

Experiencing Discovery

Hadron Collider Physics Summer School

Fermilab

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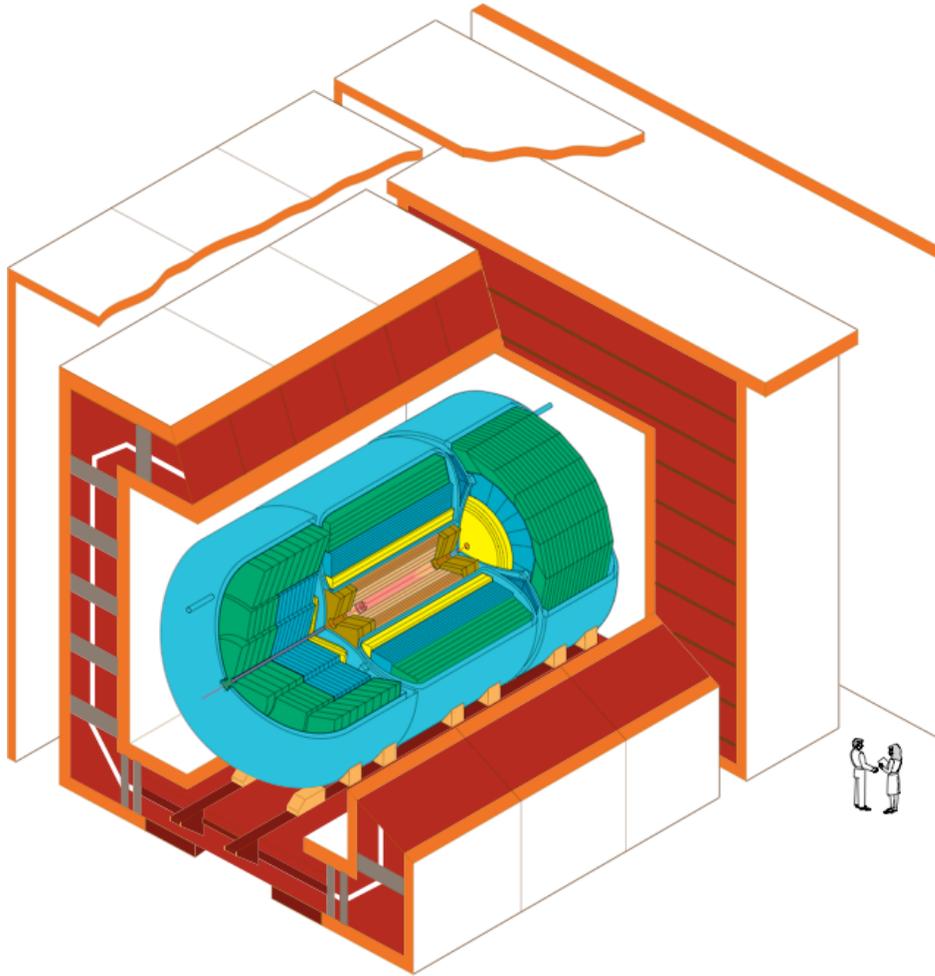
Who was I?

- Two very exciting things happened in Fall of 1992:
 1. DØ began collecting physics-quality data
 2. I started grad school at UC Berkeley
- From then until Summer of 1993, the following happened
 - Teams of physicists began the hunt for the top quark
 - I tried to solve *@#!@ Jackson problems
- By Summer 1993 I had decided to go into experimental HEP for my thesis
- By the time I was done with classes and able to go into research full-time it was spring of 1994
 - hints of top were already appearing by then
- So I'm reporting on work done almost entirely by others
 - I'll mention my role when it comes up...

Expectation of Top

- Even before the Tevatron turned on, the top quark was widely anticipated
 - the b quark had been shown to have isospin $-1/2$ -- therefore it had to have a partner!
- Even the mass of the top was constrained in the SM context
 - but those limits had been creeping up, “staying ahead” of the experimentally-excluded value
- So the situation was analogous to the Higgs boson today
 - if anything, I would argue that the top was even more of a “sure thing”

The Run I DØ Detector



Unique features:

- no B field in central region
- Finely segmented calorimeter covering out to $|\eta| \sim 4$
- Toroidal iron magnet for muon spectrometer
- Two accelerators
 - that's not a plus..

Comparing the Detectors

- The Tevatron was clearly the only place where the top quark might show up in the 1990's
- So, which of the two detectors would see it first?
- Though broadly similar, they did have significant differences
 - the advantages were as follows:

	CDF	DØ
Electron ID	✓	
Electron Energy		✓
Muon ID		✓
Muon Momentum	✓	
Jet Energy		✓
Missing E_T		✓
b Jet Tagging	✓ ✓	

The Backgrounds

- For every set of final-state objects produced in a $t\bar{t}$ event, there are other SM processes that produce the same objects. These backgrounds are:
- All-hadronic mode: QCD multijet production (HUGE)
- Lepton+jets mode: $W + \geq 4$ jet production
- Dilepton modes: Z, WW or $WZ + \geq 2$ jet production
- Note that the jets in the background tend to arise from gluon radiation
 - rate calculations were leading-order \rightarrow large uncertainties for high jet multiplicities
 - better to calculate background rates from data whenever possible

