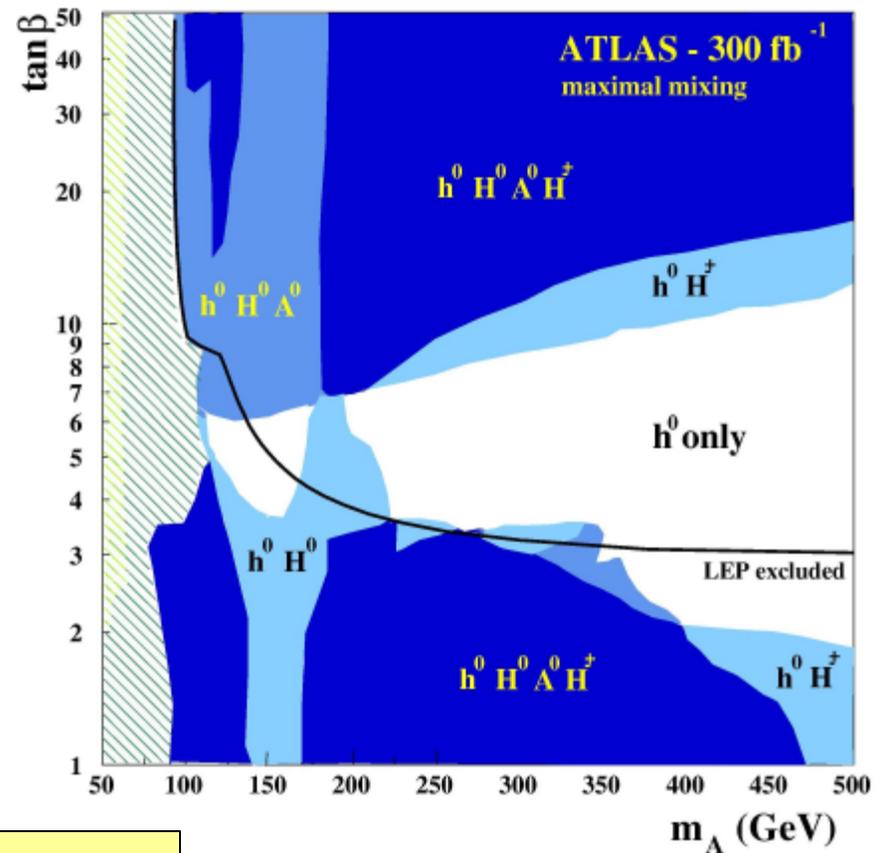
Higgs Couplings very different from SM in SUSY Models



Ratio of h coupling to b's in SUSY model to that of SM

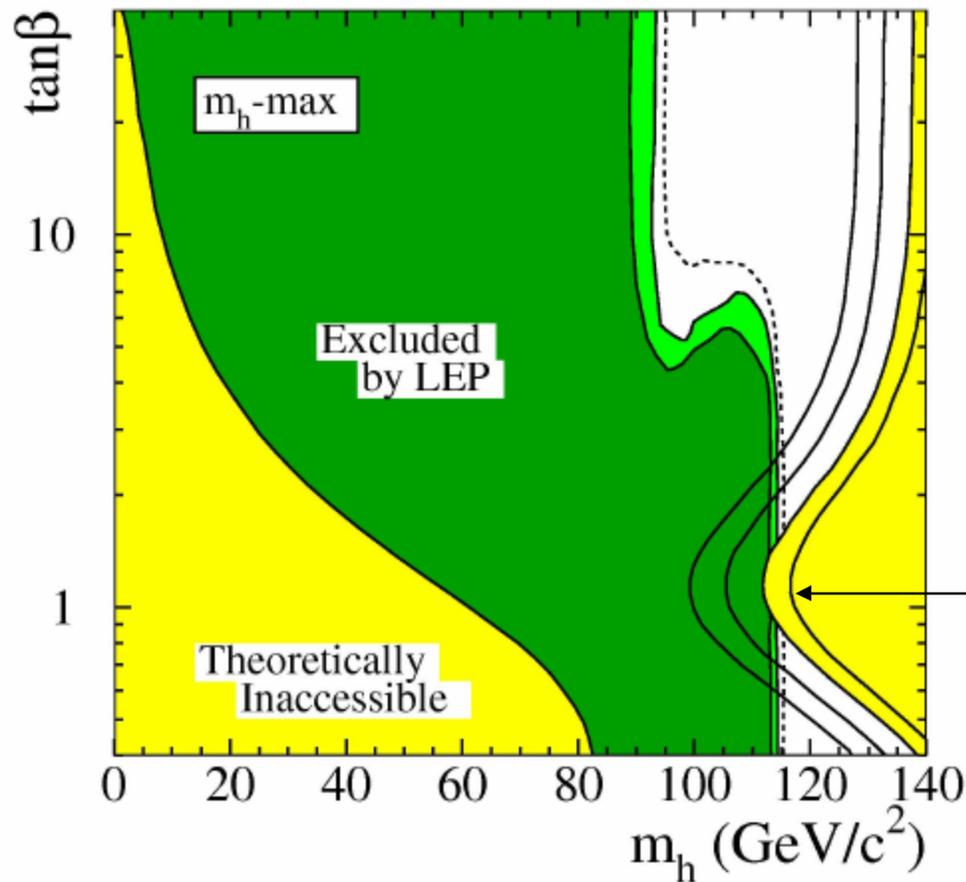
MSSM discovery

- ✎ For large fraction of M_A - $\tan\beta$ space, more than one Higgs boson is observable
- ✎ For $M_A \gg m_h$, MSSM becomes SM-like
- ✎ Plot shows regions where Higgs particles can be observed with $> 5\sigma$



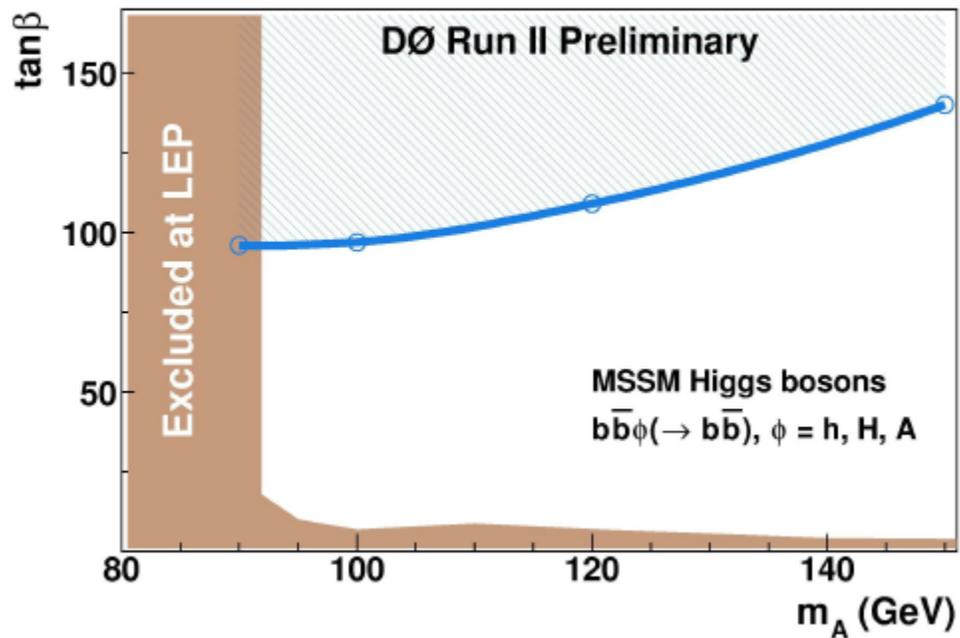
Need to observe multiple Higgs bosons and measure their couplings

Limits on SUSY Higgs from LEP



$M_t = 169.3, 174.3,$
 $179.3, 183$ GeV

New Discovery Channels in SUSY



Conclusion

The Higgs boson is the final missing link in the SM