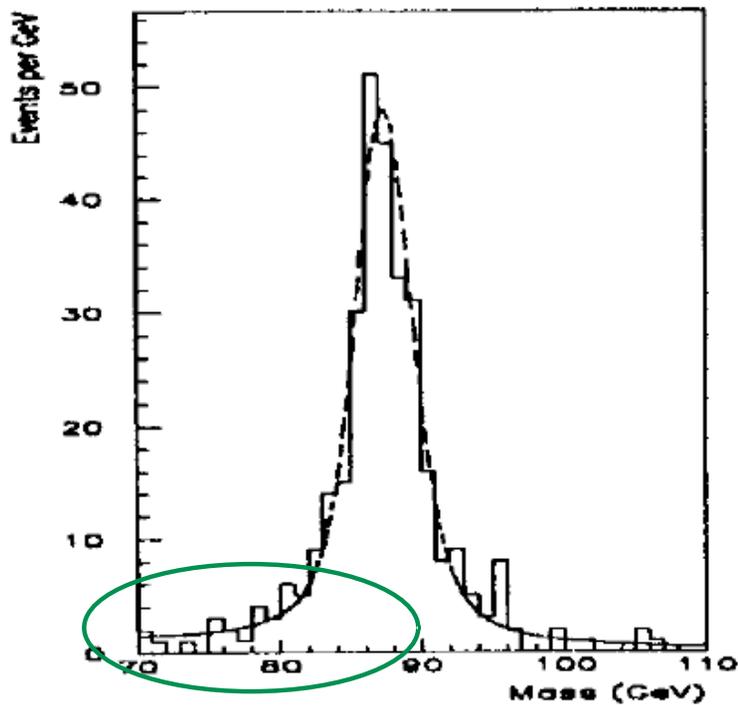
Detector Backgrounds

- In addition to the physics backgrounds on the previous slide, there are also detector backgrounds
 - a jet that appears to be a lepton
 - missing E_T coming from mismeasured jet or muon energies
- These can be studied in data directly
 - i.e. by looking at events where a jet satisfies some, but not all of the lepton ID requirements
- For the signal and physics backgrounds one needs to understand:
 - the kinematic distributions of leptons, quarks and gluons produced (theoretical model)
 - the response of DØ to those objects (calibration)

Calibrating (Leptons)

- Based on known resonance masses
 - $Z, J/\psi, \pi^0$

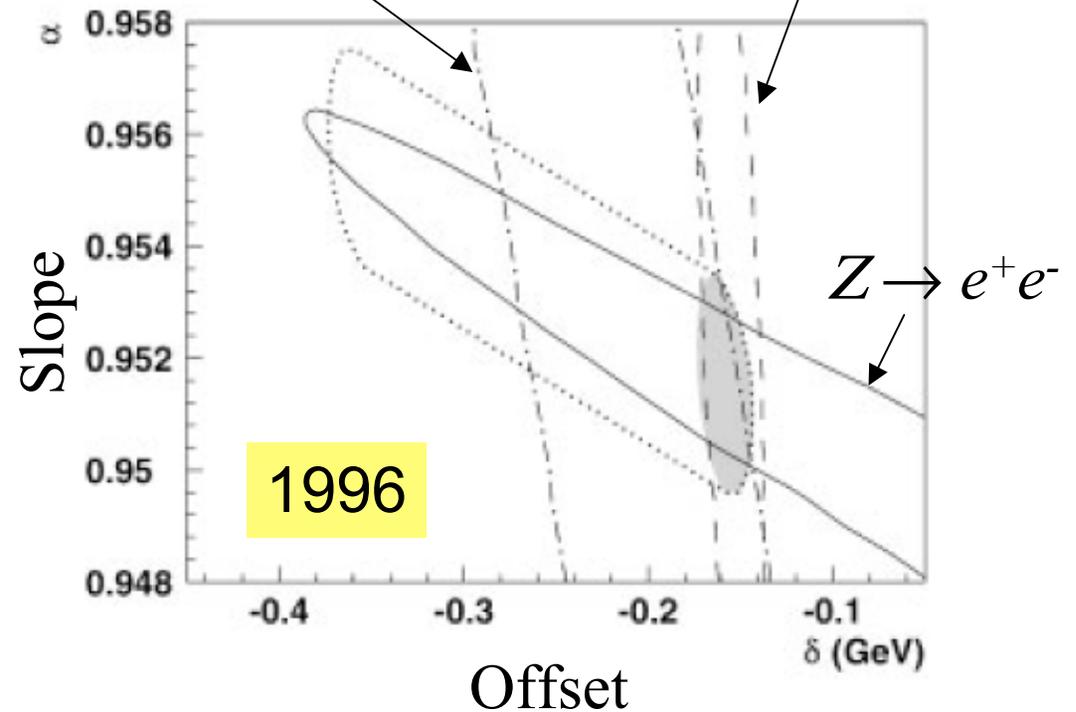
$Z \rightarrow e^+e^-$



No tail from brems

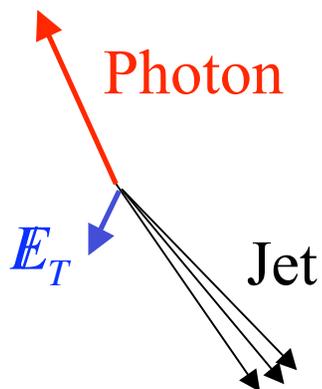
$J/\psi \rightarrow e^+e^-$

$\pi^0 \rightarrow e^+e^-$



Calibrating (Jets)

- No resonances to help out here
 - later in Run I tried to observe $W \rightarrow qq$, but no luck
- So, tie the jet calibration to the EM energy scale
 - use events with a photon recoiling against a jet
 - if calibration were perfect, would have $E_T^\gamma = E_T^{\text{jet}}$, $\cancel{E}_T = 0$
 - in real life, one sees something like:



- Hadronic response R given by:

$$R = 1 + \frac{\cancel{E}_T \cdot E_T^\gamma}{E_T^{\gamma 2}}$$

- But R is only part of the story:

Energy of particles belonging to jet within jet cone. NOT the energy of the parton that produced the jet

$$E_{\text{true}} = \frac{E_{\text{meas}} - U}{(1 - C)R}$$

Energy from electronics noise, radioactive decay, random particles

Energy leaking out of (or into) the jet cone

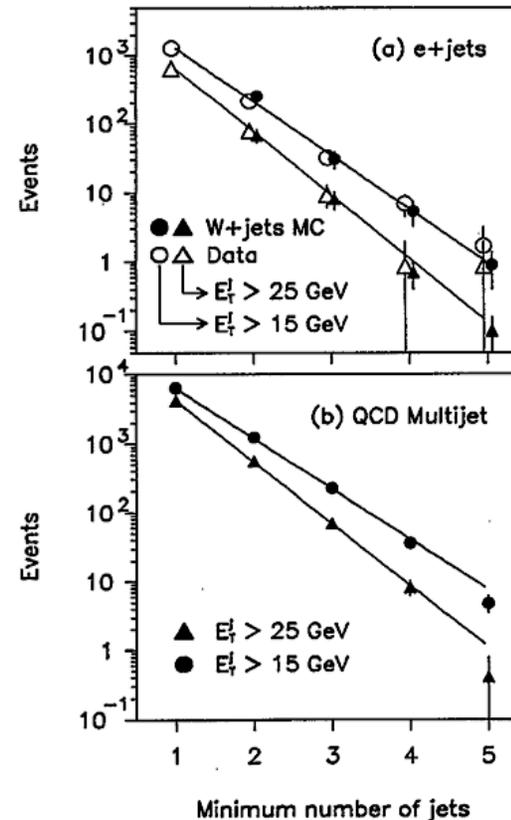
- Hadronic calibration was (and is!) one of the more complex facets of top quark measurements

Estimating the Background From Data

- “Berends scaling”

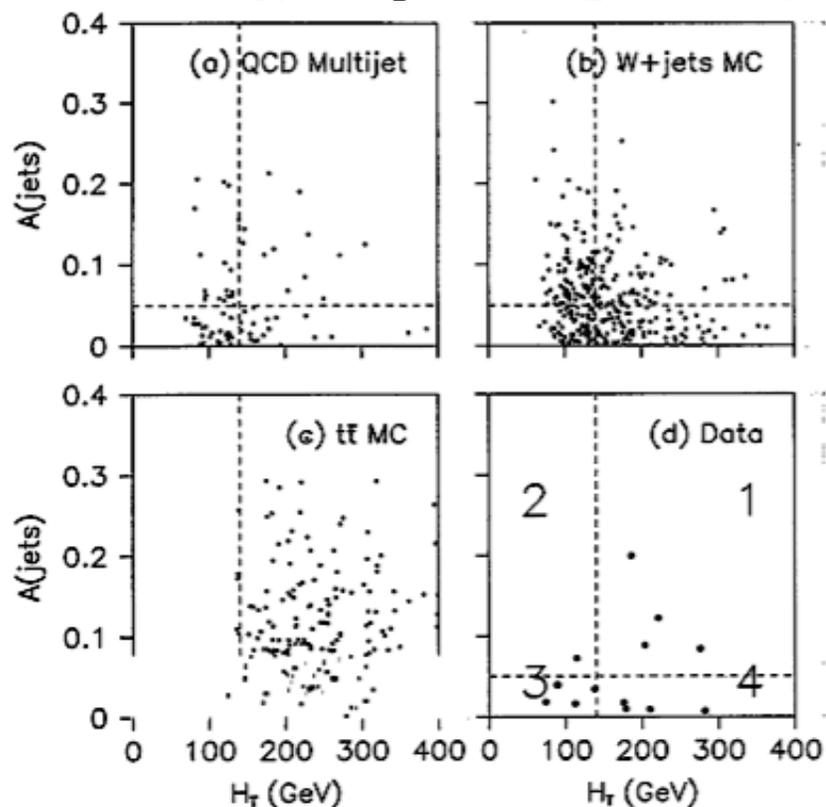
$$\text{In theory: } \frac{W + (n-1) \text{ jets}}{W + n \text{ jets}} = \frac{W + (n-2) \text{ jets}}{W + (n-1) \text{ jets}} \equiv \alpha$$

- Data agrees well with this expectation
 - for both of the major background sources in l+jets channel



Using Kinematics

- One can also look at how events are distributed in total jet energy (H_T) and aplanarity:



- *Both* methods rely on input from theory
 - Scaling rule in 1st case
 - Kinematic distributions for background in 2nd case*
- But these are two different aspects of the theory
- Results should be consistent between the two methods
 - and they were!

- Average of the two was taken as the background estimate

*EWV contributed to studies here

The First Look at DØ

- DØ's first in-depth search for the top quark was done on about 15pb^{-1} of data from 1992 and 1993
 - we get more in a good week nowadays!
- From CDF's previous search, we knew that the mass was somewhere above 91 GeV
 - means W can't decay to top, and W in top decay is on-shell
- But where? Needed a selection that was efficient across the entire mass range
 - i.e. loose kinematic cuts
 - used aplanarity to distinguish “spherical” $t\bar{t}$ events from “jet-like” backgrounds

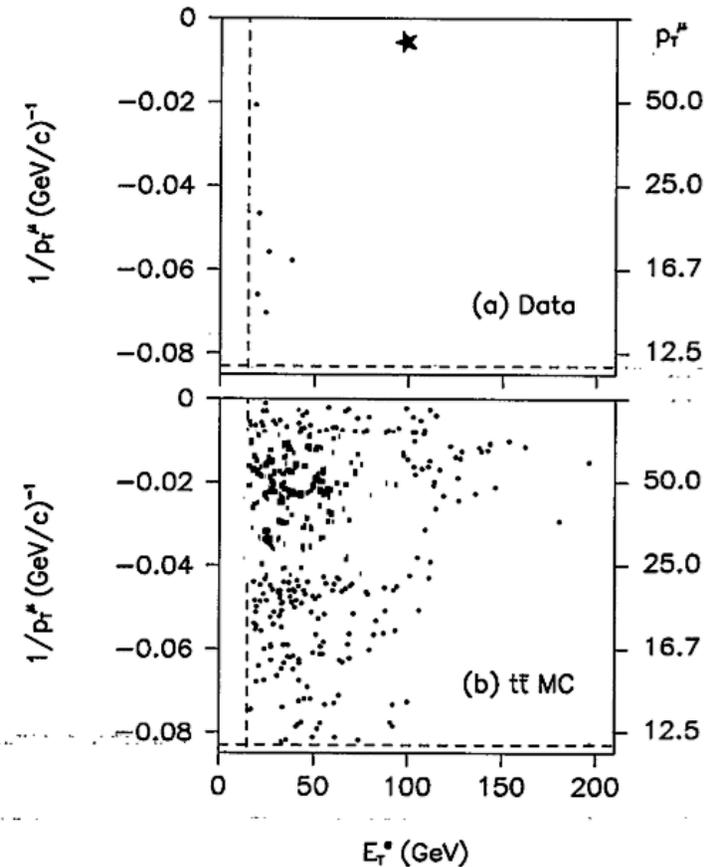
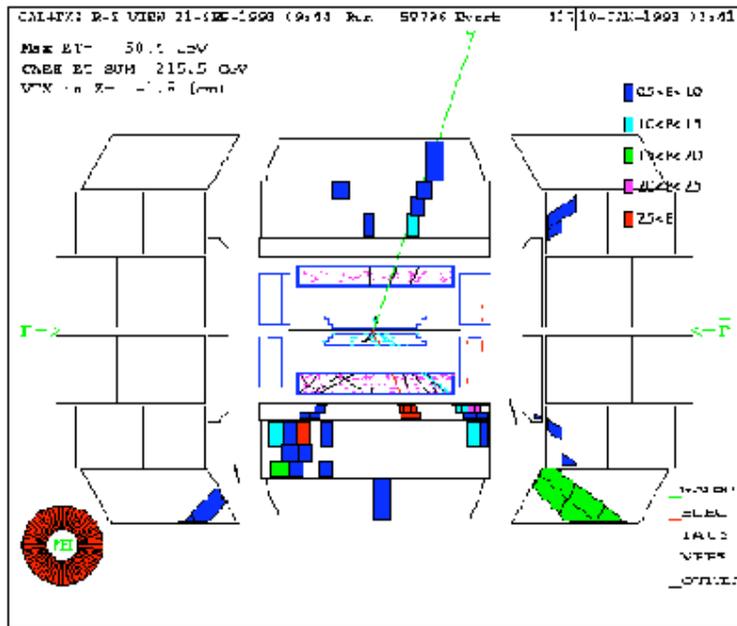
Results of the Initial Search

m_t (GeV/ c^2)	$e\mu$	ee	$e + \text{jets}$	$\mu + \text{jets}$	All	
90	$\epsilon \times B(\%)$	0.39 ± 0.10	0.16 ± 0.02	0.28 ± 0.08	0.15 ± 0.07	22.9 ± 3.6
	$\langle N \rangle$	9.4 ± 2.6	4.0 ± 0.8	6.8 ± 2.1	2.7 ± 1.3	
100	$\epsilon \times B(\%)$	0.46 ± 0.11	0.20 ± 0.03	0.44 ± 0.12	0.19 ± 0.08	17.0 ± 2.7
	$\langle N \rangle$	6.3 ± 1.7	2.8 ± 0.5	6.0 ± 1.8	1.9 ± 0.9	
120	$\epsilon \times B(\%)$	0.49 ± 0.12	0.26 ± 0.04	1.13 ± 0.22	0.61 ± 0.20	12.3 ± 1.7
	$\langle N \rangle$	2.6 ± 0.7	1.4 ± 0.3	5.9 ± 1.3	2.4 ± 0.8	
140	$\epsilon \times B(\%)$	0.54 ± 0.13	0.28 ± 0.04	1.45 ± 0.19	0.90 ± 0.27	6.7 ± 0.8
	$\langle N \rangle$	1.2 ± 0.3	0.6 ± 0.1	3.3 ± 0.6	1.6 ± 0.5	
160	$\epsilon \times B(\%)$	0.56 ± 0.14	0.29 ± 0.04	1.69 ± 0.18	0.85 ± 0.24	3.5 ± 0.4
	$\langle N \rangle$	0.6 ± 0.2	0.3 ± 0.1	1.9 ± 0.3	0.7 ± 0.2	
Physics background		0.5 ± 0.2	0.2 ± 0.1	2.1 ± 1.1	1.1 ± 0.7	3.9 ± 1.3
Fake background		0.6 ± 0.3	0.3 ± 0.1	0.3 ± 0.3	0.3 ± 0.1	1.5 ± 0.4
Total background		1.1 ± 0.4	0.5 ± 0.2	2.4 ± 1.3	1.4 ± 0.9	5.4 ± 1.4
$\int \mathcal{L} dt$ (pb $^{-1}$)		13.5 ± 1.6	13.5 ± 1.6	13.5 ± 1.6	9.8 ± 1.2	
Data		1	1	1	0	3

- Not even a hint of top
 - Masses less than 128 GeV ruled out at 95% C.L.
- Well, maybe there was *one* hint of top...

Event 417

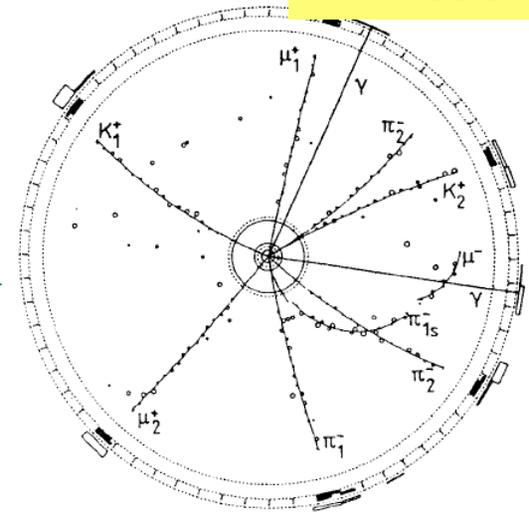
- Candidate in the $e\mu$ channel with large electron and muon energies, and large missing E_T



How Do We Discover?

ARGUS

- Many discoveries in particle physics have arisen from a single spectacular event
 - *B* meson mixing, for example



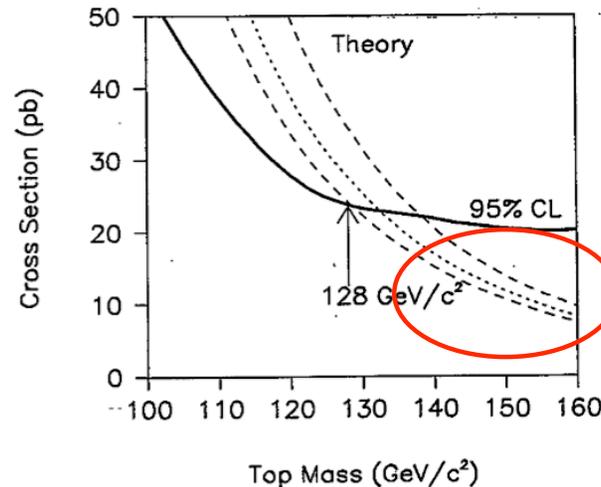
- Was Event 417 spectacular enough?
 - intense discussion within DØ
 - the event couldn't arise from the most common background

$$Z + \text{jets} \rightarrow \tau\tau + \text{jets} \rightarrow e\mu + \text{jets} + E_T$$

- but maybe $WW + \text{jets} \rightarrow e\mu + \text{jets} + E_T \dots$
- In the end, decided not to make any claims from this event
- Need to collect a set of events that collectively is very unlikely to arise from background
 - where very unlikely means $O(10^{-6})$ probability

Moving to Higher Mass

- We now knew the top wasn't less than 130 GeV
- That's a problem:



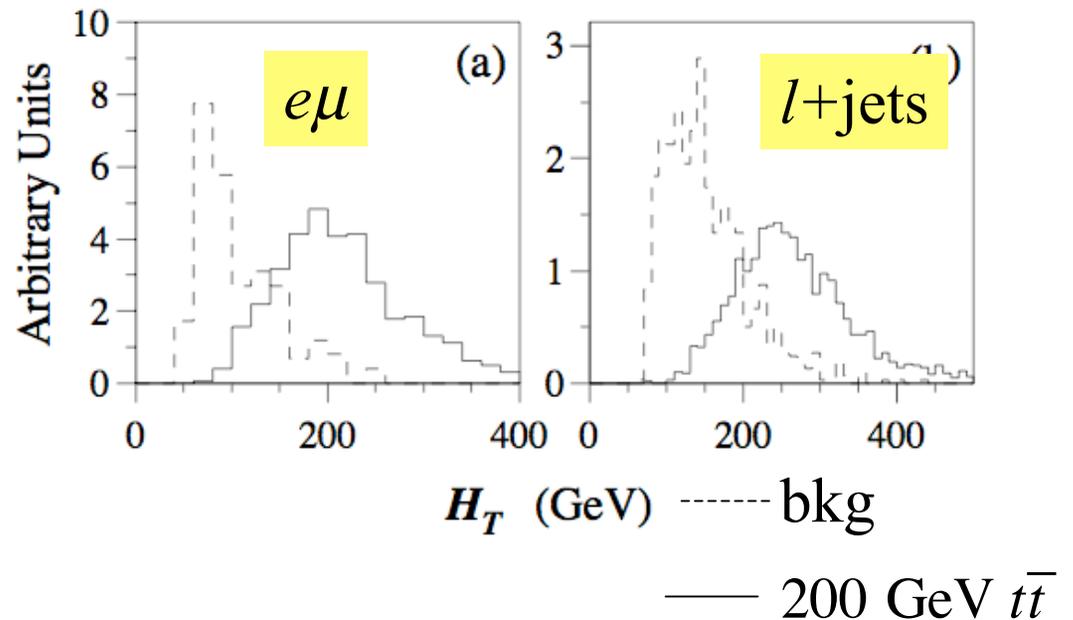
Not many events produced!

- But also an opportunity
 - the heavier the top is, the easier it is to distinguish from background
- So selection was re-optimized to look for high-mass top

How To Explore Higher Masses

- A higher top mass means that higher $\sqrt{\hat{s}}$ is required
 - more energy in $t\bar{t}$ events
 - typically shows up as higher jet transverse energies
- So the transition to searching at high mass meant requiring that either individual jet E_T 's be higher, or that their sum (H_T) be higher

- H_T turned out to be an excellent variable
- we could trade a modest amount of $t\bar{t}$ efficiency for a large reduction in background

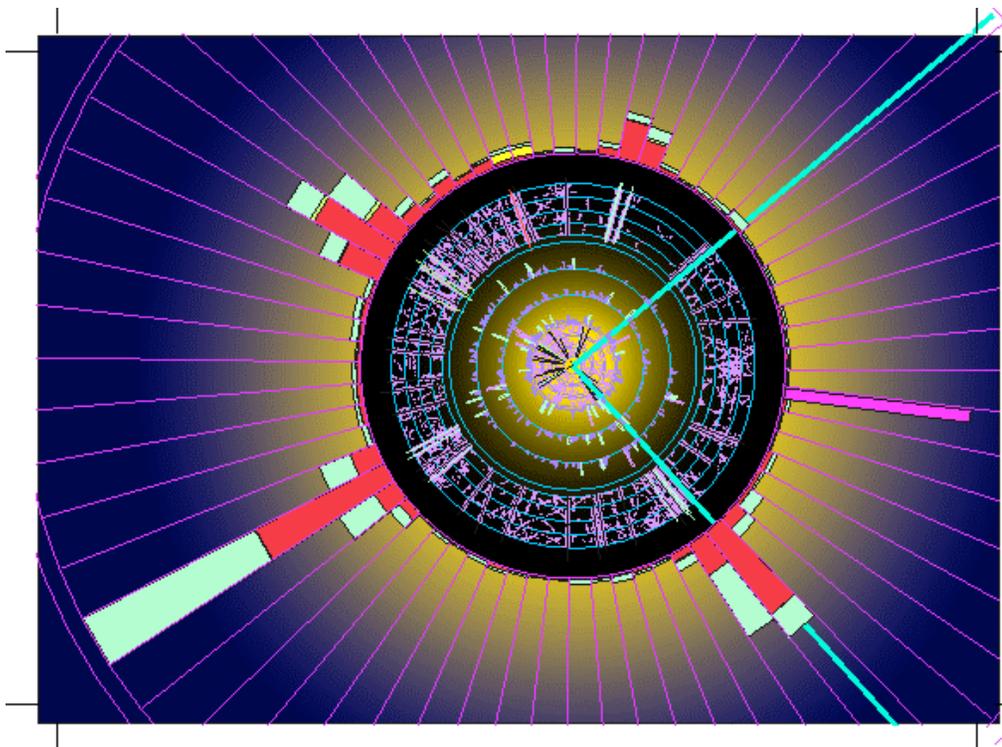


Adding More Channels

- We also increased our sensitivity to top by looking in three more decay channels:
 - Dimuon
 - Here the main challenge is limiting the Z background with poor muon momentum resolution
 - Cut on difference in dimuon momentum and missing E_T directions
 - b -tagged e +jets and μ +jets
 - But DØ didn't have a silicon tracker...

Tagging b 's Without a Vertex Detector

- While DØ couldn't see the displaced b decay vertices, it could find muons from semileptonic b decay



- Thick calorimeter and muon toroid made it essentially impossible for other charged particles to mimic this signature
- But only $\sim 20\%$ of b 's have a muon in their decay
- Means that we could tag, at most, $\sim 40\%$ of events
 - with detector efficiency, down to $\sim 20\%$
 - but reduced background allows one to open up the kinematic selection

The Evidence Mounts

- With the new selection, we find (same 15pb⁻¹ data set):

m_t (GeV/c ²)	$e\mu$ + jets	ee + jets	$\mu\mu$ + jets	e + jets	μ + jets	e + jets/ μ	μ + jets/ μ	All
140	$\epsilon \times B(\%)$	0.31 ± 0.04	0.18 ± 0.02	0.15 ± 0.02	1.1 ± 0.3	0.8 ± 0.2	0.6 ± 0.2	0.4 ± 0.1
	$\langle N \rangle$	0.72 ± 0.12	0.41 ± 0.07	0.25 ± 0.04	2.5 ± 0.7	1.3 ± 0.4	1.4 ± 0.5	0.7 ± 0.2
160	$\epsilon \times B(\%)$	0.36 ± 0.05	0.20 ± 0.03	0.15 ± 0.02	1.5 ± 0.3	1.1 ± 0.3	0.9 ± 0.2	0.5 ± 0.1
	$\langle N \rangle$	0.40 ± 0.08	0.22 ± 0.04	0.12 ± 0.02	1.7 ± 0.5	0.9 ± 0.3	1.0 ± 0.3	0.4 ± 0.1
180	$\epsilon \times B(\%)$	0.39 ± 0.05	0.21 ± 0.03	0.14 ± 0.02	1.6 ± 0.3	1.1 ± 0.3	1.1 ± 0.2	0.7 ± 0.1
	$\langle N \rangle$	0.23 ± 0.04	0.12 ± 0.02	0.06 ± 0.01	0.9 ± 0.3	0.5 ± 0.1	0.6 ± 0.1	0.3 ± 0.1
200	$\epsilon \times B(\%)$	0.40 ± 0.05	0.30 ± 0.04	0.14 ± 0.02	1.8 ± 0.4	1.3 ± 0.3	1.4 ± 0.1	0.8 ± 0.2
	$\langle N \rangle$	0.12 ± 0.02	0.09 ± 0.02	0.03 ± 0.01	0.5 ± 0.1	0.3 ± 0.1	0.4 ± 0.1	0.2 ± 0.1
Background	0.27 ± 0.14	0.15 ± 0.11	0.33 ± 0.06	1.3 ± 0.7	0.7 ± 0.5	0.6 ± 0.2	0.4 ± 0.1	3.8 ± 0.9
$\int \mathcal{L} dt$ (pb ⁻¹)	13.5 ± 1.6	13.5 ± 1.6	9.8 ± 1.2	13.5 ± 1.6	9.8 ± 1.2	13.5 ± 1.6	9.8 ± 1.2	
Data	1	0	0	2	2	2	2	9

- More events than in “all-mass” search?
- Due to lowering aplanarity cut from 0.10 to 0.05 (2 events) and lowering muon p_T cut from 15 to 12 GeV
 - so these just missed being seen in the initial search

The 50pb⁻¹ Sample

- Now we had:
 - calibrated the detector
 - developed event selections for seven different decay channels
 - demonstrated sensitivity to top production
 - seen an intriguing excess of events in the “high-mass” analysis
- Time to look at the next chunk of data!
 - total of $\sim 50\text{pb}^{-1}$
- Further optimization of selection (applying H_T in all channels, raising cut value in channels where already used) reduced background rate by nearly 4x
 - while retaining about 70% of 180 GeV top
 - cuts chosen to give best *expected* significance for top -- i.e. only look at MC when making the decision

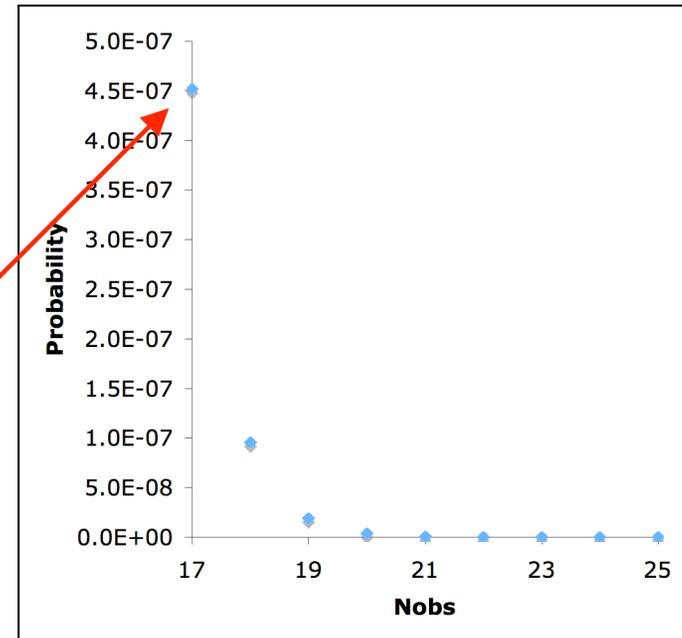
- This time, we found:

m_τ (GeV/ c^2)	$e\mu + \text{jets}$	$ee + \text{jets}$	$\mu\mu + \text{jets}$	$e + \text{jets}$	$\mu + \text{jets}$	$e + \text{jets}/\mu$	$\mu + \text{jets}/\mu$	All
140 $\epsilon \times \mathcal{B}$ (%)	0.17 ± 0.02	0.11 ± 0.02	0.06 ± 0.01	0.50 ± 0.10	0.33 ± 0.08	0.36 ± 0.07	0.20 ± 0.05	
$\langle N \rangle$	1.36 ± 0.21	1.04 ± 0.19	0.46 ± 0.08	4.05 ± 0.94	2.47 ± 0.68	2.93 ± 0.68	1.48 ± 0.42	13.80 ± 2.07
160 $\epsilon \times \mathcal{B}$ (%)	0.24 ± 0.02	0.15 ± 0.02	0.09 ± 0.02	0.80 ± 0.10	0.57 ± 0.13	0.50 ± 0.08	0.25 ± 0.06	
$\langle N \rangle$	0.94 ± 0.13	0.69 ± 0.12	0.34 ± 0.07	3.13 ± 0.54	2.04 ± 0.53	1.95 ± 0.39	0.92 ± 0.24	10.01 ± 1.41
180 $\epsilon \times \mathcal{B}$ (%)	0.28 ± 0.02	0.17 ± 0.02	0.10 ± 0.02	1.20 ± 0.30	0.76 ± 0.17	0.56 ± 0.09	0.35 ± 0.08	
$\langle N \rangle$	0.57 ± 0.07	0.40 ± 0.07	0.19 ± 0.04	2.42 ± 0.67	1.41 ± 0.36	1.14 ± 0.22	0.64 ± 0.16	6.77 ± 1.09
200 $\epsilon \times \mathcal{B}$ (%)	0.31 ± 0.02	0.20 ± 0.03	0.11 ± 0.02	1.70 ± 0.20	0.96 ± 0.21	0.74 ± 0.11	0.41 ± 0.08	
$\langle N \rangle$	0.34 ± 0.04	0.25 ± 0.05	0.11 ± 0.02	1.84 ± 0.31	0.95 ± 0.24	0.81 ± 0.16	0.41 ± 0.10	4.71 ± 0.66
Background	0.12 ± 0.03	0.28 ± 0.14	0.25 ± 0.04	1.22 ± 0.42	0.71 ± 0.28	0.85 ± 0.14	0.36 ± 0.08	3.79 ± 0.55
$\int \mathcal{L} dt$ (pb $^{-1}$)	47.9 ± 5.7	55.7 ± 6.7	44.2 ± 5.3	47.9 ± 5.7	44.2 ± 5.3	47.9 ± 5.7	44.2 ± 5.3	
Data	2	0	1	5	3	3	3	17

- Excess over background in almost every channel

Significance

- So there's clearly an excess over background
 - and it appears unlikely to be a fluctuation
 - *but how unlikely?*
- Easy to find the Poisson probability for the 3.8 expected background events to fluctuate to 17: 4.5×10^{-7} !
- Probability of 3.8 fluctuating to 17 or more is 5.7×10^{-7}
- But there's more
 - there's also the uncertainty on the expected background to consider
 - *almost entirely a systematic uncertainty*



Systematics

- Understanding the systematics is the major component of any analysis
- We had to deal with at least three major classes:

1. What does background look like?
 - theoretical understanding of object kinematics
 - vary parameters in MC
 - cross-check with data when possible

2. How does $D\emptyset$ respond?
 - trigger efficiency
 - jet energy scale and resolution
 - lepton ID efficiencies
 - missing E_T

3. How much data is there?
 - luminosity uncertainty

- By discovery paper, total uncertainty was at 15% level

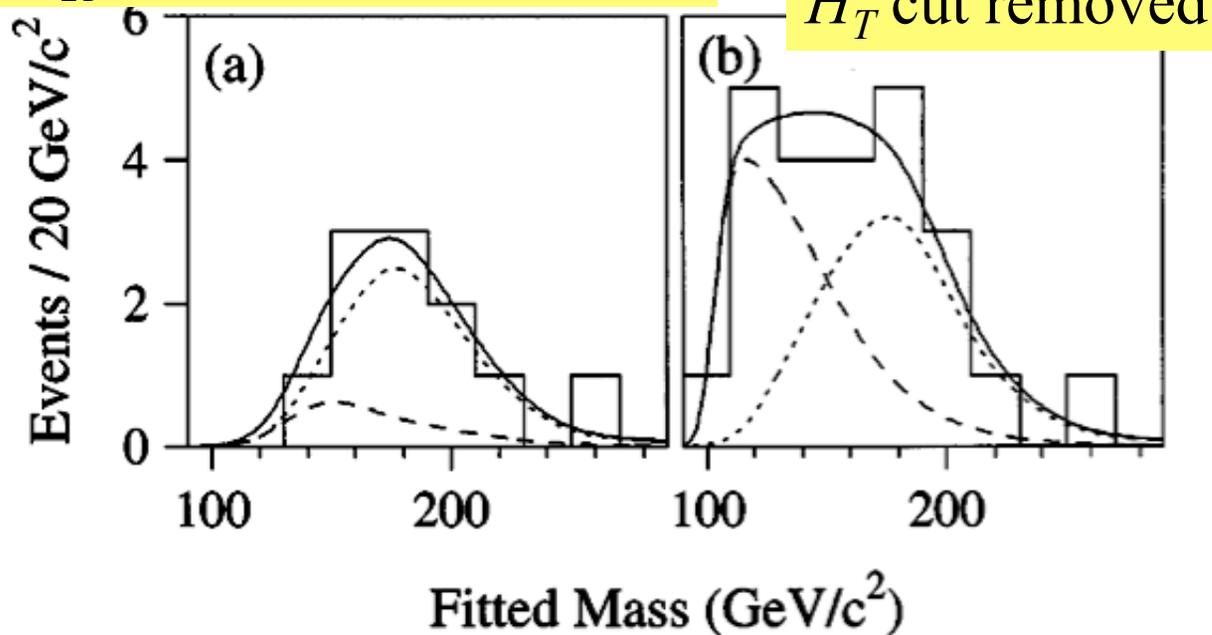
Accounting for the Systematics

- Solution: MC simulation of the background fluctuations
 - throw a Gaussian random number centered at 3.8 with width 0.6 to get expected background
 - then throw a Poisson-distributed random number based on expected background
 - count fraction of tries that result in 17 or more observed events
- Answer is 2×10^{-6} -- **Discovery!**
- Note that this assumes that all the uncertainties in the background estimate are Gaussian
 - fine for statistical component
 - not as obvious for systematics
- That's the real reason we don't trust signal probabilities until they get down to the 10^{-6} range

Circumstantial Evidence

- Beyond merely counting events, one can also look at the kinematics of the events
- The reconstructed mass in 1+jets events is the most interesting single quantity*:

199^{+19}_{-21} (stat.) ± 22 (syst.) GeV

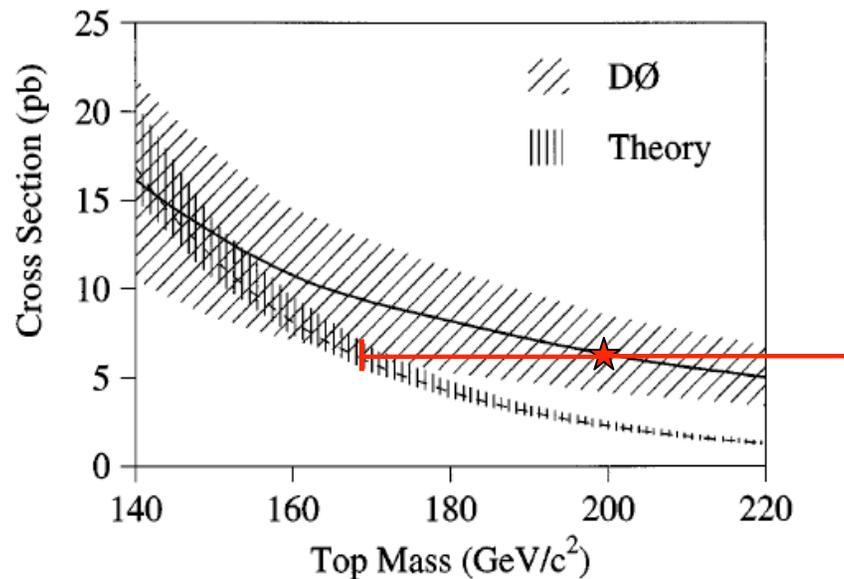


- Data looks like sum of background and $t\bar{t}$
- Not obvious how to incorporate this into overall significance

- In the end, this was taken as additional reassurance that we had top quarks in our sample

*EWV looked into mass value for dilepton events. Methods not sufficiently developed to say anything meaningful

- Best-fit top mass value was 199_{-21}^{+19} (stat.) ± 22 (syst.) GeV
- Above the previous limits, but also an indication that luck played a role in the discovery:



- More events than expected from top across range consistent with mass value

- Also looked at distribution of events across decay channels
 - highly consistent with expectation from top

Summary

- Discovering the top quark required
 - a well-calibrated detector
 - innovative methods of separating signal from background
 - combination of results from many decay channels
 - no one channel would have done it
 - committed effort from a large group of physicists
 - and a bit of luck to see it in 50 pb^{-1}
- Provided me a thorough (and intense) introduction to what's required to do a physics analysis
- Seeing something no one has seen before is definitely exciting
- But it would be even better to find something unexpected